IJMMS 2004:61, 3259–3274
PII. S0161171204402245
http://ijmms.hindawi.com

ON THE CRITICAL PERIODS OF LIÉNARD SYSTEMS WITH CUBIC RESTORING FORCES

ZHENGDONG DU

Received 23 February 2004

We study local bifurcations of critical periods in the neighborhood of a nondegenerate center of a Liénard system of the form $\dot{x} = -y + F(x)$, $\dot{y} = g(x)$, where $F(x)$ and $g(x)$ are polynomials such that $\deg(g(x)) \le 3$, $g(0) = 0$, and $g'(0) = 1$, $F(0) = F'(0) = 0$ and the system always has a center at $(0,0)$. The set of coefficients of $F(x)$ and $g(x)$ is split into two strata denoted by S_I and S_{II} and $(0,0)$ is called weak center of type I and type II, respectively. By using a similar method implemented in previous works which is based on the analysis of the coefficients of the Taylor series of the period function, we show that for a weak center of type I, at most $[(1/2)\deg(F(x))]-1$ local critical periods can bifurcate and the maximum number can be reached. For a weak center of type II, the maximum number of local critical periods that can bifurcate is at least $[(1/4)\deg(F(x))]$.

2000 Mathematics Subject Classification: 34C23, 37G10.

1. Introduction. During the last decades, there has been considerable interest in studying the generalized Liénard system of the form

$$\dot{x} = -y + F(x), \qquad \dot{y} = g(x), \tag{1.1}$$

or its equivalent form

$$\dot{x} = -y, \qquad \dot{y} = g(x) + f(x)y, \tag{1.2}$$

where $f(x) = F'(x)$. The popularity is due to at least two reasons. First, it generalizes many oscillation systems arising from applications. Second, many other systems can be transformed into the form (1.1) or (1.2) (see [1]).

One of the most studied problems is to determine the number and relative configuration of limit cycles of (1.1) in terms of the properties of $F(x)$ and $g(x)$. There is an enormous literature on this problem, see, for example, [15] for more on these issues. For the special class of system (1.2), where $f(x)$ and $g(x)$ are polynomials of degrees at most n and m, respectively, there are also extensive studies of the cyclicity $\hat{H}_{n,m}$, that is, the maximum number of small amplitude limit cycles bifurcating from the fine focus of (1.2). In [7], Christopher and Lynch give results for $\hat{H}_{n,m}$ when $f(x)$ or $g(x)$ is quadratic or cubic polynomial.

The purpose of this paper is to examine the local bifurcations of critical periods in the neighborhood of a nondegenerate center of system (1.1), where $F(x)$ and $g(x)$ are polynomials such that $\deg(g(x)) \le 3$, $g(0) = 0$, and $g'(0) = 1$, $F(0) = F'(0) = 0$ and

the system always has a center at $(0,0)$. Let $G(x)=\int_0^x g(\xi)d\xi$, the center condition for polynomial Liénard systems is given by the following theorem of Christopher (see [5]).

Theorem 1.1 [5]. *If $F(x)$ and $g(x)$ are polynomials, then system (1.1) has a nondegenerate center at $(0,0)$ if and only if $F(x)$ and $G(x)$ are both polynomials of a polynomial $M(x)$ with $M(0)=M'(0)=0$ and $M''(0)\neq 0$.*

The bifurcation of critical periods from centers of planar vector fields is an important problem, because it is closely related to the monotonicity of periods of closed orbits surrounding a center and subharmonic bifurcation for periodically forced systems. Similar to Hilbert's 16th problem, the following problem can be formulated.

Problem 1.2. Determine the maximum number $\mathscr{C}(n)$ of critical periods of polynomial systems of degree n with nondegenerate centers in terms of n only.

While Problem 1.2 is still completely open, an easier problem is proposed.

Problem 1.3. Determine the maximum number $\tilde{\mathscr{C}}(n)$ of local critical periods bifurcating from a weak center of polynomial systems of degree n in terms of n only.

In 1989, Chicone and Jacobs (see [3]) developed a general theory of solving Problem 1.3 and proved that $\tilde{\mathscr{C}}(2)=2$. However, the problem for higher-degree systems is still unsolved. A few classes of cubic systems studied in [10, 11, 13, 14] proved that $\tilde{\mathscr{C}}(3)\geq 4$. It is worth noting that some researchers have considered the global problem of bifurcations of critical periods for some specific systems, see, for example, [9]. However, there is still no general method of solving Problem 1.2.

The monotonicity of the period function of centers of system (1.2) or isochronicity has been studied by several authors (see [12] and the references therein). Recently, Christopher and Devlin [6] gave a complete classification for isochronous centers of polynomial Liénard systems of degree 34 or less. However, there are only very few studies on the number of critical periods that can bifurcate from the nondegenerate center of (1.1) for the special case where $F(x)\equiv 0$, which significantly simplifies matters (see [3, 4]).

Under our assumption, $\deg(G(x))\leq 4$ and by Theorem 1.1, both $F(x)$ and $G(x)$ are polynomials of $M(x)$, where $M(x)$ is as in Theorem 1.1. Without loss of generality, we assume that $M''(0)=2$. Thus, $M(x)=x^2+O(x^3)$ and $\deg(M(x))\leq 4$. So, we write $M(x)=x^2+b_1x^3+b_2x^4$ and $F(x)$, $G(x)$ have the form

$$F=\sum_{k=1}^{n}a_kM^k,\qquad G=\frac{1}{2}M+b_0M^2. \tag{1.3}$$

Clearly, there are only 2 possibilities: $b_0=0$ or $b_0\neq 0$ and $M(x)=x^2$. To unify these two cases, we write system (1.1) into the following form:

$$\dot{x}=-y+\sum_{k=1}^{n}a_k(x^2+b_1x^3+b_2x^4)^k,\qquad \dot{y}=x+\frac{3}{2}b_1x^2+(b_0+2b_2)x^3. \tag{1.4}$$

Thus, $F(x)$ is as in (1.3) and $g(x)$, $G(x)$ are given below:

$$g(x) = x + \frac{3}{2}b_1x^2 + (b_0 + 2b_2)x^3, \qquad G(x) = \frac{1}{2}M(x) + \frac{1}{4}b_0x^4. \tag{1.5}$$

For the coefficients $(b_0, b_1, b_2, a_1, \ldots, a_n)$ we will use the abbreviation λ so that $\lambda_k = b_k$ for $k = 0, 1, 2$ and $\lambda_{k+2} = a_k$ for $k = 1, \ldots, n$. In particular, we have $\lambda \in \mathbb{R}^{n+3}$. Then, from Theorem 1.1 we have the following lemma.

Lemma 1.4. *System (1.4) has a nondegenerate center at* $(0,0)$ *if and only if the parameter value* $\lambda \in \mathbb{R}^{n+3}$ *is in one of the following strata:*

$$S_I := \{\lambda \in \mathbb{R}^{n+3} \mid \lambda_0 \neq 0,\ \lambda_1 = \lambda_2 = 0\}, \qquad S_{II} := \{\lambda \in \mathbb{R}^{n+3} \mid \lambda_0 = 0\}. \tag{1.6}$$

We say system (1.4) has a weak center of *type I* (resp., *type II*) if the system is nonlinear and $\lambda \in S_I$ (resp., $\lambda \in S_{II}$). Our main result is the following theorem.

Theorem 1.5. (1) *System (1.4) has an isochronous center at the origin if and only if* $b_0 = (4/9)a_0^2$ *and* $b_1 = b_2 = a_1 = \cdots = a_n = 0$. *At most* $n-1$ *critical periods can bifurcate from the nonlinear isochronous center and there are perturbations to produce* $n-1$ *critical periods.*

(2) *If system (1.4) has a weak center of type I of finite order at the origin, then at most* $n-1$ *critical periods can bifurcate from the weak center and there are perturbations to produce* $n-1$ *critical periods.*

(3) *The maximum critical periods can bifurcate from a weak center of type II of system (1.4) is at least* n.

