

**ON SOLVABILITY OF BOUNDARY VALUE PROBLEM FOR ELLIPTIC EQUATIONS  
WITH BITSADZE-SAMARSKIĬ CONDITION**

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**ABSTRACT.** In this paper we investigate the solvability of a non-local problem for a linear elliptic equation, which is also known as the boundary value problem with the Bitsadze-SamarSKIĬ condition. We prove the existence and uniqueness of a classical solution to this problem. In the final part of this paper we propose an  $L^2$ -approach which gives a rise to weak solutions in a weighted Sobolev space. The crucial point in proving the existence of weak solutions is a suitable modification of the Bitsadze-SamarSKIĬ condition.

**KEY WORDS AND PHRASES.** *Elliptic equations, non-local problems, Bitsadze-SamarSKIĬ conditions.*

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1. **INTRODUCTION.** In recent years several authors have studied the solvability of non-local problems for elliptic and parabolic equations [1-9]. The importance of non-local problems appears to have been first noted in the literature by Bitsadze-SamarSKIĬ. The problem studied in these papers constitutes a direct generalization of the classical boundary value problems. The most significant feature of nonlocal problems is that the boundary condition relates values of a solution on the boundary to its values on some part of the interior of the region. This type of the boundary value problem is often referred to as the boundary value problem with the Bitsadze-SamarSKIĬ condition [7],[8]. The problem (2.1),(2.2) discussed in this article arises from the mathematical description of some processes in a plasma (see paper [6] for full account of physical aspects of non-local problems).

The paper is organized as follows. In Section 2 we give the uniqueness and existence theorem of the classical solutions of the problem (2.1), (2.2). Our method is based on the maximum principle developed in papers [4] and [5]. Section 3 contains a discussion of the solvability of the non-local problem for harmonic functions in a disc in  $R_2$ . The results of this section slightly improve the explicit formulae derived by Bitsadze (see [1] and [2]) for harmonic functions associated with some non-local problem. In a general case of a linear elliptic equation we reduce the problem (2.1), (2.2) to the solvability of the integral equation of the second kind. The final sections 4 and 5 are devoted to the study of the non-local problem for a linear elliptic equation with a parameter, whose principal part is in a divergence form. This allows us to remove some restrictions on the coefficient  $\beta$  appearing in the boundary condition (2.2). On the other hand this also suggests further extensions of the solvability of the problem (2.1), (2.2) in a weighted Sobolev space  $\tilde{W}^{1,2}(Q)$ . We adopt here the  $L^2$ -approach to the Dirichlet problem with  $L^2$ -boundary data from [13] and [10].

## 2. UNIQUENESS AND A PRIORI ESTIMATE.

We consider a linear equation of the elliptic type

$$\Delta u = \sum_{i,j=1}^n a_{ij}(x) D_{ij} u + \sum_{i=1}^n b_i(x) D_i u + c(x) u = f(x) \quad (2.1)$$

in  $Q$ , where  $Q$  is a bounded domain in  $R_n$ . The purpose of this paper is to investigate the following non-local problem: given continuous functions  $h$  and  $\beta$  defined on the boundary  $\partial Q$  of  $Q$  find a solution  $u \in C^2(Q) \cap C(\bar{Q})$  satisfying the boundary condition

$$u(x) - \beta(x) u(\phi(x)) = h(x) \text{ on } \partial Q, \quad (2.2)$$

where  $\phi$  is a given continuous mapping of  $\partial Q$  into  $Q$ .

Throughout this section we make the following assumption

(A) The coefficients of the operator  $L$  are bounded in  $Q$  and there exists a constant  $\gamma > 0$  such that

$$\gamma |\xi|^2 \leq \sum_{i,j=1}^n a_{ij}(x) \xi_i \xi_j$$

for all  $x \in Q$  and  $\xi \in R_n$ .

Moreover we assume that  $Q$  satisfies an interior sphere condition at each point of  $\partial Q$  (see [11], p. 33-35).

The uniqueness of the problem (2.1), (2.2) is a consequence of the strong maximum principle.

