

NUMERICAL SOLUTION OF INTEGRAL EQUATIONS WITH FINITE PART INTEGRALS

SAMIR A. ASHOUR

(Received 22 July 1994 and in revised form 29 March 1994)

ABSTRACT. We obtain convergence rates for several algorithms that solve a class of Hadamard singular integral equations using the general theory of approximations for unbounded operators.

Keywords and phrases. Singular integral equations.

1991 Mathematics Subject Classification. 65D05, 45L05.

1. Introduction. In several physical problems in aerodynamics, hydrodynamics, and elasticity, one encounters integral equations of the form

$$Ax = \frac{1}{\pi} \int_{-1}^1 \frac{\sqrt{1-\tau^2} x(\tau)}{(\tau-t)^2} d\tau + \frac{1}{\pi} \int_{-1}^1 \sqrt{1-\tau^2} h(t, \tau) x(\tau) d\tau = y, \quad (1.1)$$

where the first integral in (1.1) is a finite part integral [4]. Under suitable conditions on the kernel and the right-hand side, the convergence of Galerkin's method and several collocation methods, proposed by Ioakimidis [5] and Williams [9], has been discussed by Golberg [2, 3]. This author, also, used a classical Fredholm theory to establish the existence of a solution and likewise the basic tools necessary to discuss convergence. In this paper, we discuss the convergence of the mechanical quadratures method for solving (1.1) and the convergence of the least squares method for solving the Hadamard singular integral equation

$$Kx = \frac{1}{\pi} \int_{-1}^1 \frac{\sqrt{1-\tau^2} x(\tau)}{(\tau-t)^2} d\tau + T(x, t) = y, \quad (1.2)$$

where T is the given continuous operator.

2. Least squares method. Let $X = L_{2,\rho}$ denote the space of square integrable functions with respect to $\rho = \sqrt{1-t^2}$. The inner product on $L_{2,\rho}$ is given by

$$(\phi, \psi)_\rho = \frac{2}{\pi} \int_{-1}^1 \rho(t) \phi(t) \psi(t) dt \quad \text{and} \quad \|\phi\|_\rho = \sqrt{(\phi, \phi)_\rho}. \quad (2.1)$$

Let

$$U_m(t) = \frac{\sin[(m+1) \arccos t]}{\sqrt{1-t^2}}, \quad m = 0, 1, 2, \dots \quad (2.2)$$

denote the Chebyshev polynomials of the second kind. The solution x is, now, approximated by

$$x_n(t) = \sum_{k=1}^n \alpha_k U_{k-1}(t), \quad -1 \leq t \leq 1. \quad (2.3)$$

According to this method, we obtain a system of n linear algebraic equations in n unknowns

$$\sum_{i=1}^n \alpha_i (KU_{i-1}, KU_{j-1})_\rho = (y, KU_{j-1})_\rho, \quad 1 \leq j \leq n. \quad (2.4)$$

It is easy to prove that (2.4) is equivalent to

$$\begin{aligned} \sum_{k=1}^n \alpha_k \{ (TU_{k-1}, TU_{j-1})_\rho - j(TU_{k-1}, U_{j-1})_\rho - k(TU_{j-1}, U_{k-1})_\rho \} + j^2 \alpha_j \\ = (y, KU_{j-1})_\rho, \quad 1 \leq j \leq n. \end{aligned} \quad (2.5)$$

THEOREM 2.1. *If the following conditions hold*

- (i) $y \in L_{2,\rho}$, T is a continuous operator in $L_{2,\rho}$;
 - (ii) $\ker K = \{0\}$;
 - (iii) equation (1.2) has a solution $x^* \in L_{2,\rho}$ for a given $y \in L_{2,\rho}$, then for all $n \in \mathbb{N}$ equation (2.5) has a unique solution $\{\alpha_k^*\}_1^n$; and if
 - (iv) $\{KU_{i-1}\}$ is closed in $L_{2,\rho}$,
- then

$$\|r_n\|_\rho = \|y - Kx_n^*\|_\rho \rightarrow 0 \quad \text{as } n \rightarrow \infty, \quad x_n^* = \sum_{k=1}^n \alpha_k^* U_{k-1}(t). \quad (2.6)$$

PROOF. Since $\{U_{k-1}\}$ are linearly independent, then, from (ii), it follows that $\{KU_{k-1}\}$ are, also, linearly independent. Therefore, the system of equations (2.5) is non-singular and so, it has a unique solution for all n . Also, for $\beta_k \in \mathbb{R}$, we get

$$\|r_n\|_\rho = \|y - Kx_n^*\|_\rho \leq \left\| y - \sum_{i=1}^n \beta_i KU_{i-1} \right\|_\rho. \quad (2.7)$$

If condition (iv) is satisfied, then $\|r_n\| \rightarrow 0$ as $n \rightarrow \infty$. □

Now, we replace condition (ii) by the following condition:

- (ii)' $K : L_{2,\rho} \rightarrow L_{2,\rho}$ has a left bounded inverse operator K_l^{-1} .

