

TWO DIMENSIONAL LAPLACE TRANSFORMS OF GENERALIZED HYPERGEOMETRIC FUNCTIONS

R.S. DAHIYA

Department of Mathematics
Iowa State University
Ames, Iowa 50011

I.H. JOWHAR

Department of Mathematics
Florida State University
Tallahassee, Florida 32306

(Received October 16, 1986)

ABSTRACT. The object of this paper is to obtain new operational relations between the original and the image functions that involve generalized hypergeometric G-functions.

KEY WORDS AND PHRASES: Multidimensional Laplace transforms, Meijer's G-function.

1980 AMS SUBJECT CLASSIFICATION CODE: 44A30, 33A35

1. INTRODUCTION.

The integral equation

$$\phi(p, q) = pq \int_0^\infty \int_0^\infty \exp(-px - qy) f(x, y) dy dx, \quad \operatorname{Re}(p, q) > 0 \quad (1.1)$$

represents the classical Laplace transform of two variables and the functions $\phi(p, q)$ and $f(x, y)$ related by (1.1), are said to be operationally related to each other. $\phi(p, q)$ is called the image and $f(x, y)$ the original.

Symbolically we can write

$$\phi(p, q) \overset{..}{=} f(x, y) \quad \text{or} \quad f(x, y) \overset{..}{=} \phi(p, q), \quad (1.2)$$

and the symbol $\overset{..}{=}$ is called "operational".

Meijer's G-function [5] is defined by a Mellin-Barnes type integral

$$G_{u,v}^{m,n}(z \mid \begin{matrix} a_u \\ b_v \end{matrix}) = \frac{1}{2\pi i} \int_L \frac{\prod_{j=1}^m \Gamma(b_j - s) \prod_{j=1}^n \Gamma(1 - a_j + s)}{\prod_{j=m+1}^u \Gamma(1 - b_j + s) \prod_{j=n+1}^v \Gamma(a_j - s)} z^s ds \quad (1.3)$$

where m, n, u, v are integers with $v > 1$; $0 < n < u$; $0 < m < v$, the parameters a_j and b_j are such that no poles of $\Gamma(b_j - s)$; $j = 1, 2, \dots, m$ coincides with any pole of $\Gamma(1 - a_k + s)$; $k = 1, 2, \dots, n$. Thus $(a_k - b_j)$ is not a positive integer. The path L goes from $-i\infty$ to $+i\infty$ so that all poles of integrand must be simple and those of $\Gamma(b_j - s)$; $j = 1, 2, \dots, m$ lie on one side of the contour L and those of $\Gamma(1 - a_k + s)$; $k = 1, 2, \dots, n$ must lie on the other side. The integrand converges if $u + v < 2(m + n)$ and $|\arg z| < (m + n - \frac{1}{2}u - \frac{1}{2}v)\pi$. For sake of brevity a_u denotes $a_1, a_2, \dots, a_u$.

In the present paper, we propose to establish a couple of formulae for

calculating Laplace transform pairs of two dimensions that involve Meijer's G-function.

2. THE MAIN RESULTS.

- (i) $\bar{\delta} = m + n - \frac{1}{2}(u+v) > 0$, $|\arg \alpha| < \bar{\delta}\pi$,
- (ii) $0 < n < u$, $0 < m < v$, $v > 1$,
- (iii) $\operatorname{Re}(b_j + \xi) > 0$, $j = 1, 2, \dots, m$
- (iv) $\operatorname{Re}(a_k + \xi - \frac{\sigma}{2} - \frac{r}{2}) < 0$, $k = 1, 2, \dots, n$,
- (v) $a_k - b_k$ is not a positive integer, $j = 1, 2, \dots, m$, $k = 1, 2, \dots, n$,
- (vi) r represents the non-negative integer, $0, 1, 2, 3, \dots$, then

