

SOLVABILITY OF A FOURTH ORDER BOUNDARY VALUE PROBLEM WITH PERIODIC BOUNDARY CONDITIONS

CHAITAN P. GUPTA

Department of Mathematical Sciences
Northern Illinois University
DeKalb, IL 60115

(Received May 28, 1987 and in revised form July 16, 1987)

ABSTRACT. Fourth order boundary value problems arise in the study of the equilibrium of an elastaic beam under an external load. The author earlier investigated the existence and uniqueness of the solutions of the nonlinear analogues of fourth order boundary value problems that arise in the equilibrium of an elastic beam depending on how the ends of the beam are supported. This paper concerns the existence and uniqueness of solutions of the fourth order boundary value problems with periodic boundary conditions.

KEY WORDS AND PHRASES. fourth order boundary value problem, periodic boundary conditions, linear eigenvalue problem, Leray-Schauder continuation theorem, equilibrium of an elastic beam, non-trivial kernel.

AMS SUBJECT CLASSIFICATION: 34B15, 34C25

1. INTRODUCTION

Fourth order boundary value problems arise in the study of the equilibrium of an elastic beam under an external load, (e.g., see [1], [2], [3]) where the existence, uniqueness and iterative methods to construct the solutions have been studied extensively. The purpose of this paper is to study the fourth order boundary value problem with periodic boundary conditions:

$$\begin{aligned} \frac{d^4u}{dx^4} + f(u)u' + g(x,u) &= e(x), \quad x \in [0, 2\pi], \\ u(0) - u(2\pi) &= u'(0) - u'(2\pi) = u''(0) - u''(2\pi) \\ &= u'''(0) - u'''(2\pi) = 0, \end{aligned} \tag{1.1}$$

where $f: \mathbb{R} \rightarrow \mathbb{R}$ is continuous and $g: [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R}$ satisfies Caratheodory's conditions with $e \in L^1[0, 2\pi]$.

We note that the fourth order linear eigenvalue problem

$$\frac{d^4u}{dx^4} = \lambda u, \tag{1.2}$$

$u(0) - u(2\pi) = u'(0) - u'(2\pi) = u''(0) - u''(2\pi) = u'''(0) - u'''(2\pi) = 0$,
has $\lambda = n^4$, $n = 0, 1, 2, \dots$ as eigenvalues. Now the problem (1.1) is at resonance since the linear operator $Lu = \frac{d^4u}{dx^4}$ with $D(L) = \{u \in H^3(0, 2\pi) \mid u(0) = u(2\pi)$,

$u'(0) = u'(2\pi)$, $u''(0) = u''(2\pi)$, $u'''(0) = u'''(2\pi)$ has a non-tirival kernel. (See end of this introduction for the definition of $H^3(0,2\pi)$.) We shall prove that the boundary value problem (1.1) has at least one solution if $\int_0^{2\pi} e(x)dx = 0$, and there exists a constant $\rho > 0$ such that $g(x,u)u \geq 0$ for $|u| \geq \rho$. To prove the existence of a solution for the boundary value problem

$$-\frac{d^4u}{dx^4} + \alpha u' + g(x,u) = e(x), \quad x \in [0,2\pi],$$

$$\begin{aligned} u(0) - u(2\pi) &= u'(0) - u'(2\pi) = u''(0) - u''(2\pi) \\ &= u'''(0) - u'''(2\pi) = 0, \end{aligned} \quad (1.3)$$

we also need to assume that

$$\limsup_{|u| \rightarrow \infty} \frac{g(x,u)}{u} = \beta < 1, \quad \text{uniformly for a.e. } x \in [0,2\pi].$$

This is because the second eigenvalue $\lambda = 1$ of the linear eigenvalue problem (1.2) interferes with the non-linearity $g(x,u)$ in (1.3). The question of asymptotic conditions in which non-linearity $g(x,u)$ in (1.3) can interact with infinitely many eigenvalues of the eigenvalue problem (1.2) will be studied in a forthcoming paper [4].

To obtain the existence of solutions for (1.1) and (1.3), we use Mawhin's version of Leray Schauder continuation theorem as given in [5], [6], [7]. We also show that in case $f = \alpha$, where α is a constant, any two solutions of the boundary value problem (1.1), (respectively, (1.3)), differ by a constant and have a unique solution when, for example, $g(x,u)$ is strictly increasing in u for a.e. x in $[0,2\pi]$.

We note that in addition to using the classical spaces $C([0,2\pi])$, $C^k([0,2\pi])$, and $L^k(0,2\pi)$ and $L^\infty(0,2\pi)$ of continuous, k -times continuously differentiable, measurable real-valued functions whose k -th power of the absolute value is Lebesgue integrable or measurable functions which are essentially-bounded on $[0,2\pi]$ we shall use the Sobolev space $H^3(0,2\pi)$ defined by

$$H^3(0,2\pi) = \{u: [0,2\pi] \rightarrow \mathbb{R} \mid u, u', u'' \text{ abs. cont. on } [0,2\pi], \\ u''' \in L^2(0,2\pi)\}.$$

Also for $u \in L^1(0,2\pi)$ we define $\bar{u} = \frac{1}{2\pi} \int_0^{2\pi} u(x)dx$.

