

**OPERATIONAL CALCULUS FOR THE CONTINUOUS
LEGENDRE TRANSFORM WITH APPLICATIONS**

E.Y. DEEBA

Department of Applied Mathematical Sciences
University of Houston-Downtown
Houston, Texas 77002

E.L. KOH

Department of Mathematics and Statistics
University of Regina
Regina, Saskatchewan, Canada, S4S 0A2

(Received June 13, 1988 and in revised form September , 1988)

ABSTRACT. This paper develops an operational calculus for the continuous Legendre transform introduced and studied by Butzer, Stens and Wehrens [1]. It is an extension of the work done by Churchill et al [2], [3] for the discrete case. In particular, a differentiation theorem and a convolution theorem are proved and the results are applied to the solution of some boundary value problems.

KEY WORDS AND PHRASES. *Continuous Legendre Transform, Operational Calculus, Convolution, Boundary Value Problems.*

1980 AMS SUBJECT CLASSIFICATION CODES: 44A15, 93A90, 44A20.

1. **INTRODUCTION.** For a given function f belonging to an appropriate function space, the continuous Legendre transform is defined by

$$(Tf)(\lambda) = \frac{1}{2} \int_{-1}^1 P_\lambda(x) f(x) dx \quad (1)$$

where $P_\lambda(x)$ is the Legendre function and $\lambda \geq -\frac{1}{2}$. This transform has been introduced and studied by Butzer, Stens and Wehrens [1]. The discrete analog of the transform in (1) has been studied by Churchill [2] and Churchill and Dolph [3]. The object of this paper is to develop an operational calculus for the transform which is useful in solving partial differential equations whose underlying differential form is given by

$$D = \frac{d}{dx} \left[(1 - x^2) \frac{d}{dx} \right]. \quad (2)$$

In section 2 we present the background material needed in the sequel. In section 3, we derive the operational calculus for (1) including a convolution theorem and a table of transforms of some functions. In the last section we apply the results to solving some boundary value problems.

2. PRELIMINARIES. We recall basic properties of the transform $(Tf)(\lambda)$ (see [1]) and important contiguous relations that hold for the Legendre function.

The Legendre function $P_\lambda(x)$ is given by

$$P_\lambda(x) = {}_2F_1(-\lambda, \lambda + 1; 1; \frac{1-x}{2}) = \sum_{k=0}^{\infty} \frac{(-\lambda)_k (\lambda + 1)_k}{(k!)^2} \left(\frac{1-x}{2}\right)^k, \quad x \in (-1, 1].$$

Since $P_{\lambda-1}(x) = P_{-\lambda}(x)$, it suffices to consider the case $\lambda \geq -\frac{1}{2}$. $P_\lambda(x)$ satisfies the differential equation

$$Dy + \lambda(\lambda + 1)y = 0$$

where D is as given in (2). Further, it satisfies the relations $P_\lambda(1) = 1$, $P'_\lambda(1) = \frac{\lambda(\lambda+1)}{2}$, $\lim_{x \rightarrow -1^+} (1+x)P_\lambda(x) = 0$ and $\lim_{x \rightarrow -1^+} (1+x)P'_\lambda(x) = \frac{\sin \pi \lambda}{\pi}$.

The following contiguous relations (see [4]) will be useful in the derivation of the calculus for $(Tf)(\lambda)$:

$$(2\lambda + 1)xP_\lambda(x) = (\lambda + 1)P_{\lambda+1}(x) + \lambda P_{\lambda-1}(x) \quad (3)$$

and

$$(1 - x^2)P'_\lambda(x) = -\lambda xP_\lambda(x) + \lambda P_{\lambda-1}(x). \quad (4)$$

From (3) and (4) we obtain the relation

$$(1 - x^2)P'_\lambda(x) = -\frac{\lambda(\lambda + 1)}{2\lambda + 1} (P_{\lambda+1}(x) - P_{\lambda-1}(x)). \quad (5)$$