PROPOSITION 1. Let  $|\beta(x)| \leq 1$  on  $\partial Q$  and  $c(x) \leq 0$  in  $Q$  and suppose that

either

$$(a) \quad -1 \leq \beta(x_0) < 1 \text{ at some point } x_0 \in \partial Q$$

or

$$(b) \quad c(x_1) < 0 \text{ at some point } x_1 \in Q.$$

Then the problem (2.1), (2.2) has at most one solution in  $C^2(Q) \cap C(\bar{Q})$ .

PROOF. It is sufficient to show that if  $f(x) \equiv 0$  on  $Q$  and  $h(x) \equiv 0$  on  $\partial Q$  then  $u \equiv 0$  is the only solution of the problem (2.1), (2.2). It is clear that under each of the assumptions (a) or (b) any constant solution must be identically equal to 0. If  $u \not\equiv 0$  then  $u$  must be a non-constant solution and by the strong maximum principle ([11], Theorem 3.5) we may assume that

$$u(x_2) = \max_{\bar{Q}} u(x) > 0 \text{ with } x_2 \in \partial Q.$$

If  $\beta(x_2) = 0$  we get a contradiction. Therefore it remains to consider two cases

$$(i) \quad 0 < \beta(x_2) \leq 1 \quad \text{and} \quad (ii) \quad -1 \leq \beta(x_2) < 0.$$

In the first case  $u(\phi(x_2)) = \frac{u(x_2)}{\beta(x_2)} \geq u(x_2)$ , which is impossible since  $\phi(x_2) \in Q$ . In the second case (ii) we have

$$u(\phi(x_2)) = \frac{u(x_2)}{\beta(x_2)} < 0$$

and by the strong maximum principle  $u$  takes on a negative minimum at  $x_3 \in \partial Q$ , that is

$$u(x_3) = \min_{\bar{Q}} u(x) < 0$$

and we may assume that  $\beta(x_3) < 0$  since otherwise we get a contradiction. Hence

$$u(\phi(x_3)) = \frac{u(x_3)}{\beta(x_3)} > 0.$$

Now we distinguish two cases either

$$u(x_2) \leq |u(x_3)| \quad \text{or} \quad u(x_2) > |u(x_3)|.$$

We show that both cases lead to a contradiction. Indeed, in the first case we have

$$u(x_2) \leq |u(x_3)| = |\beta(x_3)| |u(\phi(x_3))| \leq |u(\phi(x_3))| = u(\phi(x_3)),$$

which is impossible. In the second case we have

$$|u(x_3)| + u(x_2) = \beta(x_2) u(\phi(x_2)) \leq |u(\phi(x_2))|.$$

Since both values  $u(x_3)$  and  $u(\phi(x_2))$  are negative  $u$  attains its negative minimum at  $\phi(x_2) \in Q$  and we arrive at a contradiction.

Inspection of the proof of Proposition 1 shows that the following version of the maximum principle holds true.

PROPOSITION 2. Suppose that  $c(x) \leq 0$  in  $Q$  and  $0 \leq \beta(x) \leq 1$  on  $\partial Q$ . Let  $Lu \leq 0$  ( $\geq 0$ ) in  $Q$ , and  $u(x) - \beta(x) u(\phi(x)) \geq 0$  ( $\leq 0$ ) on  $\partial Q$ . Then  $u(x) \geq 0$  ( $\leq 0$ ) on  $\bar{Q}$ .

As an immediate consequence we deduce an a priori estimate

THEOREM 1. Suppose that  $c(x) \leq -d$  in  $Q$  and  $0 \leq \beta(x) \leq \alpha$  where  $d > 0$  and  $0 < \alpha < 1$  are constants. If  $u$  is a solution of the problem (2.1), (2.2), then

$$|u(x)| \leq \frac{1}{d} \sup_Q |f(x)| + \frac{1}{1-\alpha} \sup_{\partial Q} |h(x)| \text{ for all } x \in \bar{Q}.$$