2. MAIN RESULTS

Let X, Y denote the Banach spaces $X = C^1[0,2\pi]$, $Y = L^1(0,2\pi)$ with usual norms and let H denote the Hilbert space $L^2(0,2\pi)$. Let Y_2 be the subspace of Y defined by

$$Y_2 = \{u \in Y \mid u = \text{constant a.e. on } [0,2\pi]\},$$

and let Y_1 be the subspace of Y such that $Y = Y_1 \oplus Y_2$. (\oplus denotes the direct sum.) We note that for $u \in Y$ we can write

$$u(x) = (u(x) - \frac{1}{2\pi} \int_0^{2\pi} u(t)dt) + \frac{1}{2\pi} \int_0^{2\pi} u(t)dt, \quad x \in [0,2\pi].$$

We define the canonical projection operators $P: Y \rightarrow Y_1$; $Q: Y \rightarrow Y_2$ as follows

$$P(u) = u(x) - \frac{1}{2\pi} \int_0^{2\pi} u(t)dt,$$

$$Q(u) = \frac{1}{2\pi} \int_0^{2\pi} u(t)dt,$$

for $u \in Y$. Clearly, $Q = I - P$, where I denotes the identity mapping on Y , and the projection operators P and Q are continuous. Now let $X_2 = X \cap Y_2$.

Clearly X_2 is a closed subspace of X . Let X_1 be the closed subspace of X such that $X = X_1 \oplus X_2$. We note that $P|X: X \rightarrow X_1$, $Q|X: X \rightarrow X_2$ are continuous. Similarly, we obtain $H = H_1 \oplus H_2$ and continuous projections $P|H: H \rightarrow H_1$, $Q|H: H \rightarrow H_2$. In the following, X , Y , H , P , Q , etc. will refer to Banach spaces, Hilbert space and the projections as defined above and we shall not distinguish between P , $P|X$, $P|H$ (resp. Q , $Q|X$, $Q|H$) and depend on the context for proper meaning.

Also for $u \in X$, $v \in Y$ let $(u, v) = \frac{1}{2\pi} \int_0^{2\pi} u(x)v(x)dx$ denote the duality pairing between X and Y . We note that for $u \in X$, $v \in Y$ where $u = Pu + Qu$, $v = Pv + Qv$, we have

$$(u, v) = (Pu, Pv) + (Qu, Qv).$$

Define a linear operator $L: D(L) \subset X \rightarrow Y$ by setting

$$\begin{aligned} D(L) = \{u \in H^3(0, 2\pi) &| u(0) = u(2\pi), u'(0) = u'(2\pi), \\ u''(0) = u''(2\pi), u'''(0) = u'''(2\pi)\} \end{aligned} \quad (2.1)$$

and for $u \in D(L)$,

$$Lu = \frac{d^4 u}{dx^4}. \quad (2.2)$$

Now, for $u \in D(L)$ we see using integration by parts and Wirtinger's inequality ([8]) that

$$(Lu, u) = \int_0^{2\pi} \frac{d^4 u}{dx^4} u dx = \int_0^{2\pi} u''^2 dx \geq \int_0^{2\pi} [(Pu)(x)]^2 dx \geq 0. \quad (2.3)$$

LEMMA 2.1: - For a given $\alpha \in \mathbb{R}$ and $h \in Y_1$, i.e. $h \in L^1(0, 2\pi)$ with $\bar{h} = Qh = 0$, the linear boundary value problem

$$\frac{d^4 u}{dx^4} + \alpha u' = h(x), \quad x \in [0, 2\pi], \quad (2.4)$$

$$u(0) = u(2\pi), \quad u'(0) = u'(2\pi), \quad u''(0) = u''(2\pi), \quad u'''(0) = u'''(2\pi),$$

has a unique solution $u(x)$ with $\bar{u} = Qu = 0$.

PROOF:- Let us set $\omega = \cos \frac{2\pi}{3} + i \sin \frac{2\pi}{3}$, $i = \sqrt{-1}$, so that $\alpha^{1/3}$, $\omega \alpha^{1/3}$, $\omega^2 \alpha^{1/3}$ are the three cube roots of $\alpha \in \mathbb{R}$. For $x \in [0, 2\pi]$, we define

$$\begin{aligned} v_1(x) &= \int_0^x h(t)dt, \quad v_2(x) = e^{-\alpha^{1/3} \omega x} \int_0^x v_1(t) e^{\alpha^{1/3} \omega t} dt, \\ v_3(x) &= e^{-\alpha^{1/3} \omega x} \int_0^x v_2(t) e^{\alpha^{1/3} \omega t} dt, \quad v(x) = e^{-\alpha^{1/3} x} \int_0^x e^{\alpha^{1/3} t} v_3(t) dt. \end{aligned}$$

Then $u(x) = C_1 + C_2 e^{-\alpha^{1/3} x} + C_3 e^{-\alpha^{1/3} \omega x} + C_4 e^{-\alpha^{1/3} \omega^2 x} + v(x)$ is such that $Ru(u(x))$ is a general solution of the equation (2.4).