The addition formula for the Legendre functions (see [4]) is given by

$$P_\lambda(\cos \alpha)P_\lambda(\cos \beta) = P_\lambda(\cos \nu) - 2 \sum_{m=1}^{\infty} \frac{\Gamma(\lambda - m + 1)}{\Gamma(\lambda + m + 1)} P_\lambda^m(\cos \alpha)P_\lambda^m(\cos \beta) \cos m\gamma \quad (6)$$

where $P_\lambda^m(\cdot)$ is the associated Legendre function and $\cos \nu = \cos \alpha \cos \beta + \sin \alpha \sin \beta \cos \gamma$ with $0 \leq \alpha, \beta \leq \pi$, $\alpha + \beta < \pi$, γ real. Formula (6) will be useful in deriving the convolution theorem. Another useful relation involving the Legendre functions is

$$\int_{-1}^1 P_\lambda(x)P_\nu(-x)dx = \frac{\sin \pi \lambda - \sin \pi \nu}{\pi(\lambda - \nu)(\lambda + \nu + 1)}, \quad \lambda \neq \nu, \lambda + \nu + 1 \neq 0. \quad (7)$$

The Legendre transform $(Tf)(\lambda)$ is a linear integral transform from $L_2(-1, 1]$ into the space $C_0(-1, 1] \cap L_2(-1, 1]$. For $f \in L_2(-1, 1]$, it was shown in [1] that $(Tf)(\lambda) = 0(\lambda^{-\frac{1}{2}})$ as $\lambda \rightarrow \infty$ and $(Tf)(\lambda - \frac{1}{2}) \in C_0(-1, 1] \cap L_2(-1, 1]$. Further, it was shown that if $f \in L_2(-1, 1] \cap C(-1, 1]$ and if $\sqrt{\lambda}(Tf)(\lambda - \frac{1}{2}) \in L_1(\mathbb{R}^+)$, then the inversion formula is given by

$$f(x) = T^{-1}((Tf)(\lambda)) = 4 \int_0^{\infty} (Tf)(\lambda - \frac{1}{2})P_{\lambda-\frac{1}{2}}(-x)\lambda \sin \pi \lambda d\lambda. \quad (8)$$

3. BASIC OPERATIONAL PROPERTIES FOR $(Tf)(\lambda)$. In this section we shall derive operational calculus for the continuous Legendre transform $(Tf)(\lambda)$ thus extending the calculus obtained by Churchill [2] and Churchill and Dolph [3] for the discrete case. We shall also derive the Legendre transform of some functions.

The first property in this direction involves the Legendre transform of the differential operator D as given in (2).

Theorem 3.1. Let f be a function such that (i) $f^{(k)} \in C(-1, 1] \cap L_2(-1, 1]$ $k = 0, 1$ (ii) $\lim_{x \rightarrow \pm 1} (1 - x^2)f(x) = \lim_{x \rightarrow \pm 1} (1 - x^2)f'(x) = 0$ and (iii) $(Tf)(\lambda)$ exists. Then

$$(T(Df))(\lambda) = -\lambda(\lambda + 1)(Tf)(\lambda). \quad (9)$$

Proof. From (1) together with successive integration by parts, we obtain

$$\begin{aligned} (T(Df))(\lambda) &= \frac{1}{2} \int_{-1}^1 P_\lambda(x) Df(x) dx \\ &= \frac{1}{2} \int_{-1}^1 P_\lambda(x) \frac{d}{dx} \left[(1 - x^2) \frac{d}{dx} f(x) \right] dx \\ &= \left[\frac{1}{2} P_\lambda(x) (1 - x^2) f'(x) - \frac{1}{2} P'_\lambda(x) (1 - x^2) f(x) \right]_{-1+}^{+1-} \\ &- \frac{1}{2} \lambda(\lambda + 1) \int_{-1}^1 P_\lambda(x) f(x) dx. \end{aligned}$$

The result follows from the facts that $P_\lambda(1) = 1$, $P'_\lambda(1) = \frac{\lambda(\lambda+1)}{2}$, $\lim_{x \rightarrow -1+} (1 + x)P_\lambda(x) = 0$ and $\lim_{x \rightarrow -1+} (1 + x)P'_\lambda(x) = \frac{\sin \pi \lambda}{\pi}$ together with the hypothesis (ii).

This basic operational property reduces a given differential equation which involves the operator D into an algebraic one or into a differential equation with one less independent variable.

Remark 3.1. (a) If, in Theorem 3.1, $D^k f = D^{k-1}(Df)$ and $f^{(k)}$ satisfy the same hypotheses, then

$$T((D^k f)(x))(\lambda) = (-1)^k \lambda^k (\lambda + 1)^\lambda (Tf)(\lambda), \quad k = 1, 2, \dots.$$

(b) We note that (9) can be cast into the form

$$\frac{1}{4}(Tf)(\lambda) - T((Df))(\lambda) = (\lambda + \frac{1}{2})^2 (Tf)(\lambda). \quad (10)$$

The second operational property involves the relationship between the transform of a given function f and the function $g(x) = \int_{-1}^x f(t) dt$.