PROOF. Let us define

$$v(x) = u(x) - \frac{1}{d} \sup_Q |f(x)| - \frac{1}{1-\alpha} \sup_{\partial Q} |h(x)|,$$

then we have

$$Lv = f - \frac{c}{d} \sup_Q |f(x)| - \frac{c}{1-\alpha} \sup_{\partial Q} |h(x)| \geq f + \sup_Q |f(x)| \geq 0$$

in  $Q$  and

$$\begin{aligned} v(x) - \beta(x) v(\phi(x)) &= h(x) - \frac{1}{d} \sup_Q |f(x)| - \frac{1}{1-\alpha} \sup_{\partial Q} |h(x)| + \\ &+ \frac{\beta(x)}{1-\alpha} \sup_{\partial Q} |h(x)| + \frac{\beta(x)}{\alpha} \sup_Q |f(x)| \leq \\ &\leq (1 - \frac{1}{1-\alpha} + \frac{\alpha}{1-\alpha}) \sup_{\partial Q} |h(x)| = 0 \end{aligned}$$

on  $\partial Q$ . Hence by Proposition 2

$$u(x) \leq \frac{1}{d} \sup_Q |f(x)| + \frac{1}{1-\alpha} \sup_{\partial Q} |h(x)|$$

on  $Q$ . Similarly we can establish the inequality

$$u(x) \geq -\frac{1}{d} \sup_Q |f(x)| - \frac{1}{1-\alpha} \sup_{\partial Q} |h(x)|$$

on  $Q$ , considering the auxiliary function

$$w(x) = u(x) + \frac{1}{d} \sup_Q |f(x)| + \frac{1}{1-\alpha} \sup_{\partial Q} |h(x)|.$$

REMARK 1. If  $c \equiv 0$  on  $Q$  and  $\beta(x) = 1$  on  $\partial Q$ , then any two solutions of the problem (2.1), (2.2) differ by a constant.

### 3. EXISTENCE OF CLASSICAL SOLUTIONS.

We commence by considering a particular case of the non-local problem (2.1), (2.2) which consists of finding a harmonic function  $u$  on  $B(0,1)$  and satisfying the boundary condition

$$u(z) - \beta u(\phi(z)) = h(z) \quad \text{on } \partial B(0,1), \quad (3.1)$$

where  $B(0,1)$  is an open disc in  $R_2$  of radius 1 centred at 0,  $\beta$  is a constant in the interval  $[-1,1]$  and  $h$  is a continuous function on  $\partial B(0,1)$ . The mapping  $\phi$  is given by  $\phi(z) = \phi_*(\delta z)$  with  $0 < \delta < 1$ , where  $\phi_*$  is a univalent analytic function on  $B(0,1)$  such that  $|\phi_*(z)| \leq 1$  in  $B(0,1)$  and  $\phi_*(0) = 0$ . By virtue of Schwarz's lemma we have

$$|\phi_*(z)| < |z| \quad \text{for all } z \in B(0,1). \quad (3.2)$$

The function  $\phi$  maps disc  $B(0,1)$  conformally and univalently onto certain set contained in  $B(0,1)$ . Letting  $\phi_0(z) = z$  and  $\phi_k(z) = \phi(\phi_{k-1}(z))$  for  $k = 1, 2, \dots$  we have

$$|\phi_k(z)| < \delta^k |z| \quad \text{in } B(0,1), \quad (3.3)$$

for  $k = 1, 2, \dots$ . Since  $u$  and  $u(\phi(z))$  are harmonic functions we have the following representation formula

$$u(z) - \beta u(\phi(z)) = \operatorname{Re} \frac{1}{\pi i} \int_{\partial B(0,1)} \left[ \frac{1}{t-z} - \frac{1}{2t} \right] h(t) dt \equiv F(z) \quad (3.4)$$

Suppose first that  $-1 < \beta < 1$ . Iterating (3.4) we get

$$u(z) = \beta^n u(\phi_n(z)) + \sum_{k=1}^n \beta^{k-1} F(\phi_{k-1}(z)). \quad (3.5)$$

It follows from (3.3) that

$$\begin{aligned} |F(\phi_n(z))| &= \left| \operatorname{Re} \frac{1}{2\pi i} \int_{\partial B(0,1)} \frac{t + \phi_n(z)}{t - \phi_n(z)} h(t) dt \right| \leq \\ &< \frac{1+\delta}{2\pi(1-\delta)} \int_0^\pi |h(e^{is})| ds \end{aligned}$$

for  $n = 1, 2, \dots$  and consequently letting  $n \rightarrow \infty$  in (3.5) we obtain

$$u(z) = \sum_{n=1}^{\infty} \rho^{n-1} F(\phi_{n-1}(z))$$

uniformly on  $\overline{B(0,1)}$ .