$$\begin{aligned} x^r (xy)^{\sigma-\xi-1} G_{u+2,v}^{m,n+1} \left(\frac{\alpha}{xy} \middle| \begin{matrix} 1-\xi, a_u, \sigma+r-\xi \\ b_v \end{matrix} \right) \\ = p^{-r} (pq)^{\xi-\sigma+1} G_{u+1,v+1}^{m+1,n+1} \left(\alpha pq \middle| \begin{matrix} 1-\xi, a_u \\ \sigma-\xi, b_v \end{matrix} \right) \end{aligned} \quad (2.1)$$

and

$$\begin{aligned} x^{\delta-r-1} \left(\frac{x}{y} \right)^{\xi-\sigma+1} G_{u+2,v}^{m,n+1} \left(\frac{\alpha x}{y} \middle| \begin{matrix} 1-\xi, a_u, \sigma+r-\xi \\ b_v \end{matrix} \right) \\ = p^{r-\delta-1} \left(\frac{q}{p} \right)^{\xi-\sigma+1} G_{u+3,v+1}^{m+1,n+2} \left(\frac{\alpha q}{p} \middle| \begin{matrix} \sigma-\xi+r-\delta, 1-\xi, a_u, \sigma-\xi+r \\ \sigma-\xi, b_v \end{matrix} \right) \end{aligned} \quad (2.2)$$

valid under the conditions:

- (a) $\operatorname{Re}(\delta + \xi - \sigma - r + b_j) > -1$, $j = 1, 2, \dots, m$,
- (b) $\operatorname{Re}(\frac{\delta}{2} + \xi - \sigma - r + a_k) < \frac{1}{4}$, $k = 1, 2, \dots, n$,
- (c) along with (i), (ii), (v) and (vi).

From (2.1) and (2.2), we propose to prove the following relations.

$$\begin{aligned} x^r (xy)^{\sigma-1} E(b_1, \dots, b_v : a_1, \dots, a_u, \sigma+r : \frac{\alpha}{xy}) \\ = p^{-r} (pq)^{1-\sigma} E(\sigma, b_1, \dots, b_v : a_1, \dots, a_u : \alpha pq), \end{aligned} \quad (2.3)$$

where $\operatorname{Re}(\sigma) > 0$, $v > u+1$ and r is a positive integer.

$$\begin{aligned} x^{\delta-r-1} \left(\frac{x}{y} \right)^{1-\sigma} E(\xi+b_1, \dots, \xi+b_v : \xi+a_1, \dots, \xi+a_u, \sigma+r : \frac{x}{y}) \\ = pq^{r-\delta} E(1+\delta-r, 1+\xi+\delta-\sigma-r+b_1, \dots, 1+\xi+\delta-\sigma-r+b_v : \\ 2+\delta-\sigma-r, 1+\xi+\delta-\sigma-r+a_1, \dots, 1+\xi+\delta-\sigma-r+a_u, 1+\delta : \frac{q}{p}) \end{aligned} \quad (2.4)$$

valid under the same conditions as (2.3).

The function appearing in (2.3) and (2.4) is MacRobert's E-function, whose

properties are given in [6], pp. 433-434, and [8].

PROOF: The generalized Stieltjes transform of a G-function is given by (see [6], p. 237)

$$\int_0^{\infty} \frac{x^{\xi-1}}{(x+\beta)^{\sigma}} G_{p,q}^{m,n} \left(\alpha x \mid \begin{matrix} a_p \\ b_q \end{matrix} \right) dx = \frac{\beta^{\xi-\sigma}}{\Gamma(\sigma)} G_{p+1,q+1}^{m+1,n+1} \left(\alpha\beta \mid \begin{matrix} 1-\xi, a_p \\ \sigma-\xi, b_q \end{matrix} \right) \quad (2.5)$$

On writing pq for β and multiplying both the sides of (2.5) by $p^{1-r} q$, we have

$$\begin{aligned} \int_0^{\infty} \frac{p^{1-r} q}{(pq+t)^{\sigma}} \cdot t^{\xi-1} G_{u,v}^{m,n} \left(\alpha t \mid \begin{matrix} a_u \\ b_v \end{matrix} \right) dt \\ = \frac{(pq)^{\xi-\sigma+1} p^{-r}}{\Gamma(\sigma)} G_{u+1,v+1}^{m+1,n+1} \left(\alpha pq \mid \begin{matrix} 1-\xi, a_u \\ \sigma-\xi, b_v \end{matrix} \right) \end{aligned} \quad (2.6)$$