Next, we compute C_1, C_2, C_3, C_4 using the boundary conditions in (2.4) and and the condition $\bar{u} = \frac{1}{2\pi} \int_0^{2\pi} u(x)dx = 0$. C_2, C_3, C_4 are computed uniquely from the three linearly independent equations

$$C_2 + C_3 + C_4 = C_2 e^{-\alpha^{1/3} 2\pi} + C_3 e^{-\alpha^{1/3} 2\pi \omega} + C_4 e^{-\alpha^{1/3} 2\pi \omega^2} + v(2\pi),$$

$$C_2 + \omega C_3 + \omega^2 C_4 = C_2 e^{-\alpha^{1/3} 2\pi} + \omega C_3 e^{-\alpha^{1/3} 2\pi\omega} + \omega^2 C_4 e^{-\alpha^{1/3} 2\pi\omega^2} - \alpha^{-1/3} v'(2\pi),$$

$$C_2 + \omega^2 C_3 + \omega C_4 = C_2 e^{-\alpha^{1/3} 2\pi} + \omega^2 C_3 e^{-\alpha^{1/3} 2\pi\omega} + \omega C_4 e^{-\alpha^{1/3} 2\pi\omega^2} + \alpha^{-2/3} v''(2\pi).$$

The constant C_1 is computed uniquely using the condition $\bar{u} = \frac{1}{2\pi} \int_0^{2\pi} u(x) dx = 0$.

In this way we get $R1$ $u(x)$ as the unique solution for (2.4). //

For $h \in Y_1$, i.e. $h \in L^1(0, 2\pi)$ with $\bar{h} = \frac{1}{2\pi} \int_0^{2\pi} h(x) dx = 0$; let $u = Kh$ be the unique solution of the problem

$$\frac{d^4 u}{dx^4} = h(x), \quad x \in [0, 2\pi],$$

$$u(0) = u(2\pi), \quad u'(0) = u'(2\pi), \quad u''(0) = u''(2\pi), \quad u'''(0) = u'''(2\pi),$$

such that $\bar{u} = \frac{1}{2\pi} \int_0^{2\pi} u(t) dt = 0$. It is immediate that the linear mapping

$K: Y_1 \rightarrow X_1$ is bounded and for $u \in Y$,

$$KP(u) \in D(L), \quad LK P(u) = P(u), \quad \text{and} \quad (KP(u), P(u)) \geq 0. \quad (2.5)$$

Let $f: \mathbb{R} \rightarrow \mathbb{R}$ be continuous and let $g: [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R}$, $(x, u) \rightarrow g(x, u)$ be such that $g(., u)$ is measurable on $[0, 2\pi]$ for each $u \in \mathbb{R}$ and $g(x, .)$ is continuous on \mathbb{R} for almost each $x \in [0, 2\pi]$. Assume, moreover, that for each $r > 0$ there exists an $\alpha_r \in L^1(0, 2\pi)$ such that $|g(x, u)| \leq \alpha_r(x)$ for a.e. $x \in [0, 2\pi]$ and all $u \in [-r, r]$. Such a g will be said to satisfy Caratheodory's conditions.

Now define $N: X \rightarrow Y$ by setting

$$(Nu)(x) = f(u(x)) u'(x) + g(x, u(x)), \quad x \in [0, 2\pi],$$

for $u \in X$. It follows easily from Arzela-Ascoli theorem that $KPN: X \rightarrow X_1$ is a well-defined compact mapping and $QN: X \rightarrow X_2$ is bounded.

For $e(x) \in Y = L^1(0, 2\pi)$, the boundary value problem (1.1) now reduces to the functional equation

$$Lu + Nu = e, \quad (2.6)$$

in X with $e \in Y$, given.