Let us now consider the case  $\rho = -1$ . It follows from (3.4) that

$$2u(0) = \operatorname{Re} \frac{1}{\pi i} \int_{\partial B(0,1)} \frac{h(t)}{2t} dt \quad (3.6)$$

and the functional equation (3.4) can be written in the form

$$\begin{aligned} u(z) + u(\phi(z)) - 2u(0) &= \operatorname{Re} \frac{1}{\pi i} \int_{\partial B(0,1)} \left[ \frac{1}{t-z} - \frac{1}{2t} - \frac{1}{2\bar{t}} \right] h(t) dt \\ &= \operatorname{Re} \frac{1}{\pi i} \int_{\partial B(0,1)} \frac{h(t)}{t(t-z)} dt \cdot z \equiv \operatorname{Re}[\tilde{\Phi}(z) \cdot z]. \end{aligned}$$

Iterating the last equation we obtain

$$u(z) = u(\phi_{2n}(z)) + \sum_{j=0}^{2n-1} (-1)^j \operatorname{Re}[\tilde{\Phi}(\phi_j(z)) \phi_j(z)] \quad (3.7)$$

and

$$u(z) = 2u(0) - u(\phi_{2n-1}(z)) + \sum_{j=0}^{2n-2} (-1)^j \operatorname{Re}[\tilde{\Phi}(\phi_j(z)) \phi_j(z)]. \quad (3.8)$$

It is easy to see that

$$|\tilde{\Phi}(\phi_n(z))| \leq \frac{\delta}{\pi(1-\delta)} \int_0^{2\pi} |h(e^{is})| ds$$

for all  $n = 1, 2, \dots$  and  $z \in \overline{B(0,1)}$ . Since  $|\phi_j(z)| < \delta^j$  in  $\overline{B(0,1)}$  the series  $\sum_{j=0}^{\infty} (-1)^j \operatorname{Re}[\tilde{\Phi}(\phi_j(z)) \phi_j(z)]$  converges uniformly on  $\overline{B(0,1)}$ . Letting  $n \rightarrow \infty$  in (3.7) and (3.8) we obtain the same limit in both cases

$$u(z) = u(0) + \sum_{j=0}^{\infty} (-1)^j \operatorname{Re}[\tilde{\Phi}(\phi_j(z)) \phi_j(z)]$$

and invoking (3.6) we get that

$$u(z) = \operatorname{Re} \frac{1}{4\pi i} \int_{\partial B(0,1)} \frac{h(t)}{t} dt + \sum_{j=0}^{\infty} (-1)^j \operatorname{Re}[\tilde{\Phi}(\phi_j(z)) \phi_j(z)].$$

Finally let  $\rho = 1$ , then

$$u(z) - u(\phi(z)) = \operatorname{Re} \frac{1}{\pi i} \int_{\partial B(0,1)} \left[ \frac{1}{t-z} - \frac{1}{2t} \right] h(t) dt.$$

By Remark 1 any two solutions differ by a constant. If  $z = 0$ , then

$$0 = u(0) - u(\phi(0)) = \operatorname{Re} \frac{1}{\pi i} \int_{\partial B(0,1)} \frac{f(t)}{t} dt,$$

which gives a necessary and sufficient condition for the solvability of the problem (3.1). Hence

$$u(z) - u(\phi(z)) = \operatorname{Re}[\bar{\Phi}(z)z].$$

Iterating this functional equation we obtain

$$u(z) = u(\phi_n(z)) + \sum_{j=1}^{n-1} \operatorname{Re}[\bar{\Phi}(\phi_j(z)) \phi_j(z)].$$

Letting  $n \rightarrow \infty$  we obtain

$$u(z) = u(0) + \sum_{j=1}^{\infty} \operatorname{Re}[\bar{\Phi}(\phi_j(z)) \phi_j(z)].$$

To determine a solution in a unique way we may impose an additional condition  $u(0) = C$ , where  $C > 0$  is a given constant. The case  $\rho = 1$  was considered by Bitsadze in [1] and [2] under the assumption that  $h$  is "Holder continuous on  $\partial Q$ ."

In a general case we reduce the problem (2.1), (2.2) to the Fredholm integral equation of the second kind.