THEOREM 2.2. *Assume that (i), (ii)', (iii), and (iv) are satisfied, then $\|x_n^* - x^*\|_\rho = O(\|r_n\|_\rho) \rightarrow 0$ as $n \rightarrow \infty$.*

PROOF. From (ii)' and (iii), we have $x^* - x_n^* = K_l^{-1}K(x^* - x_n^*) = K_l^{-1}(y - Kx_n^*)$, then $\|x_n^* - x^*\|_\rho \rightarrow 0$ as $n \rightarrow \infty$. □

3. Mechanical quadratures methods. We introduce the following method for solving (1.1): Consider the approximation x_n of x given by (2.3). Due to this method, we get the following

$$\sum_{i=1}^n \alpha_i \left\{ -i U_{i-1}(t_j) + \frac{1}{n+1} \sum_{r=1}^n (1-t_r^2) h(t_j, t_r) U_{i-1}(t_r) \right\} = y(t_j), \quad 1 \leq j \leq n, \quad (3.1)$$

where $t_j = \cos(j\pi/n+1)$.

THEOREM 3.1. *If the following conditions*

(i) $y \in C[-1, 1]$, $h \in C[-1, 1; -1, 1]$

(ii) *equation (1.1) has a unique solution $x^* \in X$ for all $y \in X$*

are satisfied, then, for n sufficiently large, equation (3.1) has a unique solution α_k^ ,*

$$\|x^* - x_n^*\|_\rho = O\{E_n(y)_C + E_n^t(h)_C + E_n^\tau(h)_C\}, \quad (3.2)$$

where $E_n^t(f)_C = \inf \{\|f - P_n\|_C, P_n \text{ is a polynomial of degree } \leq n\}$.

PROOF. Define the projection operator $P_n^t : C \rightarrow X$ by

$$P_n^t(x) = \sum_{k=1}^n x(t_k) \frac{U_n(t)}{(t-t_k)U_n'(t_k)}. \quad (3.3)$$

Thus, equation (3.1) can be written in the equivalent form

$$A_n x_n = G x_n + P_n^t H P_n^\tau(h x_n) = P_n^t y, \quad x_n, P_n y \in X_n, \quad (3.4)$$

where $X_n = \{x_n : x_n = \sum_{k=1}^n \alpha_k U_{k-1}(t), t \in [-1, 1]\}$,

$$\begin{aligned} Gx &= \frac{1}{\pi} \int_{-1}^1 \frac{\sqrt{1-\tau^2} x(\tau)}{(\tau-t)^2} d\tau, \quad x \in X, \\ Hhx &= \frac{1}{\pi} \int_{-1}^1 \sqrt{1-\tau^2} h(t, \tau) x(\tau) d\tau. \end{aligned} \quad (3.5)$$

Since $P_n^t H P_n^\tau(h x_n) = P_n^t H((P_n^\tau h) x_n)$, then, for all $x_n \in X_n$, we get

$$\begin{aligned} \|Ax_n - A_n x_n\|_\rho &= \|Hhx_n - P_n^t H P_n^\tau(h x_n)\|_\rho \\ &\leq \|(Hh - P_n^t Hh)x_n\|_\rho + \|P_n^t H(h - P_n^\tau h)x_n\|_\rho. \end{aligned} \quad (3.6)$$

According to [8], $\|P_n\| \leq C_1$, $\|x - P_n x\|_\rho \leq C_2 E_n(x)_C$, $x \in C[-1, 1]$, we get

$$\begin{aligned} \|Ax_n - A_n x_n\|_\rho &\leq C_2 E_n(Hhx_n) + C_1 \|H(h - P_n^\tau h)x_n\|_C \\ &= O\{E_n^t(h)_C + E_n^\tau(h)_C\} \|x_n\|_\rho, \end{aligned} \quad (3.7)$$

so that

$$\|A - A_n\|_{X_n \rightarrow X} = O\{E_n^t(h)_C + E_n^\tau(h)_C\} = \epsilon_n. \quad (3.8)$$