Theorem 3.2. If f is a piecewise continuous function defined on $(-1, 1)$ and $g(x) = \int_{-1}^x f(t) dt$ and if $(Tf)(\lambda)$ exists, then

$$(Tg)(\lambda) = -\frac{(Tf)(\lambda + 1) - (Tf)(\lambda - 1)}{2\lambda + 1}. \quad (11)$$

Proof. Since $D(P_\lambda(x)) = -\lambda(\lambda + 1)P_\lambda(x)$, it follows that

$$\begin{aligned} (Tg)(\lambda) &= -\frac{1}{2\lambda(\lambda + 1)} \int_{-1}^1 \frac{d}{dx} \left[(1 - x^2) \frac{d}{dx} P_\lambda(x) \right] g(x) dx \\ &= -\frac{1}{2\lambda(\lambda + 1)} (1 - x^2) P'_\lambda(x) g(x) \Big|_{-1}^1 + \frac{1}{2\lambda(\lambda + 1)} \int_{-1}^1 (1 - x^2) P'_\lambda(x) f(x) dx. \end{aligned}$$

Since $P'_\lambda(1)$ and $g(1)$ are defined, $g(-1) = 0$ and $\lim_{x \rightarrow -1+} (1 + x)P'_\lambda(x) = \frac{\sin \pi \lambda}{\pi}$, the first term is identically zero. Thus

$$(Tg)(\lambda) = \frac{1}{2\lambda(\lambda + 1)} \int_{-1}^1 (1 - x^2) P'_\lambda(x) f(x) dx.$$

The contiguous relation (5) will then imply that

$$(Tg)(\lambda) = \frac{1}{2\lambda(\lambda + 1)} \int_{-1}^1 -\frac{\lambda(\lambda + 1)}{2\lambda + 1} (P_{\lambda+1}(x) - P_{\lambda-1}(x)) f(x) dx$$

Equivalently,

$$(Tg)(\lambda) = -\frac{(Tf)(\lambda + 1) - (Tf)(\lambda - 1)}{2\lambda + 1}.$$

Remark 3.2. Similar difference relations to that of (11) can be obtained in the following situation.

(a) If $g(x) = xf(x)$ and if $(Tf)(\lambda)$ exists, then under appropriate conditions on f , one obtains

$$(Tg)(\lambda) = \frac{(\lambda + 1)(Tf)(\lambda + 1) + \lambda(Tf)(\lambda - 1)}{2\lambda + 1} \quad (12)$$

This will follow by applying the contiguous relation (3).

(b) If $g(x) = \int_{-1}^x (x-t)f(t)dt$ and if $(Tf)(\lambda)$ exists, then, again under appropriate conditions on f , the contiguous relation (5) and Theorem 3.2 yields

$$(Tg)(\lambda) = \frac{(Tf)(\lambda + 2) - 2(Tf)(\lambda) + (Tf)(\lambda - 2)}{(2\lambda + 1)^2} \quad (13)$$

The next operational property that we will derive involves the inverse of the differential operator D . We define the inverse of D , denoted by D^{-1} , by $D^{-1}(f(x)) = g(x)$ if and only if $D(g(x)) = f(x)$. If $(Tf)(\lambda)$ is known, then we want to relate $T((D^{-1}f))(\lambda)$ to the transform of f .

If, for a given function $f(x)$, $D(g(x)) = f(x)$, then on integrating twice, we obtain

$$g(x) = \int_0^x \frac{1}{1-t^2} \int_{-1}^t f(\alpha) d\alpha dt + c$$

for some constant c . If $f(x)$ is in addition an even function on $(-1, 1)$, then one can show by employing a continuity argument that $\lim_{x \rightarrow \pm 1} (1-x^2)g(x) = \lim_{x \rightarrow \pm 1} (1-x^2)g'(x) = 0$. Theorem 3.1 will then imply that

$$(Tf)(\lambda) = T((Dg))(\lambda) = -\lambda(\lambda + 1)(Tg)(\lambda).$$

Equivalently,

$$(Tg)(\lambda) = -\frac{1}{\lambda(\lambda + 1)} T((Dg))(\lambda) = -\frac{(Tf)(\lambda)}{\lambda(\lambda + 1)}.$$