It is clear from Theorem 1.5 that there are no nonlinear isochronous centers inside type II. For a weak center of type I, $\deg(F(x)) = 2n$, the maximum local critical periods that can bifurcate are $[(1/2)\deg(F(x))] - 1$ and this upper bound can be attained. For a weak center of type II, $\deg(F(x)) = 4n$, the maximum local critical periods can bifurcate is at least $[(1/4)\deg(F(x))]$.

Our approach is similar to the one implemented in [3, 10, 11, 13]. It is based on the analysis of the coefficients of the Taylor series of the period function. The Taylor coefficients of the period function have been computed and simplified by reduction modulo a Gröbner basis using Maple for low degrees of $F(x)$. This enables us to conjecture a general pattern for the ideal generated by the coefficients over the polynomial ring of the parameters. These conjectures are then proved rigorously using arguments similar to those used by Bautin in [2] to determine the structure of ideal generated by focal values of quadratic system.

This paper is organized as follows. Section 2.1 summarizes the general results by Chicone and Jacobs (see [3]). Section 2.2 summarizes a recursive algorithm to compute the period coefficients. Section 3 considers weak center of type I. Section 4 considers weak center of type II. Section 5 is the proof of Theorem 1.5.

2. Preliminary

2.1. Local critical periods of polynomial systems. Consider a family of planar ordinary differential systems with a nondegenerate center at the origin

$$\dot{x} = -y + p(x,y,\nu), \qquad \dot{y} = x + q(x,y,\nu), \tag{2.1}$$

where $p(x,y,\nu)$ and $q(x,y,\nu)$ are polynomials of degree n in x and y, the parameter $\nu = (\nu_1,\dots,\nu_m) \in \mathbb{R}^m$. The minimum period $T(r,\nu)$ of the periodic orbit passing through $(r,0)$ associated with the center of the above system for the parameter ν yields the so-called period function T. Let $P(r,\nu) = T(r,\nu) - 2\pi$. For $\nu_* \in \mathbb{R}^m$, the origin is called a *weak center of finite order* k if $P(0,\nu_*) = P'(0,\nu_*) = \cdots = P^{(2k+1)}(0,\nu_*) = 0$ and $P^{(2k+2)}(0,\nu_*) \neq 0$, where the derivatives are taken with respect to r. The origin is called an *isochronous center* if $P^{(k)}(0,\nu_*) = 0$ for all $k \geq 0$. A *critical period* is a period corresponding to a solution of the equation $P_r(r,\nu) = 0$ as ν varies. A *local critical period* is a period corresponding to a critical point of $P(r,\nu)$ which arises from a bifurcation from a weak center.

For ν_* corresponding to a weak center, the function $(r,\nu) \mapsto T(r,\nu)$ is analytic in a neighborhood of $(0,\nu_*)$ and can be represented by its Taylor series

$$T(r,\nu) = 2\pi + \sum_{k=2}^{\infty} p_k(\nu) r^{2k}, \tag{2.2}$$

for $|r|$ and $|\nu - \nu_*|$ sufficiently small. Here, the period coefficients $p_k \in \mathbb{R}[\nu_1,\dots,\nu_m]$, the Noetherian ring of polynomials in the variables $\nu_1,\dots,\nu_m$ and for any $k \geq 1$, $p_{2k+1} \in (p_2,p_4,\dots,p_{2k})$, the ideal generated by $p_2,p_4,\dots,p_{2k}$. In particular, for any ν, the first $k \geq 1$ such that $p_k(\nu) \neq 0$ is even (see [3]). The theory of Chicone and Jacobs in [3] is based on the analysis of the period coefficients. To state their theorems precisely, we first introduce the following concept.

Definition 2.1 [11]**.** Let $\{\chi_\nu\}_{\nu\in\mathbb{R}^m}$ be a family of systems with a center at the origin and associated period coefficients $p_{2k}(\nu)$. The family is said to satisfy condition $(\mathcal{P})$ if for any $\nu_* \in V(p_2,p_4,\dots,p_{2k}) := \{\nu \mid p_{2i}(\nu) = 0,\ i = 1,\dots,k\}$, $p_{2k+2}(\nu_*) \neq 0$ and any neighborhood $W \subset \mathbb{R}^m$ of ν_* in which $p_{2k+2} \neq 0$, there exists $\nu' \in W$ such that

$$p_{2k}(\nu')p_{2k+2}(\nu') < 0, \quad \text{with } \nu' \in V(p_2,p_4,\dots,p_{2k-2}). \tag{2.3}$$

The system χ_{ν_*} is said to satisfy condition $(\mathcal{P}_k)$.

The following version of the theorems of Chicone and Jacobs [3] are given by Rousseau and Toni in [11].

Finite-order bifurcation theorem. *From weak centers of finite order k at the parameter value ν_*, no more than k local critical periods bifurcate. Moreover, if the family satisfies the condition $(\mathcal{P})$ and if χ_{ν_*} satisfies the condition $(\mathcal{P}_k)$, then there are perturbations with exactly j local critical periods for any $0 \leq j \leq k$.*

ISOCHRONE BIFURCATION THEOREM. *If the vector field (2.1) has an isochronous center at the origin for the parameter value ν_* and if for each integer $n \geq 1$, the period coefficient p_{2n} is in the ideal $(p_2, p_4, \ldots, p_{2k}, p_{2k+2})$ over the ring $\mathbb{R}\{\nu_1, \ldots, \nu_m\}_{\nu_*}$ of convergent power series at ν_*, then at most k local critical periods bifurcate from the isochronous center at ν_*. Moreover, if the family satisfies the condition $(\mathcal{P})$ and if χ_{ν_*} satisfies the condition $(\mathcal{P}_k)$, then exactly j local critical periods bifurcate from the center at ν_* for any $0 \leq j \leq k$.*

2.2. The computation of the periodic coefficients. Let $g_0(x) = g(x) - x$ and transform (1.4) to polar coordinates by $x = r\cos\theta$, $y = r\sin\theta$, and eliminating time yields

$$\frac{dr}{d\theta} = \frac{F(r\cos\theta)\cos\theta + g_0(r\cos\theta)\sin\theta}{1 + (1/r)g_0(r\cos\theta)\cos\theta - (1/r)F(r\cos\theta)\sin\theta}. \tag{2.4}$$

Then, (2.4) is analytic and we assume the following expansion:

$$\frac{dr}{d\theta} = \sum_{k=2}^{\infty} A_k(\theta) r^k, \qquad \frac{dt}{d\theta} = \sum_{k=0}^{\infty} B_k(\theta) r^k, \tag{2.5}$$

where $B_0(\theta) \equiv 1$. Let γ_ξ be the closed orbit of (1.4), through $(\xi, 0)$. The period function is given by

$$T(\xi,\lambda) = \int_{\gamma_\xi} dt = \int_0^{2\pi} \frac{dt}{d\theta} d\theta = 2\pi + \int_0^{2\pi} \sum_{k=1}^{\infty} B_k(\theta) r^k d\theta, \tag{2.6}$$

where $r = r(\theta, \xi, \lambda)$ is the solution of (1.4), with the initial condition $r(0,\lambda) = \xi$. $r(\theta,\xi,\lambda)$ may be locally represented as a convergent power series in ξ:

$$r(\theta,\xi,\lambda) = \sum_{k=1}^{\infty} u_k(\theta,\lambda)\xi^k, \tag{2.7}$$

where $u_1(0,\lambda) = 1$ and $u_k(0,\lambda) = 0$ for any $k > 1$ and λ. Substituting (2.7) into (2.5) and comparing the coefficients of ξ^k, $k \geq 1$, we obtain recursive equations for u_k. For example, the first 3 equations are given by

$$\frac{du_1}{d\theta} = 0, \qquad \frac{du_2}{d\theta} = A_2 u_1^2, \qquad \frac{du_3}{d\theta} = A_3 u_1^3 + 2A_2 u_1 u_2, \tag{2.8}$$

which can be found by direct integration. From (2.6) and (2.7), we obtain a recursive algorithm for computing the period coefficients $p_2, p_4, \ldots, p_{2k}, \ldots$. The algorithm can be easily implemented in the computer algebra system Maple.