Now interpreting with the help of the known result ([4], result (2.83), p. 137), it follows

$$\begin{aligned} \frac{y^{\sigma-1}}{\Gamma(\sigma)} \left(\frac{x}{y} \right)^{\frac{\sigma+r-1}{2}} \int_0^{\infty} t^{\xi-\frac{\sigma}{2}-\frac{r}{2}-\frac{1}{2}} J_{\sigma+r-1}(2\sqrt{txy}) G_{u,v}^{m,n} \left(\alpha t \mid \begin{matrix} a_u \\ b_v \end{matrix} \right) dt \\ = \frac{(pq)^{\xi-\sigma+1}}{\Gamma(\sigma) p^r} G_{u+1,v+1}^{m+1,n+1} \left(\alpha pq \mid \begin{matrix} 1-\xi, a_u \\ \sigma-\xi, b_v \end{matrix} \right). \end{aligned} \quad (2.7)$$

Evaluating the left-hand side integral of (2.7), we get

$$\begin{aligned} x^r (xy)^{\sigma-\xi-1} G_{u+2,v}^{m,n+1} \left(\frac{\alpha}{xy} \mid \begin{matrix} 1-\xi, a_u, \sigma+r-\xi \\ b_v \end{matrix} \right) \\ = p^{-r} (pq)^{\xi-\sigma-1} G_{u+1,v+1}^{m+1,n+1} \left(\alpha pq \mid \begin{matrix} 1-\xi, a_u \\ \sigma-\xi, b_v \end{matrix} \right). \end{aligned} \quad (2.8)$$

The following result will be used in the proof of (2.2).

If $F(p,q) \stackrel{..}{=} f(x,y)$, then

$$x^{\delta-1} f\left(\frac{1}{x}, y\right) \stackrel{..}{=} \int_0^{\infty} \left(\frac{\lambda}{p}\right)^{\frac{\delta}{2}-1} J_{\delta}(2\sqrt{p\lambda}) F(\lambda, q) d\lambda. \quad (2.9)$$

From (2.8) and (2.9), we have

$$\begin{aligned} x^{\delta-r-\sigma+\xi} y^{\sigma-\xi-1} G_{u+2,v}^{m,n+1} \left(\frac{\alpha x}{y} \mid \begin{matrix} 1-\xi, a_u, \sigma+r-\xi \\ b_v \end{matrix} \right) \\ = \frac{q^{\xi-\sigma+1}}{p^{\delta/2-1}} \int_0^{\infty} \lambda^{\frac{\delta}{2}+\xi-\sigma-r} J_{\delta}(2\sqrt{p\lambda}) G_{u+1,v+1}^{m+1,n+1} \left(\alpha q \lambda \mid \begin{matrix} 1-\xi, a_u \\ \sigma-\xi, b_v \end{matrix} \right) d\lambda. \end{aligned} \quad (2.10)$$

On evaluating the right hand side integral and after some simplification, we obtain the desired result (2.2).

In (2.1) and (2.2), reducing the Meijer's G-function to MacRobert's E-function to obtain (2.3) and (2.4).

$$p^{-r}(pq)^{\frac{3}{2}-\sigma} \frac{H^{(1)}_{\sigma-\frac{1}{2}}(\sqrt{pq})}{\sigma-\frac{1}{2}} \frac{H^{(2)}(\sqrt{pq})}{\sigma-\frac{1}{2}} \stackrel{..}{=} 2\pi^{-5/2} \cos\left(\sigma - \frac{1}{2}\right)\pi x^r(xy)^{\sigma-\xi-1} \cdot G_{2,2}^{2,1}\left(\frac{1}{xy} \mid \begin{matrix} 1-\xi, \sigma+r-\xi \\ 1-\sigma-\xi, \frac{1}{2}-\xi \end{matrix} \right) \quad (3.6)$$