Thus

$$T(D^{-1}f)(\lambda) = -\frac{1}{\lambda(\lambda + 1)}(Tf)(\lambda).$$

This last relation implies that $D^{-1}f$ is the inverse Legendre transform of $-\frac{(Tf)(\lambda)}{\lambda(\lambda + 1)}$. We thus have

Theorem 3.3. If $f(x)$ is such that $f(x)$ is even on $(-1, 1)$, $f \in L_2(-1, 1] \cap C(-1, 1]$, $(Tf)(\lambda)$ exists and $\frac{(Tf)(\lambda)}{\sqrt{\lambda(\lambda + 1)}} \in L_1(\mathbb{R}^+)$, then

$$D^{-1}(f(x)) = T^{-1} \left(-\frac{(Tf)(\lambda)}{\lambda(\lambda + 1)} \right) \quad (14)$$

where the inverse transform T^{-1} is given by (8).

We shall finally develop a convolution property for the Legendre transform. In particular, we will show

Theorem 3.4. If $f(x)$ and $g(x)$ are given functions for which $(Tf)(\lambda)$ and $(Tg)(\lambda)$ respectively exist, then their product $(Tf)(\lambda)(Tg)(\lambda)$ is the transform of the function $h(x) = f(x)*g(x)$ where $h(x)$ is given by

$$h(\cos \nu) = \frac{1}{2\pi} \int_0^\pi \int_0^\pi f(\cos \alpha)g(\cos \beta) \sin \alpha d\alpha d\theta$$

where $\cos \beta = \cos \alpha \cos \nu + \sin \alpha \sin \nu \cos \theta$ with $0 \leq \alpha, \nu \leq \pi$, $\alpha + \nu < \pi$ and θ is real. The variables α , ν and β may be interpreted as the sides of a spherical triangle on the unit hemisphere and θ is the angle between the sides α and ν (see Figure 1).

Proof. From (1), we have

$$(Tf)(\lambda)(Tg)(\lambda) = \frac{1}{4} \int_{-1}^1 P_\lambda(x) f(x) dx \int_{-1}^1 P_\lambda(y) g(y) dy.$$

Set $x = \cos \alpha$ and $y = \cos \beta$. Then

$$(Tf)(\lambda)(Tg)(\lambda) = \frac{1}{4} \int_0^\pi f(\cos \alpha) \sin \alpha \int_0^\pi P_\lambda(\cos \alpha) P_\lambda(\cos \beta) g(\cos \beta) \sin \beta d\beta d\alpha.$$

The addition formula for the Legendre function (6) will yield upon an integration with respect to γ from 0 to π

$$P_\lambda(\cos \alpha) P_\lambda(\cos \beta) = \frac{1}{\pi} \int_0^\pi P_\lambda(\cos \nu) d\gamma$$

where $\cos \nu = \cos \alpha \cos \beta + \sin \alpha \sin \beta \cos \gamma$ (see figure 1).

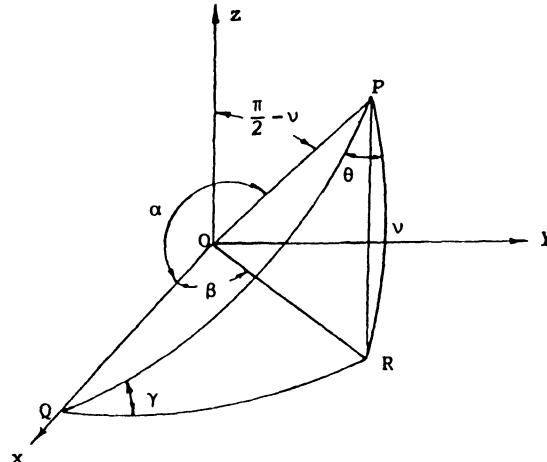


Figure 1

Thus

$$(Tf)(\lambda)(Tg)(\lambda) = \frac{1}{4\pi} \int_0^\pi f(\cos \alpha) \sin \alpha \int_0^\pi \int_0^\pi P_\lambda(\cos \nu) g(\cos \beta) \sin \beta d\gamma d\beta d\alpha.$$