According to [1], for all n such that $\|A^{-1}\| \epsilon_n < 1$, A_n has a bounded inverse and $\|A_n^{-1}\| = O(1)$, $A_n : X_n \rightarrow X_n$. Since $\|y - P_n y\|_\rho = O\{E_n(y)_C\} = \delta_n$. Finally, we have

$$\begin{aligned} \|x^* - x_n^*\|_\rho &= \|A^{-1}y - A_n^{-1}P_n y\|_\rho \\ &= O(\epsilon_n + \delta_n) \\ &= O\{E_n^t(h)_C + E_n^\tau(h)_C + E_n(y)_C\}. \end{aligned} \quad (3.9)$$

□

LEMMA 3.1. *If Q_n is an algebraic polynomial of degree $n - 1$, then*

$$\|Q_n(t)\|_C \leq n \sqrt{\frac{n}{2}} \|Q_n\|_\rho, \quad n \geq 2. \quad (3.10)$$

PROOF. One may write Q_n as $Q_n(t) = \sum_{k=1}^n C_{k-1}(Q_n) U_{k-1}(t)$, $n \in \mathbb{N}$, where $C_j(Q_n) = (Q_n, U_j)_\rho$. Since $|U_{k-1}| \leq k$ it follows that

$$\begin{aligned} \|Q_n(t)\|_C &\leq \left\{ \sum_{k=1}^n C_{k-1} | (Q_n)^2 | \right\}^{1/2} \left\{ \sum_{k=1}^n k^2 \right\}^{1/2} \\ &= \|Q_n\|_\rho \sqrt{n(n+1)(2n+1)/6} \\ &\leq n \sqrt{n/2} \|Q_n\|_\rho \\ &= \|Q_n\|_\rho O(n^{3/2}). \end{aligned} \quad (3.11) \quad \square$$

Define $W^r H_\alpha = \{x : x^{(r-1)} \text{ is absolutely continuous, } x^{(r)} \in H_\alpha\}$.

THEOREM 3.2. *Assume that conditions (i) and (ii) of Theorem 3.1 are satisfied,*

$$y(t) \in W^r H_\alpha, \quad h(t, \tau) \in W^r H_\alpha, \quad r \geq 0, \quad 0 < \alpha \leq 1, \quad (3.12)$$

then

$$\|x^* - x_n^*\|_\rho = O(n^{-r-\alpha}), \quad (3.13)$$

$$\|x^* - x_n^*\|_C = O(n^{3/2-r-\alpha}). \quad (3.14)$$

PROOF. Since $y(t) \in W^r H_\alpha$, $h(t, \tau) \in W^r H_\alpha$, then, according to [7], one has $E_n(y) = O(n^{-r-\alpha})$, $E_n^t(h) = O(n^{-r-\alpha})$, $E_n^\tau(h) = O(n^{-r-\alpha})$, $r + \alpha > 0$. This proves (3.13). It is easy to show that

$$\begin{aligned} \|x^* - x_n^*\|_C &= \sum_{k=1}^{\infty} \|x_{2^k n}^* - x_{2^{k-1} n}^*\|_C \\ &\leq \sum_{k=1}^{\infty} (2^k n)^{(3/2)} \left[\|x^* - x_{2^k n}^*\|_\rho + \|x^* - x_{2^{k-1} n}^*\|_\rho \right] \\ &\leq C_3 \sum_{k=1}^{\infty} (2^k n)^{(3/2)} (2^k n)^{-r-\alpha} = C_4 n^{-r-\alpha+3/2}, \end{aligned} \quad (3.15)$$

where C_3, C_4 are constants. This proves (3.14). \square

Define $C_\rho[-1, 1] = \{x : \sqrt{1-t^2}x \in C[-1, 1]\}$ and $\|x\|_{C_\rho} = \max \{\sqrt{1-t^2}|x(t)|\}$.