$$\begin{aligned}
& p^{-r} (pq)^{\frac{1}{2}-m} W_{k,m}(2i\sqrt{pq}) W_{k,m}(-2i\sqrt{pq}) \\
& \quad = \frac{x^r(xy)^{m+\frac{a}{2}-\frac{1}{2}}}{\sqrt{\pi} \Gamma(\frac{1}{2}-K+m) \Gamma(\frac{1}{2}-K-m)} \\
& \quad \quad \cdot G_{3,3}^{3,1} \left(\frac{1}{xy} \mid \begin{matrix} 1+\frac{a}{2}+K, 1+\frac{a}{2}-K, m+\frac{a}{2}+r+\frac{1}{2} \\ \frac{a+1}{2}-m, \frac{a}{2}+1, \frac{a+1}{2} \end{matrix} \right) \quad (3.7)
\end{aligned}$$

$$\begin{aligned}
& p^{-r} (pq)^{1-\frac{u}{2}-\frac{v}{2}} [e^{i\pi(v-u)/2} H_v^{(1)}(\sqrt{pq}) H_u^{(2)}(\sqrt{pq}) \\
& \quad - e^{i\pi(u-v)/2} H_u^{(1)}(\sqrt{pq}) H_v^{(2)}(\sqrt{pq})] \\
& \quad = 2\pi^{-5/2} (\cos u\pi - \cos v\pi) x^r (xy)^{\frac{1}{2}(u+v+w)-1} \\
& \quad \quad \cdot G_{3,3}^{3,1} \left(\frac{1}{xy} \mid \begin{matrix} \frac{w}{2}, \frac{w+1}{2}, \frac{w+u+v}{2}+r \\ \frac{w-u+v}{2}, \frac{w+u-v}{2}, \frac{w-u-v}{2} \end{matrix} \right) \quad (3.8)
\end{aligned}$$

The following four results are obtained from (2.1) by using Carlson's results [1]:

$$\begin{aligned}
& p^{1-r} [\psi(1-2a, 1-2c; 2i\sqrt{pq}) + \psi(1-2a, 1-2c; -2i\sqrt{pq})] \\
& \quad = \frac{2^{2c-2a} x^{r-1} y^{-1}}{\sqrt{\pi} \Gamma(1-2a) \Gamma(1+2c-2a)} G_{3,3}^{2,2} \left(\frac{1}{xy} \mid \begin{matrix} a, a+\frac{1}{2}, r \\ c, c+\frac{1}{2}, \frac{1}{2} \end{matrix} \right) \quad (3.9)
\end{aligned}$$

$$\begin{aligned}
& \sqrt{pq} p^{-r} [\psi(1-2a, 1-2c; 2i\sqrt{pq}) - \psi(1-2a, 1-2c; -2i\sqrt{pq})] \\
& \quad = \frac{-i 2^{2c-2a} x^{r-\frac{1}{2}} y^{-\frac{1}{2}}}{\Gamma(1-2a) \Gamma(1+2c-2a)} G_{3,3}^{2,2} \left(\frac{1}{xy} \mid \begin{matrix} a, a+\frac{1}{2}, r+\frac{1}{2} \\ c, c+\frac{1}{2}, 0 \end{matrix} \right) \quad (3.10)
\end{aligned}$$

$$\begin{aligned}
& p^{-r} (pq)^{1+c} [\exp(2i\sqrt{pq} + ic\pi) \Gamma(-2c, 2i\sqrt{pq}) + \exp(-2i\sqrt{pq} - ic\pi) \Gamma(-2c, -2i\sqrt{pq})] \\
& \quad = \pi^{-\frac{1}{2}} [\Gamma(1+2c)]^{-1} x^{r-1} y^{-1} G_{2,2}^{2,1} \left(\frac{1}{xy} \mid \begin{matrix} 0, r \\ c, c+\frac{1}{2} \end{matrix} \right) \quad (3.11)
\end{aligned}$$

where $\Gamma(a, x)$ denotes incomplete gamma function.