In the spherical triangle PQR , we have

$$\cos \beta = \cos \alpha \cos \nu + \sin \alpha \sin \nu \cos \theta.$$

Using this relation along with the sine law and transformation of co-ordinates, the double integral can be written as:

$$\int_0^\pi \int_0^\pi P_\lambda(\cos \nu) g(\cos \beta) \sin \nu d\theta d\nu.$$

Hence,

$$(Tf)(\lambda)(Tg)(\lambda) = \frac{1}{2} \int_0^\pi P_\lambda(\cos \nu) \sin \nu \left[\frac{1}{2\pi} \int_0^\pi \int_0^\pi f(\cos \alpha) g(\cos \beta) \sin \alpha d\alpha d\theta \right] d\nu.$$

The expression in the bracket is a function of ν and we then write

$$h(\cos \nu) = \frac{1}{2\pi} \int_0^\pi \int_0^\pi f(\cos \alpha) g(\cos \beta) \sin \alpha d\alpha d\theta \quad (15)$$

This may be interpreted as a convolution product of f and g and $(Th(\cos \nu))(\lambda) = (Tf)(\lambda)(Tg)(\lambda)$.

This proves Theorem 3.4.

Geometrically, the expression (15) is the mean value of $f(\cos \alpha)g(\cos \beta)$ over the unit hemisphere $x^2 + y^2 + z^2 = 1, z \geq 0$. To see this, we note that the element surface area is $dS = \sin \alpha d\alpha d\theta$.

This is clear if we identify the coordinate transformation in Figure 1 by

$$\begin{aligned} x &= \cos \alpha \\ y &= \sin \alpha \sin \theta \\ z &= \sin \alpha \cos \theta \end{aligned}$$

Thus (15) reads

$$h(\cos \nu) = \frac{1}{2\pi} \int_S \int f(\cos \alpha) g(\cos \beta) dS.$$

We will now evaluate the Legendre transform of some functions.

1. $f(x) = \text{constant} = k$

$$(Tf)(\lambda) = \begin{cases} k \frac{\sin \pi \lambda}{\pi \lambda (\lambda + 1)} & \lambda \neq 0 \\ k & \lambda = 0 \end{cases}$$

2. $f(x) = P_n(x)$. Then by (2.5) we have, for $n = 0, 1, 2, \dots$,

$$\begin{aligned} (Tf)(\lambda) &= \frac{1}{2} \int_{-1}^1 P_\lambda(x) P_n(x) = \frac{(-1)^n}{2} \int_{-1}^1 P_\lambda(x) P_n(-x) dx \\ &= \frac{\sin \pi(\lambda - n)}{2\pi(\lambda - n)(\lambda + n + 1)}, \quad \lambda \neq n, -(n + 1). \end{aligned}$$

3. $f(x) = \log(1 - x)$.

$$\begin{aligned} (Tf)(\lambda) &= \frac{1}{2} \int_{-1}^1 P_\lambda(x) \log(1 - x) dx \\ &= -\frac{1}{2\lambda(\lambda + 1)} \int_{-1}^1 \frac{d}{dx} \left[(1 - x^2) \frac{d}{dx} P_\lambda(x) \right] \log(1 - x) dx \\ &= (\log 2) \frac{\sin \pi \lambda}{\lambda(\lambda + 1)} - \frac{1}{\lambda(\lambda + 1)} - \frac{1}{2\lambda(\lambda + 1)} \int_{-1}^1 P_\lambda(x) \frac{d}{dx} \left[(1 - x^2) \frac{d}{dx} \log(1 - x) \right] dx. \end{aligned}$$

Observe that $D(\log(1 - x)) = \frac{d}{dx} \left[(1 - x^2) \frac{d}{dx} \log(1 - x) \right] = -1$. Thus

$$(Tf)(\lambda) = (\log 2) \frac{\sin \pi \lambda}{\lambda(\lambda + 1)} - \frac{1}{\lambda(\lambda + 1)} - \frac{\sin \pi \lambda}{\lambda^2(\lambda + 1)^2}$$

4. $f(\lambda) = \int_{-1}^x \frac{1}{1-t} dt$. By using 1 and 3 above, we obtain

$$(Tf)(\lambda) = \frac{1}{\lambda(\lambda+1)} + \frac{\sin \pi \lambda}{\lambda^2(\lambda+1)^2}.$$

5. $f(x) = (1 - 2tx + x^2)^{-\frac{1}{2}} = \sum_{n=0}^{\infty} t^n P_n(x)$, $-1 < t < 1$. From (2) above

$$(Tf)(\lambda) = \frac{\sin \pi \lambda}{2\pi} \sum_{n=0}^{\infty} \frac{t^n}{(\lambda - n)(\lambda + n + 1)}.$$

We finally remark that for λ equal to a non-negative integer, the results of this section yield those obtained in [2] and [3].