THEOREM 2.2:- Let $f: \mathbb{R} \rightarrow \mathbb{R}$ be continuous and let $g: [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R}$ satisfy Caratheodory's conditions. Assume that there exist real numbers a, A, r and R with $a \leq A$ and $r < 0 < R$ such that

$$g(x, u) \geq A, \quad (2.7)$$

for a.e. $x \in [0, 2\pi]$ and all $u \leq R$; and

$$g(x, u) \leq a, \quad (2.8)$$

for a.e. $x \in [0, 2\pi]$ and all $u \leq r$. Then the boundary value problem (1.1) has at least one solution for each given $e \in L^1(0, 2\pi)$ with

$$a \leq \bar{e} \leq A. \quad (2.9)$$

PROOF:- Define $g_1: [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R}$ by $g_1(x, u) = g(x, u) - \frac{1}{2}(a + A)$ and $e_1 \in L^1(0, 2\pi)$ by $e_1(x) = e(x) - \frac{1}{2}(a + A)$, so that, for a.e. $x \in [0, 2\pi]$,

using (2.7), (2.8), (2.9) we have

$$g_1(x, u) \geq \frac{1}{2} (A - a) \geq 0 \quad \text{if } u \geq R, \quad (2.10)$$

$$g_1(x, u) \leq \frac{1}{2} (a - A) \leq 0 \quad \text{if } u \leq r, \quad (2.11)$$

and

$$\frac{1}{2}(a - A) \leq \bar{e}_1 \leq \frac{1}{2}(A - a). \quad (2.12)$$

Clearly, the boundary value problem (1.1) is equivalent to

$$\frac{d^4 u}{dx^4} + f(u)u' + g_1(x, u(x)) = e_1(x), \quad x \in [0, 2\pi], \quad (2.13)$$

$$u(0) = u(2\pi), \quad u'(0) = u'(2\pi), \quad u''(0) = u''(2\pi), \quad u'''(0) = u'''(2\pi).$$

Let $N: X \rightarrow Y$ be defined by

$$(Nu)(x) = f(u(x))u'(x) + g_1(x, u(x)), \quad x \in [0, 2\pi], \quad (2.14)$$

for $u \in X$. We then see, as above, that $KPN: X \rightarrow X_1$ is a well-defined compact mapping. $QN: X \rightarrow X_2$ is bounded and the boundary value problem (2.13) is equivalent to the functional equation,

$$Lu + Nu = e_1, \quad (2.15)$$

in X with $e_1 \in Y$. Setting, $\tilde{e}_1 = KPe$, we see that to solve the functional equation (2.15) it suffices to solve the system of equations

$$\begin{aligned} Pu + KPNu &= \tilde{e}_1, \\ QNu &= \bar{e}_1, \end{aligned} \quad (2.16)$$

$u \in X$. Indeed, if $u \in X$ is a solution of (2.16) then $u \in D(L)$ and

$$\begin{aligned} LPu + LKPNu &= Lu + PNu = \tilde{e}_1 = Pe_1, \\ QNu &= \bar{e}_1 = Qe_1, \end{aligned}$$

which gives on adding that $Lu + Nu = e_1$.

Now, (2.16) is clearly equivalent to the single equation

$$Pu + QNu + KPNu = \tilde{e}_1 + \bar{e}_1, \quad (2.17)$$

which has the form of a compact perturbation of the Fredholm operator P of index zero. We can therefore apply the version given in [6] (Theorem 1, Corollary 1) or [5] (Theorem IV.4) or [7] of the Leray-Schauder Continuation theorem which ensures the existence of a solution for (2.17) if the set of solutions of the family of equations,

$$Pu + (1-\lambda)Qu + \lambda QNu + \lambda KPNu = \lambda \tilde{e}_1 + \lambda \bar{e}_1, \quad \lambda \in (0, 1), \quad (2.18)$$

is, a priori, bounded in X by a constant independent of λ . Notice that (2.18) is equivalent to the system of equations,

$$Pu + \lambda KPNu = \lambda \tilde{e}_1, \quad (2.19)$$

$$(1-\lambda)Qu + \lambda QNu = \lambda \bar{e}_1, \quad \lambda \in (0, 1).$$

Let for $\lambda \in (0, 1)$, $u_\lambda \in X$ be a solution of (2.19) so that

$$\begin{aligned} \mathbf{P}u_\lambda + \lambda \mathbf{KPNu}_\lambda &= \lambda \tilde{\mathbf{e}}_1, \\ (1-\lambda)\mathbf{Qu}_\lambda + \lambda \mathbf{QNu}_\lambda &= \lambda \bar{\mathbf{e}}_1. \end{aligned} \quad (2.20)$$

The second equation in (2.20) can now be written as

$$(1-\lambda) \int_0^{2\pi} u_\lambda(x) dx + \frac{\lambda}{2\pi} \int_0^{2\pi} g_1(x, u_\lambda(x)) dx = \lambda \bar{\mathbf{e}}_1.$$

So, if $u_\lambda(x) \geq R$ for $x \in [0, 2\pi]$ we have, using (2.10), (2.12) that

$$0 < (1 - \lambda) R + \frac{\lambda}{2} (A - a) \leq \frac{\lambda}{2} (A - a),$$

i.e.