$$\begin{aligned}
& p^{-r} [\psi(2, 2-2c; 2i\sqrt{pq}) - \psi(2, 2-2c; -2i\sqrt{pq})] \\
& \quad = -i\pi^{-\frac{1}{2}} 2^{2c} [\Gamma(1+2c)]^{-1} x^r G_{2,2}^{2,1} \left(\frac{1}{xy} \mid \begin{matrix} 0, 1+r \\ c, c+\frac{1}{2} \end{matrix} \right) \quad (3.12)
\end{aligned}$$

In the last six relations, $G_{3,3}^{3,1}$ or $G_{3,3}^{2,2}$ provides an interpretation for the symbol ${}_3F_2$, as does ψ for ${}_2F_0$. So does $G_{2,2}^{2,1}$ or $G_{2,2}^{1,1}$ for ${}_2F_1$.

3(b). The G-function expressed as a named original function.

The following results are obtained from (2.1) by using the known results [7] on pages 228-234.

$$\begin{aligned}
 & p^{\sigma-2K-1} (pq)^{K-\sigma-\frac{a}{2}} G_{1,5}^{4,1} \left(pq \mid \begin{matrix} \frac{a}{2}-K+1 \\ \sigma-K+\frac{a}{2}, \frac{a+1}{2}+m, \frac{a}{2}+1, \frac{a+1}{2}, \frac{a+1}{2}-m \end{matrix} \right) \\
 & \quad = \frac{\sqrt{\pi}}{\Gamma(2m+1)} \Gamma\left(\frac{1}{2}+K+m\right) x^{2K-\sigma+1} (xy)^{\sigma-K-1} M_{K,m}\left(\frac{2}{\sqrt{xy}}\right) W_{-K,m}\left(\frac{2}{\sqrt{xy}}\right) \quad (3.13)
 \end{aligned}$$

$$\begin{aligned}
 & p^{2K+\sigma-1} (pq)^{1-\sigma-K-\frac{a}{2}} G_{1,5}^{5,1} \left(pq \mid \begin{matrix} \frac{a}{2}+K+1 \\ \sigma+K+\frac{a}{2}, \frac{a+1}{2}+m, \frac{a+1}{2}-m, \frac{a}{2}+1, \frac{a+1}{2} \end{matrix} \right) \\
 & \quad = \sqrt{\pi} \Gamma\left(\frac{1}{2}-K+m\right) \Gamma\left(\frac{1}{2}-K-m\right) x^{1-\sigma-2K} (xy)^{\sigma+K-1} \cdot W_{K,m}\left(\frac{2i}{\sqrt{xy}}\right) W_{K,m}\left(\frac{-2i}{\sqrt{xy}}\right) \quad (3.14)
 \end{aligned}$$

$$\begin{aligned}
 & p^{\sigma-\frac{3}{2}} (pq)^{2-\sigma-\frac{w}{2}} G_{1,5}^{5,1} \left(pq \mid \begin{matrix} \frac{w}{2} \\ \sigma+\frac{w}{2}-1, \frac{w+u+v}{2}, \frac{w-u+v}{2}, \frac{w+u-v}{2}, \frac{w-u-v}{2} \end{matrix} \right) \\
 & \quad = \frac{-i\pi^{5/2}}{2(\cos u\pi - \cos v\pi)} x^{\frac{3}{2}-\sigma} (xy)^{\sigma-2} \cdot [e^{i\pi(v-u)/2} H_v^{(1)}\left(\frac{1}{\sqrt{xy}}\right) H_u^{(2)}\left(\frac{1}{\sqrt{xy}}\right) \\
 & \quad \quad - e^{i\pi(u-v)/2} H_u^{(1)}\left(\frac{1}{\sqrt{xy}}\right) H_v^{(2)}\left(\frac{1}{\sqrt{xy}}\right)] \quad (3.15)
 \end{aligned}$$