LEMMA 3.2. *If Q_n is a polynomial of degree $n - 1$, then*

$$\|Q_n(t)\|_{C_\rho} \leq \sqrt{n} \|Q_n\|_X. \quad (3.16)$$

PROOF. Since $|U_m(t)| \leq (1-t^2)^{-1/2}$, $-1 \leq t \leq 1$, $m = 0, 1, \dots$, then

$$\begin{aligned} \sqrt{1-t^2} |Q_n(t)| &\leq \sum_{k=1}^n |C_{k-1}(Q_n)| \\ &\leq \left\{ \sum_{k=1}^n |C_{k-1}(Q_n)|^2 \right\}^{1/2} \sqrt{n} = \sqrt{n} \|Q_n\|_\rho. \end{aligned} \quad (3.17) \quad \square$$

THEOREM 3.3. *If $\|x^* - x_n^*\|_X = O(n^{-m})$, $m \in \mathbb{R}$, then*

$$\|x^* - x_n^*\|_{C_\rho} = O(n^{1/2-m}), \quad m > \frac{1}{2}, \quad (3.18)$$

PROOF. Using Lemma 3.2 and the same technique as in Theorem 3.2, one obtains (3.18). \square

CONCLUSION. For the Mechanical Quadratures methods, the rate of convergence in space $C_\rho[-1, 1]$ is better than that given in space $C[-1, 1]$.

4. Approximation by degenerate kernels. We approximate $h(t, \tau)$ by

$$h_n(t, \tau) = \sum_{k=1}^n a_k(t) b_k(\tau), \quad -1 \leq t, \tau \leq 1, \quad (4.1)$$

where $\{a_k\}, \{b_k\}$ are two sets of linearly independent functions. By substituting back into (1.1), we obtain

$$A_n x = Gx + \int_{-1}^1 \sqrt{1-\tau^2} h_n(t, \tau) x(\tau) d\tau = y. \quad (4.2)$$

(4.2) can be written as

$$Gx + \sum_{k=1}^n \alpha_k a_k = y, \quad (4.3)$$

where $\alpha_k = \int_{-1}^1 \sqrt{1-\tau^2} x(\tau) b_k(\tau) d\tau$. Then the solution of (4.3) is given by

$$\begin{aligned} x &= G^{-1} y - \sum_{k=1}^n \alpha_k G^{-1} a_k, \\ G^{-1} y &= - \sum_{k=1}^{\infty} \frac{(\gamma, U_{k-1})_\rho}{k} U_{k-1}(t). \end{aligned} \quad (4.4)$$

Multiplying (4.4) by $\sqrt{1-t^2} b_j(t)$ and integrating, we get the linear system of equations

$$\alpha_j + \sum_{k=1}^n \gamma_{jk} \alpha_k = \gamma_j, \quad 1 \leq j \leq n, \quad (4.5)$$

where $\gamma_{jk} = \int_{-1}^1 \sqrt{1-t^2} b_j G^{-1}(a_k) dt$, $\gamma_j = \int_{-1}^1 \sqrt{1-t^2} b_j G^{-1}(y) dt$. Define $q = \sqrt{1-t^2} \sqrt{1-\tau^2}$.

THEOREM 4.1. *Suppose that*

- (i) $y \in L_{2,\rho}$
- (ii) $h \in L_{2,q}[-1, 1; -1, 1]$
- (iii) $\epsilon_n^2 = \int_{-1}^1 \int_{-1}^1 q(t, \tau) |h(t, \tau) - h_n(t, \tau)|^2 dt d\tau \rightarrow 0, n \rightarrow \infty$
- (iv) *equation (1.1) has a unique solution.*

Then for all n such that $q_n = \epsilon_n \|A^{-1}\| < 1$, $A : X \rightarrow X$, the linear system of equations (4.5) has a unique solution $\{\alpha_k^\}_1^n$ and the approximate solution $x_n^* = G^{-1}(y) - \sum_{k=1}^n \alpha_k^* G^{-1}(a_k)$ converges to the exact solution x^* , $\|x^* - x_n^*\|_\rho = O(\epsilon_n)$.*

PROOF. For $x \in X$, we have

$$\begin{aligned} \|Ax - A_n x\|_\rho &= \left\| \int_{-1}^1 \sqrt{1-\tau^2} [h(t, \tau) - h_n(t, \tau)] x(\tau) d\tau \right\|_\rho \\ &\leq \|x\|_\rho \|h - h_n\|_{L_{2,q}} = \epsilon_n \|x\|_\rho. \end{aligned} \quad (4.6)$$

Then $\|A - A_n\|_{X \rightarrow X} \leq \epsilon_n$, according to [6] (4.5) has a unique solution $\{\alpha_k^*\}_1^n$, $\|x^* - x_n^*\| = O(\epsilon_n)$. \square