$$0 < (1 - \lambda) R \leq 0, \text{ a contradiction.}$$

Similarly if $u_\lambda(x) \leq r$ for $x \in [0, 2\pi]$ leads to a contradiction. Hence, there exists a $\tau_\lambda \in [0, 2\pi]$ such that

$$r < u_\lambda(\tau_\lambda) < R. \quad (2.21)$$

Now, for $x \in [0, 2\pi]$ we have

$$u_\lambda(x) = u_\lambda(\tau_\lambda) + \int_{\tau_\lambda}^x u'_\lambda(s) ds,$$

so that

$$\begin{aligned} |u_\lambda(x)| &\leq \max(R, -r) + (2\pi)^{1/2} \left(\int_0^{2\pi} (u'_\lambda(s))^2 ds \right)^{1/2} \\ &\leq \max(R, -r) + (2\pi)^{1/2} \left(\int_0^{2\pi} (u''_\lambda(s))^2 ds \right)^{1/2} \end{aligned}$$

since $u_\lambda \in D(L)$, Wirtinger's inequality applies. Thus,

$$\|u_\lambda\|_X \leq C_1 \|u''_\lambda\|_H + C_2, \quad (2.22)$$

for some constants C_1, C_2 independent of λ .

Next, the first equation in (2.20) gives that

$$\mathbf{L}P u_\lambda + \lambda \mathbf{LKNu}_\lambda = \lambda \tilde{\mathbf{e}}_1,$$

i.e.

$$\mathbf{L}u_\lambda + \lambda \mathbf{PNu}_\lambda = \lambda \mathbf{Pe}_1. \quad (2.23)$$

From (2.23) and the second equation in (2.20), we get

$$\begin{aligned} (\mathbf{L}u_\lambda, \mathbf{P}u_\lambda) + \lambda (\mathbf{P}u_\lambda, \mathbf{P}u_\lambda) &= \lambda (\mathbf{Pe}_1, \mathbf{P}u_\lambda), \\ (1-\lambda)(\mathbf{Qu}_\lambda, \mathbf{Qu}_\lambda) + \lambda (\mathbf{QNu}_\lambda, \mathbf{Qu}_\lambda) &= \lambda (\bar{\mathbf{e}}_1, \mathbf{Qu}_\lambda). \end{aligned} \quad (2.24)$$

We next note that our assumptions on g_1 and (2.10), (2.12) imply that there is a constant C_3 , independent of λ such that for $u \in X$,

$$(\mathbf{Nu}, u) \geq -C_3,$$

and $(\mathbf{Lu}_\lambda, \mathbf{Pu}_\lambda) = (\mathbf{Lu}_\lambda, u_\lambda) = \int_0^{2\pi} (u''_\lambda)^2 = \|u''_\lambda\|_H^2$ since (2.3) holds. Using this we get on adding the equations in (2.24) that

$$\begin{aligned}
\|u''_\lambda\|_H^2 - C_3 &\leq (Lu_\lambda, u_\lambda) + (1-\lambda)(Qu_\lambda, Qu_\lambda) + \lambda(Nu_\lambda, u_\lambda) \\
&= \lambda(Pe_1, Pu_\lambda) + \lambda(\bar{e}_1, Qu_\lambda) \\
&\leq C_4 \|u_\lambda\|_X \\
&\leq C_4 C_1 \|u''_\lambda\|_H + C_4 C_2,
\end{aligned}$$

where C_4 is a constant independent of λ . Accordingly, there is a constant C_5 , independent of λ , such that

$$\|u''_\lambda\|_H \leq C_5,$$

which implies, using (2.22) that

$$\|u_\lambda\|_X \leq C_1 C_5 + C_2 \equiv C.$$

We have thus proved that the set of solutions of the family of equations (2.18) is bounded in X by a constant independent of $\lambda \in (0,1)$. Hence the theorem. //

REMARK 2.3:- If we take $a = A = 0$ in Theorem 2.2, then we immediately obtain the assertion made in the introduction concerning the boundary value problem (1.1).

Now, to study the boundary value problem (1.3) we define, for a given $\alpha \in \mathbb{R}$, a linear operator $L_\alpha : D(L_\alpha) \subset X \rightarrow Y$ by setting

$$\begin{aligned}
D(L_\alpha) &= \{u \in H^3(0, 2\pi) \mid u(0) = u(2\pi), u'(0) = u'(2\pi), u''(0) = u''(2\pi), \\
&\quad u'''(0) = u'''(2\pi)\}
\end{aligned} \tag{2.25}$$

and for $u \in D(L_\alpha)$,

$$L_\alpha u = -\frac{d^4 u}{dx^4} + \alpha u'. \tag{2.26}$$

It follows, using integration by parts and Wirtinger's inequality, ([8]), that

$$\begin{aligned}
(L_\alpha u, u) &= - \int_0^{2\pi} \frac{d^4 u}{dx^4} u dx + \alpha \int_0^{2\pi} u' u dx \\
&= - \int_0^2 (u'')^2 dx \geq - \int_0^{2\pi} \left(\frac{d^4 u}{dx^4} \right)^2 dx \\
&\geq - \|L_\alpha u\|_H^2.
\end{aligned} \tag{2.27}$$