3. Weak center of type I. A weak center of type I corresponds to the case where $F(x)$ is a polynomial of degree $2n$ and system (1.4) has the following form:

$$\dot{x} = -y + \sum_{k=1}^{n} a_k x^{2k}, \qquad \dot{y} = x + b_0 x^3, \tag{3.1}$$

where $b_0 \neq 0$. Direct computation with the aid of Maple yields the following lemma.

Lemma 3.1. *The period coefficient p_2 for (3.1) is given by $p_2 = (\pi/12)(4a_1^2 - 9b_0)$.*

It is convenient to split the set S_I as $S_I = S_I^A \bigcup S_I^B$, where

$$\begin{aligned} S_I^A &:= \{\lambda \in \mathbb{R}^{n+3} \mid \lambda_0 < 0,\ \lambda_1 = \lambda_2 = 0\}, \\ S_I^B &:= \{\lambda \in \mathbb{R}^{n+3} \mid \lambda_0 > 0,\ \lambda_1 = \lambda_2 = 0\}. \end{aligned} \tag{3.2}$$

Following our notation, we say system (1.4) has a weak center of *type I_A* (resp., *type I_B*) if the system is nonlinear and $\lambda \in S_I^A$ (resp., $\lambda \in S_I^B$). We consider weak centers of type I_A and type I_B separately. Obviously, $p_2 > 0$ for $b_0 < 0$ by Lemma 3.1. Thus, we have the following theorem.

Theorem 3.2. *A weak center of type I_A has order 0 and no local critical periods can bifurcate from a weak center of type I_A.*

Now, we discuss weak centers of type I_B. In this case, the system has the same form as (3.1) with $b_0 > 0$. We have the following lemma.

Lemma 3.3. *Suppose $b_0 > 0$ in (3.1). Then, if $b_0 = (4/9)a_1^2$ and $a_2 = \cdots = a_n = 0$, the origin is an isochronous center.*

Proof. Since $b_0 > 0$, the assumption of Lemma 3.3 implies that $a_1 \neq 0$. Perform coordinate transformation $(x, y) \mapsto ((2/3)a_1 x, (2/3)a_1 y)$, system (3.1) becomes

$$\dot{x} = -y + \frac{3}{2}x^2, \qquad \dot{y} = x + x^3. \tag{3.3}$$

It can be linearized by the change of coordinates $(u, v) = (-2x/(x^2 - 2y - 2), (x^2 - 2y)/(x^2 - 2y - 2))$ and has first integral $F(x, y) = -4(2x^2 - 2y - 1)/(x^2 - 2y - 2)^2$. By [8, Theorem 3.2], the origin is an isochronous center. □

To simplify the computation, we scale system (3.1) so that $b_0 = 1$. Then, system (3.1) has the form

$$\dot{x} = -y + \sum_{k=1}^{n} a_k x^{2k}, \qquad \dot{y} = x + x^3. \tag{3.4}$$

Lemma 3.4. *If $2 \le k \le n$, then the period coefficient p_{2k} of system (3.4) has the form $p_{2k} = \beta a_1 a_k + q(a_1, \ldots, a_{k-1})$, where $\beta \neq 0$ is a constant and $q(a_1, \ldots, a_{k-1}) \in \mathbb{R}[a_1, \ldots, a_{k-1}]$.*

Proof. The expansion (2.4) for system (3.4) is given by

$$\begin{aligned} \frac{dr}{d\theta} &= \frac{\sum_{k=1}^{n} (a_k \cos^{2k+1}\theta) r^{2k} + (\cos^3\theta \sin\theta) r^3}{1 - \sum_{k=1}^{n} (a_k \cos^{2k}\theta \sin\theta) r^{2k-1} + (\cos^4\theta) r^2} \\ &= \left[\sum_{k=1}^{n} (a_k \cos^{2k+1}\theta) r^{2k} + (\cos^3\theta \sin\theta) r^3 \right] \\ &\quad \cdot \left[1 + \sum_{m=1}^{\infty} (-1)^m \left(\sum_{k=1}^{n} (a_k \cos^{2k}\theta \sin\theta) r^{2k-1} + (\cos^4\theta) r^2 \right)^m \right]. \end{aligned} \tag{3.5}$$

Similarly, we have

$$\frac{dt}{d\theta} = 1 + \sum_{m=1}^{\infty} (-1)^m \left[\sum_{k=1}^{n} (a_k \cos^{2k}\theta \sin\theta) r^{2k-1} + (\cos^4\theta) r^2 \right]^m. \tag{3.6}$$

Obviously, for any $j \le 2(k-1)$, where $2 \le k \le n$, the coefficients A_j and B_j in the corresponding expansion (2.5) are independent of a_k.

On the other hand, $B_1 = a_1 \sin\theta \cos^2\theta$ and for $2 \le k \le n$,

$$\begin{gathered} A_{2k} = a_k \cos^{2k+1}\theta + \tilde{A}_{2k}, \qquad B_{2k-1} = a_k \cos^{2k}\theta \sin\theta + \tilde{B}_{2k-1}, \\ B_{2k} = 2a_1 a_k \cos^{2k+2}\theta \sin^2\theta + \tilde{B}_{2k}, \end{gathered} \tag{3.7}$$

where $\tilde{A}_{2k}$, $\tilde{B}_{2k-1}$, $\tilde{B}_{2k}$ are polynomials only depending on $a_1,\dots,a_{k-1}$. Furthermore, from (2.5) and (2.7), it is clear that $u_1,\dots,u_{2k-1}$ only depend on $a_1,\dots,a_{k-1}$. u_{2k} satisfies the following initial value problem:

$$\frac{du_{2k}}{d\theta} = A_{2k} u_1^{2k} + D_{u_{2k}}, \quad u_{2k}(0,\lambda) = 0, \tag{3.8}$$

where $D_{u_{2k}}$ only depends on $A_1,\dots,A_{2k-1}$ and $u_1,\dots,u_{2k-1}$. Solving (3.8) yields $u_{2k} = a_k P_{u_{2k}}(\cos\theta,\sin\theta) + \tilde{u}_{2k}$, where $P_{u_{2k}}(\cos\theta,\sin\theta)$ is a polynomial of $\cos\theta$ and $\sin\theta$ that, independent of $a_1,\dots,a_n$, $\tilde{u}_{2k}$, is a polynomial only depending on $a_1,\dots,a_{k-1}$. Hence, the period coefficient p_{2k} can be computed as follows:

$$p_{2k}(\lambda) = \int_0^{2\pi} (B_1 u_{2k} + (2k-1) B_{2k-1} u_1^{2k-2} u_2 + B_{2k} u_1^{2k} + \tilde{B}_u)\, d\theta, \tag{3.9}$$

where $\tilde{B}_u$ only depends on $B_2,\dots,B_{2k-2}$ and $u_1,\dots,u_{2k-1}$. Thus, p_{2k} has the form $p_{2k} = \beta a_1 a_k + q(a_1,\dots,a_{k-1})$, where $\beta \ne 0$ is a constant and $q(a_1,\dots,a_{k-1}) \in \mathbb{R}[a_1,\dots,a_{k-1}]$. □

Direct computation of period coefficients using the Gröbner base package of Maple for $n \le 8$ suggests the following Lemma which is proved rigorously.