$$\begin{aligned}
 & p^{\sigma-\frac{1}{2}} (pq)^{\frac{3}{2}-\frac{w}{2}-\sigma} G_{1,5}^{5,1} \left(pq \mid \begin{matrix} \frac{w+1}{2} \\ \sigma+\frac{w-1}{2}, \frac{w+u+v}{2}, \frac{w-u+v}{2}, \frac{w+u-v}{2}, \frac{w-u-v}{2} \end{matrix} \right) \\
 & \quad = \frac{\frac{5}{2}\pi \frac{1}{2}-\sigma}{2(\cos u\pi + \cos v\pi)} (xy)^{\sigma-\frac{3}{2}} [e^{i\pi(u-v)/2} H_v^{(1)}\left(\frac{1}{\sqrt{xy}}\right) H_u^{(2)}\left(\frac{1}{\sqrt{xy}}\right) \\
 & \quad \quad + e^{i\pi(u-v)/2} H_u^{(1)}\left(\frac{1}{\sqrt{xy}}\right) H_v^{(2)}\left(\frac{1}{\sqrt{xy}}\right)] \quad (3.16)
 \end{aligned}$$

$$\begin{aligned}
 & p^{\sigma-\frac{3}{4}} (pq)^{2-a-\sigma} G_{1,5}^{3,2} \left(pq \mid \begin{matrix} a \\ \sigma+a-b, b, c, 2a-c, 2a-b \end{matrix} \right) \\
 & \quad = 2\sqrt{\pi} x^{\frac{3}{2}-\sigma} (xy)^{\sigma-2} I_{b+c+2a}\left(\frac{1}{\sqrt{xy}}\right) K_{b-c}\left(\frac{1}{\sqrt{xy}}\right) \quad (3.17)
 \end{aligned}$$

$$\begin{aligned}
& p^{\sigma-\frac{3}{2}} (pq)^{2-\sigma} G_{1,5}^{5,1} (pq \mid \begin{matrix} 0 \\ \sigma-1, a, b, -b, -a \end{matrix}) \\
& \approx \frac{i \pi^{5/2}}{4 \sin a \pi \sin b \pi} x^{\frac{3}{2}-\sigma} (xy)^{\sigma-2} \cdot [e^{-i \pi b} H_{a-b}^{(1)}(\frac{1}{\sqrt{xy}}) H_{a+b}^{(2)}(\frac{1}{\sqrt{xy}}) \\
& \quad - e^{i \pi b} H_{a+b}^{(1)}(\frac{1}{\sqrt{xy}}) H_{a-b}^{(2)}(\frac{1}{\sqrt{xy}})] \quad (3.18)
\end{aligned}$$

$$\begin{aligned}
& \sigma^{-\frac{1}{2}} (pq)^{\frac{3}{2}-\sigma} G_{1,5}^{5,1} (pq \mid \frac{1}{2} \\
& \sigma^{-\frac{1}{2}}, a, b, -b, -a \\
& \frac{\pi x^{\frac{5}{2}}}{4 \cos a\pi \cos b\pi} (xy)^{\sigma-\frac{3}{2}} \cdot [e^{-i\pi b} H_{a-b}^{(1)}(\frac{1}{\sqrt{xy}}) H_{a+b}^{(2)}(\frac{1}{\sqrt{xy}}) \\
& - e^{i\pi b} H_{a+b}^{(1)}(\frac{1}{\sqrt{xy}}) H_{a-b}^{(2)}(\frac{1}{\sqrt{xy}})] \quad (3.19)
\end{aligned}$$

$$\begin{aligned} & p^{\sigma+2a-1} (pq)^{\frac{3}{2}-a-\sigma} G_{1,5}^{5,1}(pq \mid \\ & \quad \sigma+a-\frac{1}{2}, 0, \frac{1}{2}, b, -b) \\ & = \sqrt{\pi} \Gamma(\frac{1}{2}+b-a) \Gamma(\frac{1}{2}-b-a) x^{1-2a-\sigma} (xy)^{\sigma+a-1} w_{a,b}\left(\frac{2i}{\sqrt{xy}}\right) w_{a,b}\left(\frac{-2i}{\sqrt{xy}}\right) \end{aligned} \quad (3.20)$$

The following three results are obtained by using Carlson's results [1] on page 239 in (2.1).