**THEOREM 2.** Suppose that the assumptions of Proposition 1 hold and that

$D_{ij}^2 a_{ij}$  ( $i, j = 1, \dots, n$ ),  $D_i b_i$  ( $i = 1, \dots, n$ ) and  $c$  are "Holder continuous on  $\bar{Q}$ ". Then the problem (2.1), (2.2) admits a unique solution  $u$  in  $C^2(Q) \cap C(\bar{Q})$ .

**PROOF.** We try to find a solution in the form

$$u(x) = \int_{\partial Q} \frac{dG(x, y)}{dn_y} v(y) dS_y - \int_Q G(x, y) f(y) dy, \quad (3.9)$$

where  $v \in C(\partial Q)$  is to be determined,  $G$  is the Green function for the operator  $L$  and  $\frac{dG}{dn_y}$  denotes the conormal derivative. The boundary condition (2.2) leads to the Fredholm integral equation of the second kind.

$$\begin{aligned} v(x) - \int_{\partial Q} \beta(x) \frac{dG(\phi(x), y)}{dn_y} v(y) dS_y &= h(x) + \\ &+ \int_Q \beta(x) G(\phi(x), y) f(y) dy. \end{aligned} \quad (3.10)$$

Since  $\phi(\partial Q) \subset Q$  the kernel  $\beta(x) \frac{dG(\phi(x), y)}{dn_y}$  is continuous function on  $\partial Q \times \partial Q$ . By Proposition 1 the homogeneous equation corresponding to (3.10) has only trivial solution. Hence by the Fredholm alternative there exists

a unique solution  $v \in L^2(\partial Q)$  which by the continuity of the kernel belongs to  $C(\partial Q)$ . Consequently the formula (3.9) gives a solution to the problem (2.1), (2.2).

#### 4. ENERGY ESTIMATE.

In this section we consider the elliptic equation in the form

$$\begin{aligned} \Delta u + \lambda u = - \sum_{i,j=1}^n D_i(a_{ij}(x)) D_j u + \sum_{i=1}^n b_i(x) D_i u + c(x)u + \\ + \lambda u = f(x) \text{ in } Q, \end{aligned} \quad (4.1)$$

with the boundary condition (2.2).

Throughout this section we assume that  $\beta$  is a continuous function on  $\partial Q$  and  $\phi: \partial Q \rightarrow Q$  is a  $C^1$ -mapping with the positive Jacobian. Further we assume that  $a_{ij}$ ,  $D_i a_{ij}$ ,  $b_i$  ( $i, j = 1, \dots, n$ )  $c$  and  $f$  are "Holder continuous on  $Q$  and that  $Q$  is a bounded domain with the boundary of class  $C^2$ .

The objective of this section is to show that the problem (4.1), (2.2) has a unique solution for large values of the parameter  $\lambda$ .

For small  $\delta > 0$  we define  $Q_\delta = Q \cap \{x; \min_{y \in \partial Q} |x-y| > \delta\}$ .

According to Lemma 14.16 in [11] (p. 355), the distance  $r(x) = \text{dist}(x, \partial Q)$  belongs to  $C^2(\bar{Q} - Q_\delta)$  if  $\delta_0$  is sufficiently small. Denote by  $\rho(x)$  the extension of the function  $r(x)$  into  $\bar{Q}$  satisfying the following properties

$\rho(x) = r(x)$  for  $x \in \bar{Q} - Q_\delta$ ,  $\rho \in C^2(\bar{Q})$ ,  $\rho(x) > \frac{3\delta_0}{4}$  in  $Q_\delta$ ,  $\gamma_1^{-1} r(x) \leq \rho(x) \leq \gamma_1 r(x)$  in  $Q$  for some constant  $\gamma_1 > 0$ ,  $\partial Q_\delta = \{x; \rho(x) = \delta\}$  for  $\delta \in (0, \delta_0)$  and finally  $\partial Q = \{x; \rho(x) = 0\}$ .