Lemma 3.5. *For system (3.4), the period coefficients p_{2k} $(k > 1)$, reduced modulo the ideal generated by $p_2,\dots,p_{2k-2}$ and omitting the constant factor, are given by $p_{2k} = a_1 a_k$ for $1 < k \le n$ and $p_{2k} = 0$ for $k > n$. In particular, $p_{2k} \in (p_2,\dots,p_{2n})$, the ideal of the polynomial ring $\mathbb{R}[a_1,a_2,\dots,a_n]$.*

Proof. We prove Lemma 3.5 by induction on k. By direct computation, we find that p_4 modulo the ideal generated by p_2 is given by $p_4 = (2\pi/3) a_1 a_2$. Thus, Lemma 3.5 is true for $k = 2$.

Now, assume that $2 \le k \le n$ and $p_4,\dots,p_{2k-2}$, with each reduced modulo the ideal generated by the previous coefficients, and omitting the constant factor are given by $p_{2j} = a_1 a_j$ $(j = 2,\dots,k-1)$. By Lemma 3.4, $p_{2k} = \beta a_1 a_k + q(a_1,\dots,a_{k-1})$, where $\beta \ne 0$ is a constant and $q(a_1,\dots,a_{k-1}) \in \mathbb{R}[a_1,\dots,a_{k-1}]$.

Because p_2 is a quadratic polynomial of a_1, $q(a_1,\dots,a_{k-1})$ can be written into the form $q(a_1,\dots,a_{k-1}) = q_0(a_1,\dots,a_{k-1})p_2 + a_1q_1(a_2,\dots,a_{k-1}) + q_2(a_2,\dots,a_{k-1})$. Thus, we have

$$p_{2k} = \beta a_1a_k + q_0(a_1,\dots,a_{k-1})p_2 + a_1q_1(a_2,\dots,a_{k-1}) + q_2(a_2,\dots,a_{k-1}). \tag{3.10}$$

By Lemma 3.3, when $p_2 = 0$ and $a_2 = \cdots = a_k = 0$, $p_{2k} = 0$. Thus, $q_1(0,\dots,0) = q_2(0,\dots,0) = 0$. Hence, they have the following form:

$$\begin{aligned} q_1(a_2,\dots,a_{k-1}) &= \sum_{j=2}^{k-1} a_jS_j(a_2,\dots,a_{k-1}),\\ q_2(a_2,\dots,a_{k-1}) &= \sum_{j=2}^{k-1} a_jR_j(a_2,\dots,a_{k-1}), \end{aligned} \tag{3.11}$$

where $S_j(a_2,\dots,a_{k-1})$ and $R_j(a_2,\dots,a_{k-1})$ are polynomials of $a_2,\dots,a_{k-1}$. By induction hypothesis, $a_1a_j = p_{2j}$ $(j = 2,\dots,k-1)$. Thus, (3.10) becomes

$$p_{2k} = \beta a_1a_k + q_0p_2 + \sum_{j=2}^{k-1} p_{2j}S_j + \sum_{j=2}^{k-1} a_jR_j. \tag{3.12}$$

On the other hand, we have $\mu p_2 + \gamma a_1^2 = 1$, where $\mu = -4/3\pi$, $\gamma = 4/9$. Therefore, by induction hypothesis, for each a_j $(2 \le j \le k-1)$, we have

$$a_j = \mu a_jp_2 + \gamma a_1(a_1a_j) = \mu a_jp_2 + \gamma a_1p_{2j}. \tag{3.13}$$

Substituting (3.13) into (3.12) yields

$$p_{2k} = \beta a_1a_k + q_0p_2 + \sum_{j=2}^{k-1} p_{2j}S_j + \sum_{j=2}^{k-1} (\mu a_jp_2 + \gamma a_1p_{2j})R_j. \tag{3.14}$$

Thus, p_{2k}, reduced modulo $(p_2,\dots,p_{2(k-1)})$ and omitting the constant factor, is a_1a_k. This completes the proof of Lemma 3.5 for $1 < k \le n$.

Now, assume that $k > n$. Then, p_{2k} has the form $p_{2k} = V_0(a_1,\dots,a_n)p_2 + a_1V_1(a_2,\dots,a_n) + V_2(a_2,\dots,a_n)$, where V_0, V_1, V_2 are polynomials. By Lemma 3.3, when $p_2 = 0$ and $a_2 = \cdots = a_n = 0$, $p_{2k} = 0$. Thus, for $j = 1,2$, $V_j(0,\dots,0) = 0$. Using the same method as above and the result about $p_2,\dots,p_{2n}$ we just proved, it is straightforward to show that p_{2k}, reduced modulo the ideal generated by $p_2,\dots,p_{2n}$, is zero. Hence, for any $k \ge 1$, $p_{2k} \in (p_2,\dots,p_{2n})$. □

Lemma 3.5 describes the simple structure of the period coefficients which enables us to prove the following theorem.

THEOREM 3.6. *If the origin is a weak center of type I_B, then it is an isochronous center of (3.1) if and only if $b_0 = (4/9)a_1^2$ and $a_2 = \cdots = a_n = 0$. If the origin is a weak center of finite order, then its order is at most $n-1$, at most $n-1$ local critical periods can bifurcate and there are perturbations with exactly j critical periods for each $j \le n-1$. Moreover, at most $n-1$ local critical periods can bifurcate from the isochronous center and there are perturbations to produce the maximum number of critical periods. Here, all perturbations of parameters are within S_I^B.*

PROOF. We first prove that the origin is an isochronous center if and only if $b_0 = (4/9)a_1^2$ and $a_2 = \cdots = a_n = 0$. The sufficiency of the condition has been proved in Lemma 3.3. Now, assume that the origin is an isochronous center. By Lemma 3.1, $p_2 = 0$ if and only if $b_0 = (4/9)a_1^2 > 0$. Perform coordinate transformation $(x,y) \mapsto ((2/3)a_1x,(2/3)a_1y)$, hence system (3.1) becomes

$$\dot{x} = -y + \sum_{k=1}^{n} \tilde{a}_k x^{2k}, \qquad \dot{y} = x + x^3, \tag{3.15}$$

where $\tilde{a}_k = (3/2a_1)^{2k-1}a_k$. System (3.15) has the same form as (3.4). Hence, by Lemma 3.5, $p_4 = \cdots = p_{2n} = 0$ if and only if $\tilde{a}_2 = \cdots = \tilde{a}_n = 0$. This implies that $a_2 = \cdots = a_n = 0$. Thus, the necessity of the condition is also proved.

To discuss the local critical periods, note that $b_0 > 0$, we can scale system (3.1) so that $b_0 = 1$. It suffices to consider (3.4) only.

If the origin is a weak center of finite order and $p_2 = 0$, then $a_1^2 = 9/4$ and there must be an integer k, such that $1 < k \le n$ and $a_2 = \cdots = a_{k-1} = 0$ and $a_k \ne 0$. Thus, by Lemma 3.5, $p_2 = \cdots = p_{2(k-1)} = 0$, $p_{2k} = a_1a_k \ne 0$ (with a nonzero constant factor omitted). That is, the origin is a weak center of order $k-1$. Hence, the maximum order of the weak center is $n-1$ and at most $n-1$ local critical periods can bifurcate.