$$\begin{aligned}
& p^{\sigma-\frac{3}{2}} (pq)^{2-\sigma} G_{1,5}^{3,1} (pq \mid \begin{matrix} 0 \\ \sigma-1, c, c+\frac{1}{2}, d, d+\frac{1}{2} \end{matrix}) \\
& = \frac{\Gamma(1+2c) \pi^{-\frac{1}{2}} x^{\frac{3}{2}-\sigma}}{\Gamma(1+2c-2d) 2^{2d+1}} (xy)^{\sigma-c-2} [e^{ic\pi} {}_1F_1(1+2c; 1+2c-2d; \frac{-2i}{\sqrt{xy}}) \\
& \quad + e^{-ic\pi} {}_1F_1(1+2c; 1+2c-2d; \frac{2i}{\sqrt{xy}})] \quad (3.21)
\end{aligned}$$

$$\begin{aligned}
& p^{\sigma-\frac{3}{2}} (pq)^{2-\sigma} G_{1,5}^{5,1} \left(\begin{matrix} 0 \\ \sigma-1, c, c+\frac{1}{2}, d, d+\frac{1}{2} \end{matrix} \middle| \right) \\
& = \sqrt{\pi} 2^{-2d} \Gamma(1+2c) \Gamma(1+2d) \cdot x^{\frac{3}{2}-\sigma} (xy)^{\sigma-c-2} [e^{ic\pi} \psi(1+2c, 1+2c-2d; \frac{2i}{\sqrt{xy}}) \\
& \quad + e^{-ic\pi} \psi(1+2c, 1+2c-2d; \frac{-2i}{\sqrt{xy}})] \quad (3.22)
\end{aligned}$$

$$\begin{aligned}
& p^{\frac{5}{2}-\sigma-\delta} \left(\frac{p}{p} \right)^{2-\sigma} G_{3,5}^{3,2} \left(\frac{q}{p} \mid \begin{matrix} \frac{1}{2}-\delta, 0, \frac{1}{2} \\ \sigma-1, c, c+\frac{1}{2}, d, d+\frac{1}{2} \end{matrix} \right) \\
& = \frac{\Gamma(1+2c)}{\Gamma(1+2c-2d)} \frac{\pi^{-1/2}}{2^{2d+1}} x^{c+\delta-\frac{1}{2}} y^{\sigma-c-2} \cdot [e^{ic\pi} {}_1F_1(1+2c; 1+2c-2d; -2i\sqrt{\frac{x}{y}}) \\
& \quad + e^{-ic\pi} {}_1F_1(1+2c; 1+2c-2d; 2i\sqrt{\frac{x}{y}})] \quad (3.28)
\end{aligned}$$

$$\begin{aligned}
& p^{\frac{1}{2}-\delta} q^{2-\sigma} G_{3,5}^{5,2} \left(\frac{q}{p} \mid \begin{matrix} \frac{1}{2}-\delta, 0, \frac{1}{2} \\ \sigma-1, c, c+\frac{1}{2}, d, d+\frac{1}{2} \end{matrix} \right) \\
& = 2^{-2d} \sqrt{\pi} \frac{\Gamma(1+2c)\Gamma(1+2d)}{\Gamma(1+2c-2d)} x^{\delta+c-\frac{1}{2}} y^{\sigma-c-2} \cdot [e^{ic\pi} \psi(1+2c, 1+2c-2d; 2i\sqrt{\frac{x}{y}}) \\
& \quad + e^{-ic\pi} \psi(1+2c, 1+2c-2d; -2i\sqrt{\frac{x}{y}})] \quad (3.29)
\end{aligned}$$

$$\begin{aligned}
& p^{\frac{3}{2}-\sigma-\delta} \left(\frac{q}{p} \right)^{\frac{3}{2}-\sigma} G_{3,5}^{5,2} \left(\frac{q}{p} \mid \begin{matrix} -\delta, \frac{1}{2}, 0 \\ \sigma-\frac{1}{2}, c, c+\frac{1}{2}, d, d+\frac{1}{2} \end{matrix} \right) \\
& = \frac{\sqrt{\pi}i}{2^{2d}} \frac{\Gamma(1+2c)\Gamma(1+2d)}{\Gamma(1+2c-2d)} x^{c+\delta} y^{\sigma-c-\frac{3}{2}} \cdot [e^{ic\pi} \psi(1+2c, 1+2c-2d; 2i\sqrt{\frac{x}{y}}) \\
& \quad - e^{-ic\pi} \psi(1+2c, 1+2c-2d; -2i\sqrt{\frac{x}{y}})] \quad (3.30)
\end{aligned}$$