REFERENCES

- [1] B. G. Gabdulkhayev, *Optimalnye approksimatsii reshenii lineinykh zadach. [Optimal approximations of solutions of linear problems]*, Kazan. Gos. Univ., Kazan, 1980 (Russian). MR 83k:65047.
- [2] M. A. Golberg, *The convergence of several algorithms for solving integral equations with finite-part integrals*, J. Integral Equations **5** (1983), no. 4, 329-340. MR 85h:65273. Zbl 529.65079.
- [3] ———, *The convergence of several algorithms for solving integral equations with finite part integrals. II*, J. Integral Equations **9** (1985), no. 3, 267-275. MR 87j:65162a. Zbl 606.65090.
- [4] J. Hadamard, *Lectures on Cauchy's problem in linear partial differential equations*, Dover Publications, New York, 1952. MR 14,474f. Zbl 049.34805.
- [5] N. I. Ioakimidis, *Two methods for the numerical solution of Bueckner's singular integral equation for plane elasticity crack problems*, Comput. Methods Appl. Mech. Engrg. **31** (1982), no. 2, 169-177. MR 83i:73052. Zbl 484.73074.
- [6] L. V. Kantorovic and G. P. Akilov, *Funktsionalnyi analiz v normirovannykh prostranstvakh. [Functional analysis in normed spaces]*, Gosudarstv. Izdat. Fis. Mat. Lit., Moscow, 1959 (Russian). MR 22#9837. Zbl 127.06102.
- [7] G. Meinardus, *Approximation of functions: Theory and numerical methods.*, Expanded translation of the German edition. Translated by Larry L. Schumaker. Springer Tracts in Natural Philosophy, vol. 13, Springer-Verlag, New York, 1967. MR 36#571. Zbl 152.15202.
- [8] A. F. Timan, *Teorij pribli+enij funktsii deistvitel'nogo peremennogo. [Theory of approximation of functions of a real variable]*, Gosudarstv. Izdat. Fiz. Mat. Lit., Moscow, 1960 (Russian). MR 22#8257. Zbl 117.29001.
- [9] D. E. Williams, *Some mathematical methods in three-dimensional subsonic flutter-derivatives theory*, British Ministry of Aviation Reports and Memo, No.3302 (1961).

ASHOUR: DEPARTMENT OF MATHEMATICS, CAIRO UNIVERSITY, GIZA, EGYPT

Special Issue on Time-Dependent Billiards

Call for Papers

This subject has been extensively studied in the past years for one-, two-, and three-dimensional space. Additionally, such dynamical systems can exhibit a very important and still unexplained phenomenon, called as the Fermi acceleration phenomenon. Basically, the phenomenon of Fermi acceleration (FA) is a process in which a classical particle can acquire unbounded energy from collisions with a heavy moving wall. This phenomenon was originally proposed by Enrico Fermi in 1949 as a possible explanation of the origin of the large energies of the cosmic particles. His original model was then modified and considered under different approaches and using many versions. Moreover, applications of FA have been of a large broad interest in many different fields of science including plasma physics, astrophysics, atomic physics, optics, and time-dependent billiard problems and they are useful for controlling chaos in Engineering and dynamical systems exhibiting chaos (both conservative and dissipative chaos).

We intend to publish in this special issue papers reporting research on time-dependent billiards. The topic includes both conservative and dissipative dynamics. Papers discussing dynamical properties, statistical and mathematical results, stability investigation of the phase space structure, the phenomenon of Fermi acceleration, conditions for having suppression of Fermi acceleration, and computational and numerical methods for exploring these structures and applications are welcome.

To be acceptable for publication in the special issue of Mathematical Problems in Engineering, papers must make significant, original, and correct contributions to one or more of the topics above mentioned. Mathematical papers regarding the topics above are also welcome.

Authors should follow the Mathematical Problems in Engineering manuscript format described at <http://www.hindawi.com/journals/mpe/>. Prospective authors should submit an electronic copy of their complete manuscript through the journal Manuscript Tracking System at <http://mts.hindawi.com/> according to the following timetable:

Manuscript Due	December 1, 2008
First Round of Reviews	March 1, 2009
Publication Date	June 1, 2009

Guest Editors

Edson Denis Leonel, Departamento de Estatística, Matemática Aplicada e Computação, Instituto de Geociências e Ciências Exatas, Universidade Estadual Paulista, Avenida 24A, 1515 Bela Vista, 13506-700 Rio Claro, SP, Brazil ; edleonel@rc.unesp.br

Alexander Loskutov, Physics Faculty, Moscow State University, Vorob'evy Gory, Moscow 119992, Russia; loskutov@chaos.phys.msu.ru