We, next, use lemma 2.1 to define a bounded linear mapping $K_\alpha : Y_1 \rightarrow X_1$ by setting $u = K_\alpha h$ for a given $h \in Y_1$, where $u \in X_1$ (so that $\bar{u} = Qu = 0$) is the unique solution of the boundary value problem

$$-\frac{d^4 u}{dx^4} + \alpha u' = h(x), \quad x \in [0, 2\pi], \tag{2.28}$$

$$u(0) = u(2\pi), \quad u'(0) = u'(2\pi), \quad u''(0) = u''(2\pi), \quad u'''(0) = u'''(2\pi).$$

The bounded linear mapping $K_\alpha : Y_1 \rightarrow X_1$ defined in this way has the following properties:

(i) for $u \in Y$, $K_\alpha P(u) \in D(L_\alpha)$, $L_\alpha K_\alpha P(u) = P(u)$ and

$$(K_\alpha P(u), P(u)) \geq - \|P(u)\|_H^2, \quad (\text{in view of (2.27)}); \tag{2.29}$$

(ii) if $g : [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R}$ satisfies Caratheodory's conditions and $N : X \rightarrow Y$ is defined by setting

$$(Nu)(x) = g(x, u(x)), \quad x \in [0, 2\pi]$$

then $K_\alpha^P N : X \rightarrow X_1$ is a well-defined compact mapping and $QN : X \rightarrow X_2$ is bounded.

Theorem 2.4: Let $\alpha \in \mathbb{R}$ be given and $g : [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R}$ satisfy Caratheodory's conditions. Assume that there exist real numbers a, A, r, R with $a \leq A$, and $r < 0 < R$ such that

$$g(x, u) \geq A, \quad (2.30)$$

for a.e. $x \in [0, 2\pi]$, and all $u \geq R$; and

$$g(x, u) \leq a, \quad (2.31)$$

for a.e. $x \in [0, 2\pi]$, and all $u \leq r$. Suppose, further, that

$$\limsup_{|u| \rightarrow \infty} \left| \frac{g(x, u)}{u} \right| = \beta < 1 \quad (2.32)$$

uniformly for a.e. $x \in [0, 2\pi]$. Then the boundary value problem (1.3) has at least one solution for each given $e \in L^2[0, 2\pi]$ with

$$a \leq \bar{e} \leq A. \quad (2.33)$$

Proof:- As in the proof of Theorem 2.2, define $g_1 : [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R}$ by $g_1(x, u) = g(x, u) - \frac{1}{2}(a + A)$ and $e_1 \in L^2(0, 2\pi)$ by $e_1(x) = e(x) - \frac{1}{2}(a + A)$. Then for a.e. $x \in [0, 2\pi]$,

$$g_1(x, u) \geq \frac{1}{2}(A - a) \geq 0 \quad \text{if } u \geq R, \quad (2.34)$$

$$g_1(x, u) \leq \frac{1}{2}(a - A) \leq 0 \quad \text{if } u \leq r, \quad (2.35)$$

$$\limsup_{|u| \rightarrow \infty} \left| \frac{g_1(x, u)}{u} \right| = \beta < 1, \quad (2.36)$$

uniformly, and

$$\frac{1}{2}(a - A) \leq \bar{e}_1 \leq \frac{1}{2}(A - a). \quad (2.37)$$

Also the boundary value problem (1.3) is equivalent to

$$-\frac{d^4 u}{dx^4} + \alpha u' + g_1(x, u) = e_1(x), \quad x \in [0, 2\pi], \quad (2.38)$$

$$u(0) = u(2\pi), \quad u'(0) = u'(2\pi), \quad u''(0) = u''(2\pi), \quad u'''(0) = u'''(2\pi).$$

Next, let $N : X \rightarrow Y$ be defined by

$$Nu(x) = g_1(x, u(x)), \quad x \in [0, 2\pi],$$

for $u \in X$. Choosing, now, $\varepsilon > 0$ such that $\beta + \varepsilon < 1$, we see, using the fact that g , satisfies Caratheodory's conditions and (2.34), (2.35), (2.36), that there exists a constant $C(\varepsilon) > 0$ such that

$$(Nu, u) \geq \frac{1}{\beta + \varepsilon} \|Nu\|_H^2 - C(\varepsilon), \quad (2.39)$$

for $u \in X$. Also, $K_\alpha^P N : X \rightarrow X_1$ is a well-defined compact mapping and $QN : X \rightarrow X_2$ is bounded.