REFERENCES

- [1] Carlson, B.C., Some extensions of Lardner's relations between ${}_0F_3$ and Bessel functions. SIAM Journal of Mathematical Analysis. Vol. 1, No. 2, May 1970.
- [2] Dahiya, R.S., Computation of two-dimensional Laplace transforms. Rendiconti di Matematica, Univ. of Rome. Vol. (3) 8, Series VI (1975), pp. 805-813.
- [3] Dahiya, R.S. and Egwurube, Michael, Theorems and Applications of two dimensional Laplace transforms. Ukrainian Mathematical Journal (accepted for publication).
- [4] Ditkin, V.A. and Prudnikov, A.P., Operational calculus in two variables and its applications. Pergamon Press (English Edition), (1962).
- [5] Erde'lyi, A., Magnus, W., Oberhettinger, F., and Tricomi, F.G., Higher Transcendental Functions, Vol. 1, McGraw Hill, New York, pp. 207, 209, 215 (1953).
- [6] Erde'lyi, A., Magnus, W., Oberhettinger, F., and Tricomi, F.G., Tables of Integral Transforms, Vol. 2, McGraw Hill, New York (1954).
- [7] Luke, Y.L., The special functions and their approximations, Vol. 1, Academic Press, New York (1969).
- [8] MacRobert, T.M., Functions of a complex variables, 5th edition, London (1962).
- [9] Voelker, D. and Doetsch, G., Die Zweidimensionale Laplace transformation. Verlag Birkhäuser, Basel (1950).

Special Issue on Time-Dependent Billiards

Call for Papers

This subject has been extensively studied in the past years for one-, two-, and three-dimensional space. Additionally, such dynamical systems can exhibit a very important and still unexplained phenomenon, called as the Fermi acceleration phenomenon. Basically, the phenomenon of Fermi acceleration (FA) is a process in which a classical particle can acquire unbounded energy from collisions with a heavy moving wall. This phenomenon was originally proposed by Enrico Fermi in 1949 as a possible explanation of the origin of the large energies of the cosmic particles. His original model was then modified and considered under different approaches and using many versions. Moreover, applications of FA have been of a large broad interest in many different fields of science including plasma physics, astrophysics, atomic physics, optics, and time-dependent billiard problems and they are useful for controlling chaos in Engineering and dynamical systems exhibiting chaos (both conservative and dissipative chaos).

We intend to publish in this special issue papers reporting research on time-dependent billiards. The topic includes both conservative and dissipative dynamics. Papers discussing dynamical properties, statistical and mathematical results, stability investigation of the phase space structure, the phenomenon of Fermi acceleration, conditions for having suppression of Fermi acceleration, and computational and numerical methods for exploring these structures and applications are welcome.

To be acceptable for publication in the special issue of Mathematical Problems in Engineering, papers must make significant, original, and correct contributions to one or more of the topics above mentioned. Mathematical papers regarding the topics above are also welcome.

Authors should follow the Mathematical Problems in Engineering manuscript format described at <http://www.hindawi.com/journals/mpe/>. Prospective authors should submit an electronic copy of their complete manuscript through the journal Manuscript Tracking System at <http://mts.hindawi.com/> according to the following timetable:

Manuscript Due	March 1, 2009
First Round of Reviews	June 1, 2009
Publication Date	September 1, 2009

Guest Editors

Edson Denis Leonel, Department of Statistics, Applied Mathematics and Computing, Institute of Geosciences and Exact Sciences, State University of São Paulo at Rio Claro, Avenida 24A, 1515 Bela Vista, 13506-700 Rio Claro, SP, Brazil; edleonel@rc.unesp.br

Alexander Loskutov, Physics Faculty, Moscow State University, Vorob'evy Gory, Moscow 119992, Russia; loskutov@chaos.phys.msu.ru