**THEOREM 3.** There exist positive constants  $\lambda_0$ ,  $C$  and  $d$  such that if  $u$  is a solution in  $C^2(Q) \cap C(\bar{Q})$  of the problem (4.1), (2.2) for  $\lambda \geq \lambda_0$  then

$$\begin{aligned} \int_Q |Du(x)|^2 r(x) dx + \int_Q u(x)^2 r(x) dx + \sup_{0 < \delta \leq d} \int_{\partial Q_\delta} u(x)^2 dS_x \leq \\ \leq C \left( \int_{\partial Q} h(x)^2 dS_x + \int_Q f(x)^2 dx \right). \end{aligned}$$

**Proof.** We follow the proof of Theorem 5 in [13]. Multiplying (4.1) by



$$v(x) = \begin{cases} u(x) (\rho(x) - \delta) & \text{on } Q_\delta \\ 0 & \text{on } Q - Q_\delta \end{cases}$$

and integrating by parts we obtain

$$\begin{aligned} \frac{1}{2} \int_{\partial Q_\delta} \sum_{i,j=1}^n a_{ij}(x) D_i \rho D_j \rho u^2 dS_x &= -\frac{1}{2} \int_{Q_\delta} \sum_{i,j=1}^n D_i (a_{ij}(x) D_j \rho) u^2 dx + \\ &+ \int_{Q_\delta} \sum_{i=1}^n b_i(\lambda) D_i u \cdot u(\rho - \delta) dx + \int_{Q_\delta} (c(x) + \lambda) u^2 (\rho - \delta) dx - \int_{Q_\delta} f u (\rho - \delta) dx. \end{aligned}$$

Applying Hölder's inequality we easily obtain

$$\begin{aligned} \sup_{0 < \delta \leq d} \int_{\partial Q_\delta} u^2 dS_x &\leq C_1 \left[ \int_Q |Du|^2 \rho dx + \int_Q u^2 dx + \right. \\ &\left. + \lambda \int_Q u^2 \rho dx + \int_Q f^2 dx \right], \end{aligned} \quad (4.2)$$

where  $d$  and  $C_1$  are positive constants.

Similarly

$$\int_Q |Du|^2 \rho dx + \lambda \int_Q u^2 \rho dx \leq C_2 \left[ \int_Q u^2 dx + \int_Q f^2 dx + \int_{\partial Q} u^2 dS_x \right] \quad (4.3)$$

for some  $C_2 > 0$ . It follows from (2.2) that

$$\begin{aligned} \int_Q |Du|^2 \rho dx + \lambda \int_Q u^2 \rho dx &\leq C_3 \left[ \int_Q u^2 dx + \int_Q f^2 dx + \right. \\ &\left. + \int_{\partial Q} h^2 dS_x + \int_{\partial Q} u(\phi(x))^2 dS_x \right], \end{aligned} \quad (4.4)$$

where  $C_3 > 0$ . The estimates (4.2) and (4.4) yield that

$$\begin{aligned} \int_Q |Du|^2 \rho dx + \lambda \int_Q u^2 \rho dx + \sup_{0 < \delta < d} \int_{\partial Q_\delta} u^2 dS_x &\leq \\ &\leq C_4 \left[ \int_Q u^2 dS_x + \int_Q f^2 dx + \int_{\partial Q} h^2 dS_x + \int_{\partial Q} u(\phi(x))^2 dS_x \right], \end{aligned} \quad (4.5)$$

for some  $C_4 > 0$ . Since  $\phi$  is a  $C^1$ -mapping with the positive Jacobian it is obvious that

$$\int_{\partial Q} u(\phi(x))^2 dS_x \leq C_5 \int_{\phi(\partial Q)} u^2 dS_x \leq C_6 \left[ \int_{Q_\phi} |Du|^2 dx + \int_{Q_\phi} u^2 dx \right], \quad (4.6)$$

where  $C_5 > 0$  and  $C_6 > 0$  are constants and  $Q_\phi$  is a domain containing  $\phi(\partial Q)$  with  $\text{dist}(Q_\phi, \partial Q) > 0$ . Consequently by virtue of the Caccioppoli inequality

we have

$$\int_Q u(\phi(x))^2 dS_x \leq C_7 \left[ \int_Q u^2 dx + \int_Q f^2 dx \right] \quad (4.7)$$

for some  $C_7 > 0$ . We now observe that

$$\int_Q u^2 dx \leq \frac{1}{d_1} \int_Q u^2 \rho dx + d \sup_{0 < \delta \leq d} \int_{\partial Q} u^2 dS_x, \quad (4.8)$$

where  $d_1 = \inf_{\bar{Q}_d} \rho(x)$ . Choosing  $d$  sufficiently small and  $\lambda$  sufficiently large we easily derive the desired estimate from (4.5), (4.6), (4.7) and (4.8).