Again, we see as in the proof of Theorem 2.2, that the boundary value problem (2.38) is equivalent to the system of equations

$$\begin{aligned} \mathbf{P}u + K_\alpha \mathbf{P}N\mathbf{u} &= \tilde{\mathbf{e}}_1 = K_\alpha \mathbf{P}\mathbf{e}_1, \\ \mathbf{Q}N\mathbf{u} &= \bar{\mathbf{e}}_1. \end{aligned} \quad (2.40)$$

Further, it suffices to prove that the set of solutions of the family of equations

$$\begin{aligned} \mathbf{P}u + \lambda K_\alpha \mathbf{P}N\mathbf{u} &= \lambda \tilde{\mathbf{e}}_1 \\ (1 - \lambda) \mathbf{Q}u + \lambda \mathbf{Q}N\mathbf{u} &= \lambda \bar{\mathbf{e}}_1, \quad \lambda \in (0,1) \end{aligned} \quad (2.41)$$

is, a priori, bounded in X by a constant independent of $\lambda \in (0,1)$.

Let, now, for $\lambda \in (0,1)$, $u_\lambda \in X$ be a solution of (2.41) so that

$$\begin{aligned} \mathbf{P}u_\lambda + \lambda K_\alpha \mathbf{P}N\mathbf{u}_\lambda &= \lambda \mathbf{e}_1, \\ (1 - \lambda) \mathbf{Q}u_\lambda + \lambda \mathbf{Q}N\mathbf{u}_\lambda &= \lambda \bar{\mathbf{e}}_1. \end{aligned} \quad (2.42)$$

It, now, follows from the second equation in (2.42), in a manner similar to deriving the estimate (2.22) in the proof of Theorem 2.2, that

$$\|Qu_\lambda\| \leq \|u_\lambda\|_X \leq C_1 \|L_\alpha u_\lambda\|_H + C_2, \quad (2.43)$$

for some constants C_1, C_2 independent of $\lambda \in (0,1)$.

Also, we have from (2.42) that

$$(\mathbf{P}u_\lambda, \mathbf{P}N\mathbf{u}_\lambda) + \lambda(K_\alpha \mathbf{P}N\mathbf{u}_\lambda, \mathbf{P}N\mathbf{u}_\lambda) = \lambda(\tilde{\mathbf{e}}_1, \mathbf{P}N\mathbf{u}_\lambda),$$

$$(1 - \lambda) \|Qu_\lambda\|^2 + \lambda(Qu_\lambda, QN\mathbf{u}_\lambda) = \lambda(\bar{\mathbf{e}}_1, Qu_\lambda).$$

These equations then give us, in view of (2.29) and (2.39), that

$$\begin{aligned} \frac{1}{\beta + \varepsilon} \|Nu_\lambda\|_H^2 - \|P\mathbf{N}u_\lambda\|_H^2 - C(\varepsilon) &\leq (Nu_\lambda, u_\lambda) + \lambda(K_\alpha \mathbf{P}N\mathbf{u}_\lambda, \mathbf{P}N\mathbf{u}_\lambda) \\ &\leq (\tilde{\mathbf{e}}_1, \mathbf{P}N\mathbf{u}_\lambda) + (\bar{\mathbf{e}}_1, Qu_\lambda). \end{aligned}$$

Using, now, the facts that $\|\mathbf{P}v\|_H \leq \|v\|_H$ for $v \in X$, and $\beta + \varepsilon < 1$, we see that these exist constants C_3, C_4 independent of $\lambda \in (0,1)$ such that

$$\|Nu_\lambda\|_H \leq C_3 \|Qu_\lambda\|_H^{1/2} + C_4. \quad (2.44)$$

Now, the first equation in (2.42) gives that

$$L_\alpha u_\lambda + \lambda P\mathbf{N}u_\lambda = \lambda \mathbf{P}\mathbf{e}_1,$$

so that

$$\begin{aligned} \|L_\alpha u_\lambda\|_H &\leq \lambda \|\mathbf{P}\mathbf{e}_1 - P\mathbf{N}u_\lambda\|_H \leq \|\mathbf{P}\mathbf{e}_1\|_H + \|P\mathbf{N}u_\lambda\|_H \\ &\leq \|\mathbf{P}\mathbf{e}_1\|_H + \|Nu_\lambda\|_H \\ &\leq C_3 \|Qu_\lambda\|_H^{1/2} + C_4 + \|\mathbf{P}\mathbf{e}_1\|_H. \end{aligned}$$

(2.43) and (2.45) now imply that there exist a constant C_5 , independent of $\lambda \in (0,1)$, such that

$$\|Qu_\lambda\| \leq c_5$$

and

$$\|u_\lambda\|_X \leq c_1 c_3 \sqrt{c_5} + c_1 c_4 + c_1 \|p_{e_1}\|_H + c_2 = c.$$

This completes the proof of the theorem //.

Remark 2.5:- The analogue of Theorem 2.4 when α in (1.3) is replaced by $f(u)$, where $f : \mathbb{R} \rightarrow \mathbb{R}$ is a given continuous function will be treated in a forthcoming paper [4].