4. **APPLICATIONS.** In this section we consider some applications of the Legendre transform. We consider problems arising in heat conduction and in potential theory.

A. **Heat Conduction Problem.** Consider a non-homogeneous bar with extremities at $x = \pm 1$ and is insulated at these end points. Let $u(x, t)$ be the temperature of the bar at position x at time t . The one dimensional heat equation with prescribed initial temperature is given by

$$\begin{aligned} \frac{\partial}{\partial x} \left(k \frac{\partial}{\partial x} u(x, t) \right) &= \rho c \frac{\partial u}{\partial t}(x, t) \\ u(x, 0) &= g(x), \quad -1 < x < 1 \end{aligned}$$

where k , ρ and c are physical constants representing thermal conductivity, density and specific heat respectively. We assume that the thermal conductivity k is given by $k = \alpha(1 - x^2)$, α being a real constant. The above equation reads

$$\begin{aligned} \frac{\partial}{\partial x} \left((1 - x^2) \frac{\partial}{\partial x} u(x, t) \right) &= \frac{\rho c}{\alpha} \frac{\partial u}{\partial t}(x, t) \\ u(x, 0) &= g(x), \quad -1 < x < 1. \end{aligned}$$

If $U(\lambda, t) = T(u(x, t))(\lambda)$ and $G(\lambda) = (Tu(x, 0))(\lambda)$, then, by Theorem 3.1, we obtain upon the application of the transform

$$\begin{aligned} \frac{d}{dt} U(\lambda, t) &= -\frac{\alpha}{\rho c} \lambda(\lambda+1)U(\lambda, t) \\ U(\lambda, 0) &= G(\lambda). \end{aligned}$$

The solution is given by

$$U(\lambda, t) = G(\lambda) e^{-\frac{\alpha}{\rho c}(\lambda+1)\lambda t}$$

Now $u(x, t)$ can be obtained by either employing the inversion formula (8) or the convolution theorem. By employing the inversion formula and under the assumption that $u(x, t) \in C(-1, 1) \cap L_2(-1, 1)$ and $\sqrt{\lambda} U(\lambda - \frac{1}{2}, t) \in L_1(\mathbb{R}^+)$, one obtains

$$u(x, t) = 4 \int_0^{\infty} G(\lambda - \frac{1}{2}) e^{-\frac{\alpha}{\rho c}(\lambda^2 - \frac{1}{4})t} P_{\lambda}(-x) \lambda \sin \pi \lambda d\lambda$$

On the other hand the convolution property (Theorem 3.4) will yield

$$u(\cos \nu, t) = \frac{1}{2\pi} \int_0^{\pi} \int_0^{\pi} g(\cos \alpha) f(\cos \beta) \sin \alpha d\alpha d\theta$$

where α , β , θ are as in Figure 1 and $\cos \beta = \cos \alpha \cos \nu + \sin \alpha \sin \nu \cos \theta$ and f is the inverse transform of $e^{-\frac{\alpha}{\rho c}(\lambda^2 - \frac{1}{4})t}$. That is, by (8)

$$f(x) = e^{\frac{\alpha}{\rho c} t} 4 \int_0^{\infty} e^{-\frac{\alpha}{\rho c}(\lambda - \frac{1}{2})^2} P_{\lambda}(-x) \lambda \sin \pi \lambda d\lambda.$$

B. Dirichlet Problem for the Unit Sphere (see[2]) Consider the problem of determining the potential $v(r, \cos \theta)$ in the interior of a unit sphere with a prescribed potential $f(\cos \theta)$ on $r = 1$, $0 \leq \theta \leq \pi$. The Laplace equation defining this potential is

$$\nabla^2 v = \frac{1}{r}(rv)_{rr} + \frac{1}{r^2 \sin \theta}(\sin \theta v_\theta)_\theta = 0.$$

If $x = \cos \theta$, then the equation reduces to

$$\begin{aligned} r(rv)_{rr} + ((1-x^2)v_x)_x &= 0 \\ v(1, x) &= f(x), \quad -1 \leq x \leq 1. \end{aligned}$$