Repeating the argument of Theorem 2 we deduce the following

**THEOREM 4.** There exists a positive constant  $\lambda_0$  such that for every  $\lambda \geq \lambda_0$  the problem (4.1), (2.2) admits a unique solution in  $C^2(Q) \cap C(\bar{Q})$ .

## 5. WEAK SOLUTIONS.

The energy estimate from Section 4 shows that one can expect solutions of the problem (4.1), (2.2) in a weighted Sobolev space defined by

$$W^{1,2}_r(Q) = \{u \in W^{1,2}_{loc}(Q); \int_Q |Du(x)|^2 r(x) dx + \int_Q u(x)^2 dx < \infty\}$$

and equipped with the norm

$$\|u\|_{W^{1,2}_r}^2 = \int_Q [|Du(x)|^2 r(x) + u(x)^2] dx.$$

We recall briefly that a function  $u$  is said to be a weak (generalized) solution of (4.1) if  $u \in W^{1,2}_{loc}(Q)$  and it satisfies

$$\int_Q \left[ \sum_{i,j=1}^n a_{ij} D_i u D_j v + \sum_{i=1}^n b_i D_i u \cdot v + (c+\lambda)uv \right] dx = \int_Q f v dx \quad (5.1)$$

for each  $v \in W^{1,2}_r(Q)$  with compact support in  $Q$ .

To proceed further we need some terminology. It follows from the regularity of the boundary  $\partial Q$  that there exists a number  $\delta_0$  such that for  $\delta \in (0, \delta_0)$  the domain  $Q_\delta$  (defined in Section 4) with the boundary  $\partial Q_\delta$  possesses the following property: to each  $x_0 \in \partial Q$  there exists a unique point  $x_\delta(x_0) \in \partial Q_\delta$  such that  $x_\delta(x_0) = x_0 - \delta \nu(x_0)$ , where  $\nu(x_0)$  is the outward normal to  $\partial Q$  at  $x_0$ . The above relation gives a one-to-one mapping of class  $C^1$ , of  $\partial Q$  onto  $\partial Q_\delta$ .

It is known that elements of the space  $\tilde{W}^{1,2}(Q)$ , in general, do not have traces on the boundary  $\partial Q$  (see [12]). However, by Theorem 4 in [13], if  $u \in \tilde{W}^{1,2}(Q)$  is a solution of (4.1) then there exists a function  $\zeta \in L^2(\partial Q)$  such that

$$\lim_{\delta \rightarrow 0} \int_{\partial Q} [u(x_\delta) - \zeta(x)]^2 dS_x = 0.$$

Therefore, as in the paper [13], we adopt the following  $L^2$ -approach to the problem (4.1), (2.2).

Let  $h \in L^2(\partial Q)$ . A weak solution  $u \in W_{loc}^{1,2}(Q)$  of (4.1) is a solution of the non-local problem with the boundary condition (2.2) if

$$\lim_{\delta \rightarrow 0} \int_{\partial Q} [u(x_\delta) - \beta(x) u(\phi(x)) - h(x)]^2 dS_x = 0. \quad (5.2)$$

It follows from Theorem 1 in [13] that if  $u \in W_{loc}^{1,2}(Q)$  is a solution of the problem (4.1), (2.2) (with the boundary condition (2.2) understood in the sense of (5.2) then  $u \in \tilde{W}^{1,2}(Q)$ . We mention also that  $u(\phi(x))$  is understood in the sense of trace, which is well defined since  $u \in W_{loc}^{1,2}(Q)$  (see [14], chap. 6).

We now are in a position to establish the existence result in  $\tilde{W}^{1,2}(Q)$  of the problem (4.1), (2.2).

**THEOREM 5.** Let  $h \in L^2(\partial Q)$ . Then there exists a positive constant  $\lambda_0$  such that for  $\lambda \geq \lambda_0$  the problem (4.1), (2.2) in  $W_{loc}^{1,2}(Q)$  admits a unique solution.