Now, assume that the origin is a weak center of order $n-1$. Then, $a_1^2 = 9/4$, $a_2 = \cdots = a_{n-1} = 0$, and $a_n \ne 0$. Let $\nu_* = (a_1,0,\ldots,0,a_n)$, where $a_1 = 3/2$ or $-3/2$. Let a_1' be $3/2$ or $-3/2$. It is straightforward to see that the algebraic surfaces $p_{2(n-2)}(a_1', 0,\ldots,0,a_{n-2}',a_{n-1}',a_n') = 0$ and $p_{2(n-1)}(a_1',0,\ldots,0,a_{n-2}',a_{n-1}',a_n') = 0$ intersect transversally at their common roots for $(a_{n-2}',a_{n-1}') \in (-\infty,\infty)\times(-\infty,\infty)$. In fact, from the expressions for the period coefficients given in Lemma 3.5, the determinant of the Jacobian matrix of $p_{2(n-2)}(a_1',0,\ldots,0,a_{n-2}',a_{n-1}',a_n')$ and $p_{2(n-1)}(a_1',0,\ldots,0,a_{n-2}',a_{n-1}', a_n')$ is given by $a_1'^2$, which is not zero. This guarantees that in the neighborhood of ν_*, there exists a perturbation ν' such that $p_{2(n-2)}(\nu')p_{2(n-1)}(\nu') < 0$ with $p_{2n}(\nu') \ne 0$. This implies that, in the neighborhood of such perturbation, we may choose $\tilde{\nu}_*$ such that the system satisfies the condition $(\mathcal{P}_k)$ with $k = n-1$. Thus, by the finite-order bifurcation theorem, there are perturbations with exactly j critical periods for each $j \le n-1$.

Denote by $\tilde{\nu}_* = (\delta+\epsilon_1,\epsilon_2,\ldots,\epsilon_n)$ the perturbation of the isochronous center, where δ is $3/2$ or $-3/2$. Denote by $\tilde{p}_{2k}(\tilde{\nu}_*)$ the perturbed period coefficients. Using the same method as in the proof of Lemma 3.5, we find the perturbed period coefficients, with each reduced modulo the ideal generated by the previous coefficients, are given by $\tilde{p}_2(\tilde{\nu}_*) = (\pi/3)\epsilon_1(2\delta+\epsilon_1)$, $\tilde{p}_{2k}(\tilde{\nu}_*) = \epsilon_k(\delta+\epsilon_1)$ for $2 \le k \le n$, $\tilde{p}_{2k}(\tilde{\nu}_*) = 0$ for $k > n$.

Clearly, for any $k \geq 1$, $\tilde{p}_{2k}(\tilde{\nu}_*) \in (\tilde{p}_2(\tilde{\nu}_*), \ldots, \tilde{p}_{2n}(\tilde{\nu}_*))$, the ideal of the Noetherian ring $\mathbb{R}\{a_1, \ldots, a_n\}_{\nu_*}$ of convergent power series at $\nu_* = (\delta, 0, \ldots, 0)$. Thus, by the isochrone bifurcation theorem, at most $n-1$ local critical periods can bifurcate from the isochronous center. Similar to the argument in the previous paragraph, the maximum number can be reached. □

4. Weak center of type *II*. A weak center of type *II* corresponds to the following system:

$$\dot{x} = -y + \sum_{k=1}^{n} a_k (x^2 + b_1 x^3 + b_2 x^4)^k, \qquad \dot{y} = x + \frac{3}{2} b_1 x^2 + 2 b_2 x^3. \tag{4.1}$$

Executing our Maple program yields the following lemma.

Lemma 4.1. *The period coefficient p_2 for (4.1) is given by $p_2 = (\pi/24)(8a_1^2 + 45b_1^2 - 36b_2)$.*

We split S_{II} into 3 subsets as $S_{II} = S_{II}^A \bigcup S_{II}^B \bigcup S_{II}^C$, where

$$\begin{aligned} S_{II}^A &:= \{\lambda \in \mathbb{R}^{n+3} \mid \lambda_0 = 0,\ \lambda_2 \leq 0\}, \\ S_{II}^B &:= \{\lambda \in \mathbb{R}^{n+3} \mid \lambda_0 = \lambda_1 = 0,\ \lambda_2 > 0\}, \\ S_{II}^C &:= \{\lambda \in \mathbb{R}^{n+3} \mid \lambda_0 = 0,\ \lambda_1 \neq 0,\ \lambda_2 > 0\}. \end{aligned} \tag{4.2}$$

We say system (4.1) has a weak center of *type II_A* (resp., *type II_B* or *type II_C*) if the system is nonlinear and $\lambda \in S_{II}^A$ (resp., $\lambda \in S_{II}^B$ or $\lambda \in S_{II}^C$). We immediately have the following theorem.

Theorem 4.2. *A weak center of type II_A cannot be an isochronous center and no local critical periods can bifurcate from a weak center of type II_A.*

Proof. If $b_2 \leq 0$ and $b_1^2 + b_2^2 \neq 0$, then $p_2 > 0$ by Lemma 4.1 and the order of the weak center is 0. Thus, no local critical periods can bifurcate by finite order bifurcation theorem. If $b_1 = b_2 = 0$, then the weak center is of type I_A, there are no local critical periods that can bifurcate by Theorem 3.2. □

Thus, in the following we only need to discuss weak centers of types II_B and II_C.

4.1. Weak center of type II_B. For a weak center of type II_B, we can scale system (4.1) so that $b_2 = 1/2$. Thus, we only need to consider the following system:

$$\dot{x} = -y + \sum_{k=1}^{n} a_k \left(x^2 + \frac{x^4}{2}\right)^k, \qquad \dot{y} = x + x^3. \tag{4.3}$$

We first prove the following lemma.

Lemma 4.3. *A weak center of type II_B cannot be an isochronous center.*

PROOF. Clearly, system (4.3) can be written into the form

$$\dot{x} = -y + \sum_{k=1}^{n} c_k x^{2k}, \qquad \dot{y} = x + x^3, \tag{4.4}$$

where

$$c_k = \sum_{j=0}^{[k/2]} \frac{1}{2^j} \frac{(k-j)!}{j!(k-2j)!} a_{k-j}. \tag{4.5}$$

System (4.4) has the same form as (3.4). By Theorem 3.6, if the origin is an isochronous center, then $c_2 = \cdots = c_{2n} = 0$ and $c_1^2 = 9/4$. But $c_2 = \cdots = c_{2n} = 0$ implies that $a_1 = a_2 = \cdots = a_n = 0$. Hence, $c_1 = 0$. This is a contradiction. □

The following Lemma characterizes the ideal generated by the period coefficients.

LEMMA 4.4. *For system (4.3), the period coefficients p_{2k} $(k > 1)$, reduced modulo the ideal generated by $p_2, \ldots, p_{2k-2}$ and omitting the constant factor, are given by $p_{2k} = a_1 a_k + \alpha_k$ for $1 < k \le n$, $p_{2(n+1)} = \alpha_{n+1}$ and $p_{2k} = 0$ for $k > n+1$, where $\alpha_1, \ldots, \alpha_{n+1}$ are nonzero constants. In particular, $p_{2k} \in (p_2, \ldots, p_{2(n+1)})$, the ideal of the polynomial ring $\mathbb{R}[a_1, a_2, \ldots, a_n]$.*

PROOF. We first prove the result about p_{2k} for $1 < k \le n$ by induction on k. Direct computation yields $p_2 = (\pi/12)(4a_1^2 - 9)$, p_4 reduced modulo the ideal generated by p_2 and omitting the constant factor is $p_4 = a_1 a_2 + 9/8$.