Remark 2.6:- If $f(u) \equiv \alpha$, $\alpha \in \mathbb{R}$ given and $g(x,u)$ is strictly increasing in u for a.e. $x \in [0, 2\pi]$ then it is easy to see that the boundary value problem (1.1) has exactly one solution. Similarly if $g(x,u)$ is strictly increasing in u and there is a $\beta < 1$, such that

$$(g(x, u_1) - g(x, u_2))(u_1 - u_2) \geq \beta(g(x, u_1) - g(x, u_2))^2,$$

for a.e. x in $[0, 2\pi]$, then the boundary value problem (1.3) has exactly one solution.

Remark 2.7:- If we take $a = A = 0$ in Theorem 2.4, we immediately obtain the assertion concerning the boundary value problem (1.3) in the introduction.

Acknowledgement:- The author is grateful to the referee for his suggestions to add references [1], [2], [3] to the bibliography and for editorial improvement of this paper.

REFERENCES

1. AGARWAL, R.P. Boundary Value Problems for Higher Order Differential Equations, World Scientific, Singapore, Philadelphia, PA, 1986.
2. AGARWAL, R.P., CHOW, Y.M. Iterative Methods for a Fourth Order Boundary Value Problem, *Jour. Comp. Appl. Math.* 10 (1984), 203-207.
3. USMANI, R.A. Solving Boundary Value Problems in Plate Deflection Theory, *Simulation*, December 1981, 195-206.
4. GUPTA, C.P. Asymptotic conditions for the solvability of a fourth order boundary value problem with periodic boundary conditions. (Under preparation).
5. MAWHIN, J. Topological degree methods in nonlinear boundary value problems. CBMS - Regional Conference Series in Maths. No. 40, (1979) American Mathematical Society, Providence, R.I.
6. MAWHIN, J. Landesman - Lazer type problems for nonlinear equations. *Conf. Sem. Mat. Univ. Bari* no. 147, (1977).
7. MAWHIN, J. Compacite', monotonie, et convexité dans l'étude de problèmes aux limites semi-linéaires. *Sem. Anal. Moderne* No. 19, (1981) Université de Sherbrooke, Québec, CANADA.
8. HARDY, G.H., LITTLEWOOD, J.E., and POLYA, G. Inequalities, Cambridge University Press, London and New York, 1952.
9. GUPTA, C.P. Existence and uniqueness theorems for the bending of an elastic beam equation. *Applicable Analysis* (to appear).
10. GUPTA, C.P. Existence and uniqueness theorems for the bending of an elastic beam equation at resonance. *Jour. Math. Anal. & Appl.* (to appear).

Special Issue on Time-Dependent Billiards

Call for Papers

This subject has been extensively studied in the past years for one-, two-, and three-dimensional space. Additionally, such dynamical systems can exhibit a very important and still unexplained phenomenon, called as the Fermi acceleration phenomenon. Basically, the phenomenon of Fermi acceleration (FA) is a process in which a classical particle can acquire unbounded energy from collisions with a heavy moving wall. This phenomenon was originally proposed by Enrico Fermi in 1949 as a possible explanation of the origin of the large energies of the cosmic particles. His original model was then modified and considered under different approaches and using many versions. Moreover, applications of FA have been of a large broad interest in many different fields of science including plasma physics, astrophysics, atomic physics, optics, and time-dependent billiard problems and they are useful for controlling chaos in Engineering and dynamical systems exhibiting chaos (both conservative and dissipative chaos).

We intend to publish in this special issue papers reporting research on time-dependent billiards. The topic includes both conservative and dissipative dynamics. Papers discussing dynamical properties, statistical and mathematical results, stability investigation of the phase space structure, the phenomenon of Fermi acceleration, conditions for having suppression of Fermi acceleration, and computational and numerical methods for exploring these structures and applications are welcome.

To be acceptable for publication in the special issue of Mathematical Problems in Engineering, papers must make significant, original, and correct contributions to one or more of the topics above mentioned. Mathematical papers regarding the topics above are also welcome.

Authors should follow the Mathematical Problems in Engineering manuscript format described at <http://www.hindawi.com/journals/mpe/>. Prospective authors should submit an electronic copy of their complete manuscript through the journal Manuscript Tracking System at <http://mts.hindawi.com/> according to the following timetable:

Manuscript Due	December 1, 2008
First Round of Reviews	March 1, 2009
Publication Date	June 1, 2009

Guest Editors

Edson Denis Leonel, Departamento de Estatística, Matemática Aplicada e Computação, Instituto de Geociências e Ciências Exatas, Universidade Estadual Paulista, Avenida 24A, 1515 Bela Vista, 13506-700 Rio Claro, SP, Brazil ; edleonel@rc.unesp.br

Alexander Loskutov, Physics Faculty, Moscow State University, Vorob'evy Gory, Moscow 119992, Russia; loskutov@chaos.phys.msu.ru