**PROOF.** Let  $h_m$  be a sequence in  $C^1(\partial Q)$  such that  $\lim_{m \rightarrow \infty} \int_{\partial Q} (h_m - h)^2 dS_x = 0$ . Let  $\lambda_0$  be a constant from Theorem 4 and assume that  $\lambda \geq \lambda_0$ . For each  $m \geq 1$  Theorem 4 guarantees the existence of the unique solution

$u_m \in C^2(Q) \cap C(\bar{Q})$  of the problem (4.1), (2.2) with  $h = h_m$ . Moreover we have for each  $u_m$

$$\int_Q |Du_m|^2 dx + \int_Q u_m^2 dx \leq C \left( \int_Q f^2 dx + \int_{\partial Q} h_m^2 dS_x \right),$$

where  $C > 0$  is a constant independent of  $m$ . Since the sequence  $u_m$  is bounded in  $\tilde{W}^{1,2}(Q)$ , there exists a subsequence, which we relabel as  $u_m$ , converging weakly in  $\tilde{W}^{1,2}(Q)$  to a function  $u$ . By Theorem 4.11 in [15],  $\tilde{W}^{1,2}(Q)$  is compactly embedded in  $L^2(Q)$  and therefore we may assume that  $u_m$

converges to  $u$  in  $L^2(Q)$ . It is obvious that  $u$  is a weak solution of (4.1). By Theorem 4 in [13]  $u$  has a trace  $\zeta \in L^2(\partial Q)$  in the sense of  $L^2$ -convergence, that is

$$\lim_{\delta \rightarrow 0} \int_{\partial Q} [u(x_\delta) - \zeta(x)]^2 dS_x = 0.$$

To complete the proof we show that  $\zeta(x) = \beta(x) u(\phi(x)) + h(x)$  a.e. on  $\partial Q$ .

Let  $\Phi \in C^1(Q)$ . It is easy to show  $\Phi(x) \rho(x)$  is a legitimate test function in (5.1) and integrating by parts we obtain

$$\begin{aligned} \int_{\partial Q} \zeta \Phi \sum_{i,j=1}^n a_{ij} D_i \rho D_j \rho dS_x &= - \int_Q \sum_{i,j=1}^n D_i (a_{ij} D_j \rho \Phi) u dx + \\ &+ \int_Q \sum_{i=1}^n b_i D_i u \Phi dx + \int_Q (c+\lambda) u \Phi \rho dx - \int_Q f \Phi u \rho dx \equiv \\ &\equiv F(u). \end{aligned} \quad (5.3)$$

Similarly

$$\int_{\partial Q} [h_m(x) + \beta(x) u_m(\phi(x))] \sum_{i,j=1}^n a_{ij} D_i \rho D_j \rho dS_x \equiv F(u_m). \quad (5.4)$$

Applying the estimate (4.6) and the obvious analogue of the energy estimate to  $u_p - u_q$  we obtain

$$\begin{aligned} \int_{\partial Q} [u_p(\phi(x)) - u_q(\phi(x))]^2 dS_x &\leq C \left[ \int_Q |Du_p - Du_q|^2 dx + \right. \\ &+ \left. \int_Q |u_p - u_q|^2 dx \right] \leq \frac{C}{k} \int_Q |Du_p - Du_q|^2 \rho dx + C \int_Q |u_p - u_q|^2 dx \leq \\ &\leq \tilde{C} \int_{\partial Q} |h_p - h_q|^2 dS_x, \end{aligned}$$

where  $k = \inf_{\partial Q} \rho(x)$  and  $\tilde{C} > 0$  is a constant independent of  $p$  and  $q$ . Hence by the continuity of weak solutions on  $Q$  we may assume that  $u_m(\phi(x)) \rightarrow u(\phi(x))$  as  $m \rightarrow \infty$  in  $L^2(\partial Q)$ . Combining this with the fact that  $F(u_m) \rightarrow F(u)$  as  $m \rightarrow \infty$  we deduce from (5.3) and (5.4) that  $\zeta(x) = \beta(x) u(\phi(x)) + h(x)$  a.e. on  $\partial Q$  and this completes the proof.

REMARK 2. It is worth noting that the non-local problem of the type (4.1), (2.2) has been studied in [9] for the higher order elliptic equations. The corresponding boundary datum  $h$  in [9] belongs to the space  $H^{\frac{1}{2}}(\partial Q)$ . Since  $H^{\frac{1}{2}}(\partial Q)$  is a proper subspace  $L^2(\partial Q)$ , Theorem 5 cannot be deduced from the results of the paper [9].

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