Now, assume that $2 \le k \le n$ and $p_4, \ldots, p_{2(k-1)}$ with each reduced modulo the ideal generated by previous coefficients and omitting the constant factor, are given by $p_{2j} = a_1 a_j + \alpha_j$ $(j = 2, \ldots, k-1)$, where α_j is nonzero constant. Similar to the proof of Lemma 3.4, p_{2k} has the form $p_{2k} = \beta a_1 a_k + R(a_1, \ldots, a_{k-1})$, where β is a nonzero constant and $R(a_1, \ldots, a_{k-1})$ is a polynomial. Since $p_2 = (\pi/12)(4a_1^2 - 9)$, we can write $R(a_1, \ldots, a_{k-1}) = R_1(a_1, \ldots, a_{k-1}) p_2 + a_1 R_2(a_2, \ldots, a_{k-1}) + R_3(a_2, \ldots, a_{k-1})$, where R_1, R_2, R_3 are polynomials. We claim that $R_2(0, \ldots, 0) = 0$.

In fact, if we set $a_2 = \cdots = a_n = 0$, then system (4.3) has the form

$$\dot{x} = -y + a_1 \left(x^2 + \frac{x^4}{2} \right), \qquad \dot{y} = x + x^3. \tag{4.6}$$

Let $T(\xi, a_1)$ be the minimum period of the closed orbit of (4.6) passing through $(\xi, 0)$. Then, $T(\xi, -a_1)$ is the minimum period of the closed orbit of the following system passing through $(\xi, 0)$:

$$\dot{x} = -y - a_1 \left(x^2 + \frac{x^4}{2} \right), \qquad \dot{y} = x + x^3. \tag{4.7}$$

But (4.7) can be transformed to system (4.6) by the scaling $(x, y) \mapsto (-x, -y)$ and this scaling does not change the period of the closed orbit passing through $(\xi, 0)$. Thus, $T(\xi, a_1) = T(\xi, -a_1)$. Hence, when $a_2 = \cdots = a_n = 0$, the period coefficients are

functions of a_1^2. Thus, $R(a_1,0,\dots,0) = R_1(a_1,0,\dots,0)p_2 + a_1R_2(0,\dots,0) + R_3(0,\dots,0)$ is a polynomial of a_1^2. Since $p_2 = (\pi/12)(4a_1^2-9)$, it is clear that $R_1(a_1,0,\dots,0)$ does not have terms of a_1 with odd degree. Thus, $R_2(0,\dots,0) = 0$.

Hence, R_2 has the form

$$R_2(a_2,\dots,a_{k-1}) = a_2S_2(a_2,\dots,a_{k-1}) + \cdots + a_{k-1}S_{k-1}(a_2,\dots,a_{k-1}), \tag{4.8}$$

where $S_2,\dots,S_{k-1}$ are polynomials. Therefore, p_{2k} can be written as

$$p_{2k} = \beta a_1a_k + R_1p_2 + \sum_{i=2}^{k-1} a_1a_iS_i(a_2,\dots,a_{k-1}) + R_3(a_2,\dots,a_{k-1}). \tag{4.9}$$

By induction assumption, $a_1a_i = p_{2i} - \alpha_i$, so (4.9) can be written in the form

$$p_{2k} = \beta a_1a_k + R_1p_2 + \sum_{i=2}^{k-1} S_ip_{2i} + Q(a_2,\dots,a_{k-1}), \tag{4.10}$$

where $Q(a_2,\dots,a_{k-1})$ is a polynomial and can be written as $Q(a_2,\dots,a_{k-1}) = \kappa + Q_2(a_2,\dots,a_{k-1})$, where Q_2 is a polynomial such that $Q_2(0,\dots,0) = 0$ and κ is a constant.

On the other hand, $p_2 = (\pi/12)(4a_1^2-9)$, which implies that $\mu p_2 + \gamma a_1^2 = 1$, where $\mu = -4/3\pi$, $\gamma = 4/9$. Therefore, by induction hypothesis, for each a_j $(2 \le j \le k-1)$, we have

$$a_j = \mu a_jp_2 + \gamma a_1(a_1a_j) = \mu a_jp_2 + \gamma a_1p_{2j} + \tau_ja_1, \tag{4.11}$$

where $\tau_j = -\gamma\alpha_j$ is a constant. Thus, for any monomial $a_2^{j_2}\cdots a_{k-1}^{j_{k-1}}$ with $m := j_2 + \cdots + j_{k-1} > 0$, after reducing modulo the ideal generated by $p_2,\dots,p_{2(k-1)}$, it equals τa_1^m for some constant τ. Since $Q_2(a_2,\dots,a_{k-1})$ is a combination of such type of monomials, it is clear that, after reducing modulo the ideal generated by $p_2,\dots,p_{2(k-1)}$, $Q_2(a_2,\dots,a_{k-1}) = \tilde{Q}(a_1)$, a polynomial of a_1. Combining this with (4.10), it is obvious that, after reducing modulo the ideal generated by $p_2,\dots,p_{2(k-1)}$, $p_{2k} = \beta a_1a_k + \tilde{Q}(a_1) + \kappa$. Similar to the argument above, $\tilde{Q}(a_1)$ is a polynomial of a_1^2. Hence, it can be written in the form $\tilde{Q}(a_1) = \hat{Q}(a_1)p_2 + \iota$, where ι is a constant and $\hat{Q}(a_1)$ is a polynomial. Therefore, the reduced p_{2k} is $p_{2k} = \beta a_1a_k + \iota + \kappa = \beta(a_1a_k + \alpha_k)$, where $\alpha_k = (\iota+\kappa)/\beta$ is a constant. If we omit the constant factor, then $p_{2k} = a_1a_k + \alpha_k$. It is trivial to show that α_k is nonzero: just set $p_2 = 0$ and $a_2 = \cdots = a_n = 0$, simple computation shows that $p_4,\dots,p_{2n}$, which are multiples of $\alpha_2,\dots,\alpha_n$, are all nonzeros.

Now consider $p_{2(n+1)}$, it can be written as

$$p_{2(n+1)} = V_1(a_1,\dots,a_n)p_2 + a_1V_2(a_2,\dots,a_n) + V_3(a_2,\dots,a_n), \tag{4.12}$$

where V_1, V_2, V_3 are polynomials. Following the same reasoning as above, we can prove that the reduced $p_{2(n+1)}$ is a nonzero constant α_{n+1}. Since $p_{2(n+1)}$ reducing modulo the ideal generated by $p_2,\dots,p_{2n}$ is a nonzero constant, it is clear that whenever $k > n+1$, p_{2k} is reduced to zero. It follows that $p_{2k} \in (p_2,\dots,p_{2(n+1)})$ for any $k \geq 1$. The lemma is thus proved. □

From Lemma 4.4, we have the following theorem.

THEOREM 4.5. *If the origin is a weak center of type II_B, then it cannot be an isochronous center and it is a weak center of order at most n. At most n local critical periods can bifurcate from the weak center of order n and there are perturbations to produce exactly j critical periods for each $j \leq n$. Here all perturbations of parameters are within S_{II}^B.*

PROOF. For a weak center of type II_B, it suffices to consider system (4.3). We have proved in Lemma 4.3 that the origin cannot be an isochronous center. From Lemma 4.4, it is easy to see that the origin is a weak center of order at most n. Now, set $a_1 = 3/2$ or $a_1 = -3/2$, then $p_2 = 0$. If we set $a_k = -\alpha_k/a_1$ for $2 \leq k \leq n$, where α_k is the nonzero constant stated in Lemma 4.4, then we have $p_2 = \cdots = p_{2n} = 0$ and $p_{2(n+1)} \neq 0$. Thus, there exists a weak center of order n. By finite order bifurcation theorem, at most n critical periods can bifurcate from the weak center of order n.