If $V(r, \lambda)$ and $F(\lambda)$ denote respectively the Legendre transform of $v(r, x)$ and $f(x)$, then, upon applying the transform to the underlying equation, we obtain

$$\begin{aligned} r \frac{d^2}{dr^2}(rV(r, \lambda)) - \lambda(\lambda + 1)V(r, \lambda) &= 0, \\ V(1, \lambda) &= F(\lambda). \end{aligned}$$

The solution of this equation is given by

$$V(r, \lambda) = c_1 r^\lambda + c_2 r^{-(\lambda+1)}.$$

In order to apply the inversion formula (8) we need to have $v(r, x) \in L_2(-1, 1) \cap C(-1, 1)$ and $\sqrt{\lambda} V(r, \lambda) \in L_1(\mathbb{R}^+)$. This will imply that $c_2 = 0$ and $v(1, \lambda) = F(\lambda)$ will imply that $c_1 = F(\lambda)$. Hence the solution is given by

$$v(r, \lambda) = F(\lambda)r^\lambda$$

and

$$v(r, x) = r \int_0^\infty F(\lambda - \frac{1}{2}) r^{\lambda - \frac{1}{2}} P_\lambda(-x) \lambda \sin \pi \lambda d\lambda.$$

ACKNOWLEDGEMENT. The work of the first author was partially supported by a research grant from the University of Houston-Downtown. The second author was partially supported by NSERC of Canada under Grant A-7184.

REFERENCES

1. Butzer, P.L., R.L. Stens and M. Wehrens, The continuous Legendre transform, its inverse transform and applications, Internat. J. Math and Math. Sciences, Vol. 3, No. 1 (1980), 47-67.
2. Churchill, R.V., The operational calculus of the Legendre transform, J. Math. & Phys., 33(1954), 165-177.
3. Churchill, R.V. and C.L. Dolph, Inverse transforms of product of Legendre transforms, Proc. Amer. Math. Soc., 5(1954), 93-100.
4. Erdelyi, A., et al, Higher Transcendental Functions, Vol. 1, McGraw Hill, N.Y., 1953.

Special Issue on Time-Dependent Billiards

Call for Papers

This subject has been extensively studied in the past years for one-, two-, and three-dimensional space. Additionally, such dynamical systems can exhibit a very important and still unexplained phenomenon, called as the Fermi acceleration phenomenon. Basically, the phenomenon of Fermi acceleration (FA) is a process in which a classical particle can acquire unbounded energy from collisions with a heavy moving wall. This phenomenon was originally proposed by Enrico Fermi in 1949 as a possible explanation of the origin of the large energies of the cosmic particles. His original model was then modified and considered under different approaches and using many versions. Moreover, applications of FA have been of a large broad interest in many different fields of science including plasma physics, astrophysics, atomic physics, optics, and time-dependent billiard problems and they are useful for controlling chaos in Engineering and dynamical systems exhibiting chaos (both conservative and dissipative chaos).

We intend to publish in this special issue papers reporting research on time-dependent billiards. The topic includes both conservative and dissipative dynamics. Papers discussing dynamical properties, statistical and mathematical results, stability investigation of the phase space structure, the phenomenon of Fermi acceleration, conditions for having suppression of Fermi acceleration, and computational and numerical methods for exploring these structures and applications are welcome.

To be acceptable for publication in the special issue of Mathematical Problems in Engineering, papers must make significant, original, and correct contributions to one or more of the topics above mentioned. Mathematical papers regarding the topics above are also welcome.

Authors should follow the Mathematical Problems in Engineering manuscript format described at <http://www.hindawi.com/journals/mpe/>. Prospective authors should submit an electronic copy of their complete manuscript through the journal Manuscript Tracking System at <http://mts.hindawi.com/> according to the following timetable:

Manuscript Due	December 1, 2008
First Round of Reviews	March 1, 2009
Publication Date	June 1, 2009

Guest Editors

Edson Denis Leonel, Departamento de Estatística, Matemática Aplicada e Computação, Instituto de Geociências e Ciências Exatas, Universidade Estadual Paulista, Avenida 24A, 1515 Bela Vista, 13506-700 Rio Claro, SP, Brazil ; edleonel@rc.unesp.br

Alexander Loskutov, Physics Faculty, Moscow State University, Vorob'evy Gory, Moscow 119992, Russia; loskutov@chaos.phys.msu.ru