Denote $\nu_* = (\delta, -\alpha_2/\delta,\dots,-\alpha_n/\delta)$, where $\delta = 3/2$ or $\delta = -3/2$. Then, $p_2(\nu_*) = \cdots = p_{2n}(\nu_*) = 0$ and $p_{2(n+1)}(\nu_*) \neq 0$. For any neighborhood $W \subset \mathbb{R}^n$ of ν_*, let $\nu' \in W$, $\nu' = (a_1',\dots,a_n')$, such that $p_2(\nu') = \cdots = p_{2n-2}(\nu') = 0$. Then, $p_{2n}(\nu') = a_1'a_n' + \alpha_n$ and $p_{2n+2}(\nu') = \alpha_{n+1}$. Obviously, we can always pick a_1' near δ and a_n' near $-\alpha_n/\delta$ such that $p_{2n}(\nu')p_{2n+2}(\nu') < 0$. Thus, the system satisfies condition $(\mathcal{P}_k)$ with $k = n$. by finite order bifurcation theorem, there are perturbations to produce exactly j critical periods for each $j \leq n$. □

4.2. Weak center of type II_C. For a weak center of type II_C, we again scale system (4.1) so that $b_2 = 1/2$. So, system (4.1) has the following form:

$$\dot{x} = -y + \sum_{k=1}^{n} a_k \left(x^2 + b_1x^3 + \frac{x^4}{2}\right)^k, \qquad \dot{y} = x + \frac{3}{2}b_1x^2 + x^3. \tag{4.13}$$

Clearly, $G(x) = \int_0^x g(\xi)d\xi = M/2$. In [6], Christopher and Devlin proved that if $F(x) = f_1M + \cdots + f_rM^r$ and $G(x) = g_1M + \cdots + g_pM^p$, then a necessary condition for the origin to be an isochronous center is that $p = 2r$. From this result, we immediately have the following lemma.

LEMMA 4.6. *A weak center of type II_C cannot be an isochronous center.*

It is much harder to describe the ideal generated by the period coefficients. Based on computation for $n \leq 6$, we believe that for any $k \geq 1$, the period coefficient $p_{2k} \in (p_2,\dots,p_{2(n+2)})$, the ideal of the polynomial ring $\mathbb{R}[b_1,a_1,\dots,a_n]$. Although we are unable to rigorously prove this, we are able to estimate the lower bound of the maximum number of local critical periods. First, we have the following lemma.

LEMMA 4.7. *For system (4.13), the period coefficients p_{2k} ($k \geq 1$) are polynomials of b_1^2.*

PROOF. Performing the coordinate change $(x,y) \mapsto (x, y - F(x))$, system (4.13) can be rewritten into the form

$$\begin{aligned} \dot{x} &= -y, \\ \dot{y} &= x + \frac{3}{2}b_1x^2 + x^3 + (2x + 3b_1x^2 + 2x^3)\sum_{k=1}^{n} ka_k\left(x^2 + b_1x^3 + \frac{x^4}{2}\right)^{k-1} y, \end{aligned} \tag{4.14}$$

where

$$F(x) = \sum_{k=1}^{n} a_k\left(x^2 + b_1x^3 + \frac{x^4}{2}\right)^k. \tag{4.15}$$

The coordinate change is nonsingular near the origin and it does not change the periods of the closed orbits near the origin.

Let $T(\xi, b_1, a_1, \ldots, a_n)$ be the minimum period of the closed orbit of (4.13) passing through $(\xi, 0)$. Then, $T(\xi, -b_1, a_1, \ldots, a_n)$ is the minimum period of the closed orbit of the following system passing through $(\xi, 0)$:

$$\begin{aligned} \dot{x} &= -y, \\ \dot{y} &= x - \frac{3}{2}b_1x^2 + x^3 + (2x - 3b_1x^2 + 2x^3)\sum_{k=1}^{n} ka_k\left(x^2 - b_1x^3 + \frac{x^4}{2}\right)^{k-1} y. \end{aligned} \tag{4.16}$$

But (4.16) can be transformed to (4.14) via $(x,y) \mapsto (-x,-y)$. Thus, $T(\xi, b_1, a_1, \ldots, a_n) = T(\xi, -b_1, a_1, \ldots, a_n)$. Hence, $T(\xi, b_1, a_1, \ldots, a_n)$ is a function of b_1^2; namely, the period coefficients are polynomials of b_1^2. □

Note that for $b_1 = 0$, system (4.13) is the same as system (4.3). We have the following theorem.

THEOREM 4.8. *The maximum number of local critical periods which can bifurcate from a weak center of type II_C is at least n.*

PROOF. Let $T(\xi, b_1, a_1, \ldots, a_n)$ be the minimum period of the closed orbit passing through $(\xi, 0)$. By Lemma 4.7,

$$T(\xi, b_1, a_1, \ldots, a_n) = T(\xi, 0, a_1, \ldots, a_n) + b_1^2\tilde{T}(\xi, b_1^2, a_1, \ldots, a_n), \tag{4.17}$$

where $\tilde{T}(\xi, b_1^2, a_1, \ldots, a_n)$ is an analytic function. Clearly, $T(\xi, 0, a_1, \ldots, a_n)$ is identical to the period function of system (4.3). Let

$$P(\xi, b_1, a_1, \ldots, a_n) = T(\xi, b_1, a_1, \ldots, a_n) - 2\pi, \tag{4.18}$$

then

$$\begin{aligned} P(\xi,b_1,a_1,\dots,a_n) &= P(\xi,0,a_1,\dots,a_n)+b_1^2\tilde{T}(\xi,b_1^2,a_1,\dots,a_n),\\ P_\xi(\xi,b_1,a_1,\dots,a_n) &= P_\xi(\xi,0,a_1,\dots,a_n)+b_1^2\tilde{T}_\xi(\xi,b_1^2,a_1,\dots,a_n). \end{aligned} \tag{4.19}$$

By Theorem 4.5, the function $\xi \mapsto P_\xi(\xi,0,a_1,\dots,a_n)$ can have at most n zeros near $\xi=0$ and there exists $a_1',\dots,a_n'$ such that $P_\xi(\xi,0,a_1',\dots,a_n')$ has n zeros near $\xi=0$. Furthermore, following exactly the same line as in the proof of the finite order bifurcation theorem given in [3] (i.e., [3, Theorem 2.1]), we may select $a_1',\dots,a_n'$ and construct the n zeros $\xi_1,\dots,\xi_n$ such that $0<\xi_1<\cdots<\xi_n<\eta$ for some η and on each pair of the open intervals (ξ_{i-1},ξ_i), (ξ_i,ξ_{i+1}) $(1\le i\le n)$, $P_\xi(\xi,0,a_1',\dots,a_n')$ has different signs. Here, we set $\xi_0=0$ and $\xi_{n+1}=\eta$. Thus, there exists $\xi_0',\dots,\xi_n'$ such that $0<\xi_0'<\cdots<\xi_n'<\eta$ and for any $1\le i\le n$, $P_\xi(\xi_{i-1}',0,a_1',\dots,a_n')P_\xi(\xi_i',0,a_1',\dots,a_n')<0$. Since $\tilde{T}_\xi(\xi,b_1^2,a_1',\dots,a_n')$ is continuous (actually analytic) on $[0,\eta]$, it is easy to see that for sufficiently small $|b_1|$, $P_\xi(\xi_{i-1}',b_1,a_1',\dots,a_n')P_\xi(\xi_i',b_1,a_1',\dots,a_n')<0$. Hence, there exists $\tilde{\xi}_i\in(\xi_{i-1}',\xi_i')$ such that $P_\xi(\tilde{\xi}_i,b_1,a_1',\dots,a_n')=0$. So, $P_\xi(\xi,b_1,a_1',\dots,a_n')$ has at least n zeros. □

By finite order bifurcation theorem, no more than q local critical periods can bifurcate from a weak center of order q. Hence, we immediately obtain the following corollary to Theorem 4.8.

Corollary 4.9. *There are weak centers of type II_C with order at least n.*

If our conjecture that $p_{2k}\in(p_2,\dots,p_{2(n+2)})$ is true, then since the origin cannot be an isochronous center, the origin is a weak center of order at most n and there are at most n local critical periods that can bifurcate from the weak center of order n. By Theorem 4.8, the maximum number of critical periods can be attained.

5. Proof of Theorem 1.5

Proof of Theorem 1.5. By Theorems 3.2, 3.6, 4.2, 4.5, and Lemma 4.6, system (1.4) has a nonlinear isochronous center at the origin if and only if $b_0=(4/9)a_0^2$ and $b_1=b_2=a_1=\cdots=a_n=0$ and $a_0\neq0$. If $a_0=0$, then the origin is a linear isochronous center. For a nonlinear isochronous center, $b_0\neq0$. Since we only consider small perturbation of parameter values, any perturbation of the parameters corresponding to the isochronous center remains in S_I^B. Thus, assertion (1) of Theorem 1.5 is true by Theorem 3.6.

Note that under small perturbation of parameter values, a weak center of type I_A (resp., I_B) is still a weak center of type I_A (resp., I_B). Thus, assertion (2) of Theorem 1.5 is true by Theorems 3.2 and 3.6. Assertion (3) of Theorem 1.5 is clear by Theorem 4.8. □

We remark that a weak center of type II can be perturbed to become a weak center of type I, but this will not increase the number of local critical periods by (2) of Theorem 1.5.

Acknowledgments. The author thanks the referees for their useful remarks and suggestions. This work is partially supported by Science Foundation for Young Scholars of Sichuan University, Grant no. 0020105505017.

References

[1] W. A. Albarakati, N. G. Lloyd, and J. M. Pearson, *Transformation to Liénard form*, Electron. J. Differential Equations **2000** (2000), no. 76, 1–11.

[2] N. N. Bautin, *On the number of limit cycles which appear with the variation of coefficients from an equilibrium position of focus or center type*, American Math. Soc. Translation **1954** (1954), no. 100, 1–19.

[3] C. Chicone and M. Jacobs, *Bifurcation of critical periods for plane vector fields*, Trans. Amer. Math. Soc. **312** (1989), no. 2, 433–486.

[4] S.-N. Chow and J. A. Sanders, *On the number of critical points of the period*, J. Differential Equations **64** (1986), no. 1, 51–66.

[5] C. Christopher, *An algebraic approach to the classification of centers in polynomial Liénard systems*, J. Math. Anal. Appl. **229** (1999), no. 1, 319–329.

[6] C. Christopher and J. Devlin, *On the classification of Liénard systems with amplitude-independent periods*, J. Differential Equations **200** (2004), no. 1, 1–17.

[7] C. Christopher and S. Lynch, *Small-amplitude limit cycle bifurcations for Liénard systems with quadratic or cubic damping or restoring forces*, Nonlinearity **12** (1999), no. 4, 1099–1112.

[8] P. Mardešić, C. Rousseau, and B. Toni, *Linearization of isochronous centers*, J. Differential Equations **121** (1995), no. 1, 67–108.

[9] V. G. Romanovski and M. Han, *Critical period bifurcations of a cubic system*, J. Phys. A **36** (2003), no. 18, 5011–5022.

[10] C. Rousseau and B. Toni, *Local bifurcation of critical periods in vector fields with homogeneous nonlinearities of the third degree*, Canad. Math. Bull. **36** (1993), no. 4, 473–484.

[11] ———, *Local bifurcations of critical periods in the reduced Kukles system*, Canad. J. Math. **49** (1997), no. 2, 338–358.

[12] M. Sabatini, *On the period function of Liénard systems*, J. Differential Equations **152** (1999), no. 2, 467–487.

[13] B. Toni, *Small amplitude critical periods in a cubic polynomial system*, Southwest J. Pure Appl. Math. (1999), no. 2, 1–12.

[14] W. Zhang, X. Hou, and Z. Zeng, *Weak centers and bifurcation of critical periods in reversible cubic systems*, Comput. Math. Appl. **40** (2000), no. 6-7, 771–782.

[15] Z. F. Zhang, T. R. Ding, W. Z. Huang, and Z. X. Dong, *Qualitative Theory of Differential Equations*, Translations of Mathematical Monographs, vol. 101, American Mathematical Society, Rhode Island, 1992.

Zhengdong Du: Department of Mathematics, Sichuan University, Chengdu, Sichuan 610064, China

E-mail address: zdu85@yahoo.com

Journal of Applied Mathematics and Decision Sciences

Special Issue on
Intelligent Computational Methods for Financial Engineering

Call for Papers

As a multidisciplinary field, financial engineering is becoming increasingly important in today's economic and financial world, especially in areas such as portfolio management, asset valuation and prediction, fraud detection, and credit risk management. For example, in a credit risk context, the recently approved Basel II guidelines advise financial institutions to build comprehensible credit risk models in order to optimize their capital allocation policy. Computational methods are being intensively studied and applied to improve the quality of the financial decisions that need to be made. Until now, computational methods and models are central to the analysis of economic and financial decisions.

However, more and more researchers have found that the financial environment is not ruled by mathematical distributions or statistical models. In such situations, some attempts have also been made to develop financial engineering models using intelligent computing approaches. For example, an artificial neural network (ANN) is a nonparametric estimation technique which does not make any distributional assumptions regarding the underlying asset. Instead, ANN approach develops a model using sets of unknown parameters and lets the optimization routine seek the best fitting parameters to obtain the desired results. The main aim of this special issue is not to merely illustrate the superior performance of a new intelligent computational method, but also to demonstrate how it can be used effectively in a financial engineering environment to improve and facilitate financial decision making. In this sense, the submissions should especially address how the results of estimated computational models (e.g., ANN, support vector machines, evolutionary algorithm, and fuzzy models) can be used to develop intelligent, easy-to-use, and/or comprehensible computational systems (e.g., decision support systems, agent-based system, and web-based systems)

This special issue will include (but not be limited to) the following topics:

- **Computational methods**: artificial intelligence, neural networks, evolutionary algorithms, fuzzy inference, hybrid learning, ensemble learning, cooperative learning, multiagent learning
- **Application fields**: asset valuation and prediction, asset allocation and portfolio selection, bankruptcy prediction, fraud detection, credit risk management
- **Implementation aspects**: decision support systems, expert systems, information systems, intelligent agents, web service, monitoring, deployment, implementation

Authors should follow the Journal of Applied Mathematics and Decision Sciences manuscript format described at the journal site http://www.hindawi.com/journals/jamds/. Prospective authors should submit an electronic copy of their complete manuscript through the journal Manuscript Tracking System at http://mts.hindawi.com/, according to the following timetable:

Manuscript Due	December 1, 2008
First Round of Reviews	March 1, 2009
Publication Date	June 1, 2009

Guest Editors

Lean Yu, Academy of Mathematics and Systems Science, Chinese Academy of Sciences, Beijing 100190, China; Department of Management Sciences, City University of Hong Kong, Tat Chee Avenue, Kowloon, Hong Kong; yulean@amss.ac.cn

Shouyang Wang, Academy of Mathematics and Systems Science, Chinese Academy of Sciences, Beijing 100190, China; sywang@amss.ac.cn

K. K. Lai, Department of Management Sciences, City University of Hong Kong, Tat Chee Avenue, Kowloon, Hong Kong; mskklai@cityu.edu.hk